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SPECTRAL ANALYSIS OF THE
ENERGY EXCHANGE AT OWS NOVEMBER

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THESIS

Spectral Analysis of the
Energy Exchange at OWS NOVEMBER

by

Kevin Merle Rabe

March 1975

Thesis Advisor:

R. H. Bourke

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SECURITY CLASSIFICATION OF THIS PAGE (When Data Entered)

REPORT DOCUMENTATION PAGE		READ INSTRUCTIONS BEFORE COMPLETING FORM
1. REPORT NUMBER	2. GOVT ACCESSION NO.	3. RECIPIENT'S CATALOG NUMBER
4. TITLE (and Subtitle) Spectral Analysis of the Energy Exchange at OWS NOVEMBER		5. TYPE OF REPORT & PERIOD COVERED Master's Thesis March 1975
7. AUTHOR(s) Kevin Merle Rabe		6. PERFORMING ORG. REPORT NUMBER
9. PERFORMING ORGANIZATION NAME AND ADDRESS Naval Postgraduate School Monterey, CA 93940		10. PROGRAM ELEMENT, PROJECT, TASK AREA & WORK UNIT NUMBERS
11. CONTROLLING OFFICE NAME AND ADDRESS Naval Postgraduate School Monterey, CA 93940		12. REPORT DATE March 1975
14. MONITORING AGENCY NAME & ADDRESS (if different from Controlling Office)		13. NUMBER OF PAGES 108
16. DISTRIBUTION STATEMENT (of this Report)		15. SECURITY CLASS. (of this report) UNCLASSIFIED
17. DISTRIBUTION STATEMENT (of the abstract entered in Block 20, if different from Report)		
18. SUPPLEMENTARY NOTES		
19. KEY WORDS (Continue on reverse side if necessary and identify by block number) Spectral Analysis OWS November Heat Budget Sea-Surface Temperature		
20. ABSTRACT (Continue on reverse side if necessary and identify by block number) Oceanographic and meteorological data from Ocean Weather Station (OWS) NOVEMBER were analyzed by statistical and spectral methods in order to describe the nature and periodicities of the air-ocean energy exchange. Salinity data from 1968 through 1970 were compared with daily observed parameters commonly associated with changes in the salinity. Regression analysis showed surface salinity to be highly		

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Spectral Analysis of the
Energy Exchange at OWS NOVEMBER

by

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Submitted in partial fulfillment of the
requirements for the degree of

MASTER OF SCIENCE IN OCEANOGRAPHY

from the
NAVAL POSTGRADUATE SCHOOL
March 1975

ABSTRACT

Oceanographic and meteorological data from Ocean Weather Station (OWS) NOVEMBER were analyzed by statistical and spectral methods in order to describe the nature and periodicities of the air-ocean energy exchange. Salinity data from 1968 through 1970 were compared with daily observed parameters commonly associated with changes in the salinity. Regression analysis showed surface salinity to be highly correlated with surface pressure which is interpreted as the surface salinity responding to a baroclinic transport. At this location none of the classical parameters usually associated with fluctuations in the surface salinity were highly correlated.

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LIST OF SYMBOLS

SST	Sea-surface temperature ($^{\circ}\text{C}$)
PS	Surface pressure (mb)
W	Wind speed (m sec^{-1})
W^2	Wind speed squared
E	Evaporation (cm)
P	Precipitation (cm)
T_a	Air temperature ($^{\circ}\text{C}$)
T_w	Water temperature ($^{\circ}\text{C}$)
CL	Cloud cover, in tenths of sky covered
e_w	Vapor pressure over water (mb)
e_a	Vapor pressure of the air (mb)
C_D	The drag coefficient
S	Salinity ($^{\circ}/\text{oo}$)
QS	Solar radiation ($\text{cal cm}^{-2} \text{ day}^{-1}$)
QB	Back radiation ($\text{cal cm}^{-2} \text{ day}^{-1}$)
QE	Heat of evaporation ($\text{cal cm}^{-2} \text{ day}^{-1}$)
QC	Convective heat ($\text{cal cm}^{-2} \text{ day}^{-1}$)
QN	Net radiation ($\text{cal cm}^{-2} \text{ day}^{-1}$)

ACKNOWLEDGMENTS

The author wishes to thank Dr. Robert Bourke for his assistance and guidance, both during the course of this research and the preparation of this thesis, and also, Dr. Dale Leipper and Dr. Edward Thornton for their constructive review of the text. Thanks are also extended to various personnel at both the Environmental Prediction Research Facility and Fleet Numerical Weather Central through whose assistance the opportunity to conduct this research was made possible. They are: Captain G. D. Hamilton, past Commanding Officer at EPRF, Captain R. C. Sherar, present Commanding Officer, Dr. Taivo Laevastu, my immediate supervisor, and Captain C. Barreau, formerly Chief of Development at FNWC.

I. INTRODUCTION

A. OBJECTIVE

Sea-surface temperature is one of the most important parameters in any study of air-sea interaction. Anomalies of this parameter, measured over periods of a week to several months, have been linked to gross atmospheric effects. However, few studies have investigated the time response of the sea-surface temperature over long periods to those factors which govern its fluctuations.

Until the more recent decades such a study would have been limited by both instrumentation and general paucity of data. This resulted in reliance upon monthly or annual means to describe the state of the ocean. The advent of the ocean weather station program in the 1940's, and recently, the use of satellites has greatly increased the amount and frequency of data available to the researcher. Also, the advent of high speed computers has made the use of time series analysis routines relatively easy and inexpensive to use.

The past limitations made statistical studies of changes of surface temperature on time scales of a day or less, difficult. However, daily fluctuations can be quite informative. If advection is ignored, the sea-surface temperature at OWS NOVEMBER was found to fluctuate primarily in response to the day to day changes in air temperature, the amount of

insolation, wind-wave mixing, and the amount of cloud cover as well as the elevation of the sun. The magnitude of these diurnal fluctuations can be as large as the monthly fluctuations.

Sea-surface temperature fluctuations are important in many oceanographic parameters, such as the stability of the water column. This, in turn, is of great consequence in sound propagation. Anomalies of sea-surface temperature, likewise, have been linked to climatic variations over the ocean.

The salinity is important in that it affects the density of sea water. Thus, it affects both vertical stability and sound propagation. Together with the water temperature, one can infer transport characteristics of a region based on water mass analysis.

This study examines changes occurring at the surface of the ocean over time scales varying from one day to several years. The two parameters of primary interest are the sea-surface temperature and surface salinity. The area of study is Ocean Weather Station (OWS) NOVEMBER located at 30° North and 140° West. The data used are synoptic data collected by the United States Coast Guard at this position. The data are of two types: hydrographic Nansen cast data and shipboard meteorological observations.

The thesis is organized to first present correlations of the primary parameters versus various meteorological parameters. These correlations provide an indication of

those parameters which can be used as predictors for the sea-surface temperature and salinity.

Secondly, a spectral analysis of the data is made. From this, the synoptic atmospheric mechanisms which cause the fluctuations of the sea-surface temperature are identified. Periods of high spectral energy, along with their corresponding phase and coherence relationships are examined to gain insight into the time scales associated with each mechanism.

B. REVIEW OF THE LITERATURE

1. The Subtropic Region

Ocean Weather Station NOVEMBER, located at 30° North and 140° West, lies within what Tully (1964) defines as the Pacific Subtropical Region. Figure 1 shows its location in comparison to other major features of the North Pacific Ocean. The region is characterized by three dominant conditions: an excess of evaporation over precipitation throughout the year, the occurrence of convective mixing, and seasonal heating and cooling patterns which result in a pronounced seasonal thermocline.

The excess of evaporation over precipitation (E-P), has a distinct effect on the other two traits. Tully (1964) estimated E-P to be approximately $80 \text{ cm}^{-1} \text{ yr}$ in the vicinity of OWS NOVEMBER. This leads to a removal of fresh water from the surface and acts as a driving force for the convective mixing. This mixing then contributes to the erosion of the seasonal thermocline.

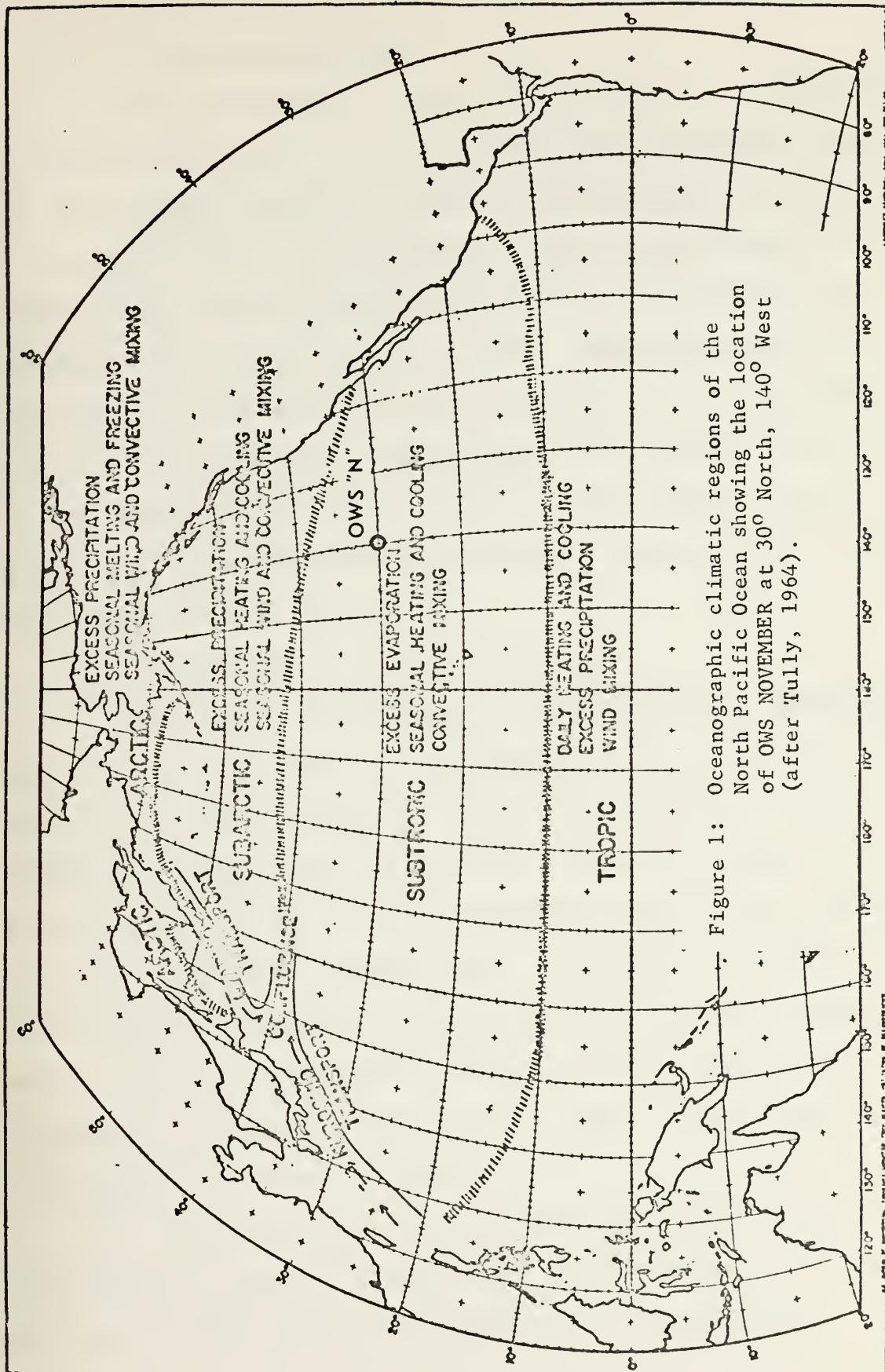


Figure 1: Oceanographic climatic regions of the North Pacific Ocean showing the location of OWS NOVEMBER at 30° North, 140° West (after Tully, 1964).

2. Surface Circulation

The subtropical region of the Pacific is dominated by the Pacific Gyre. Figures 2 and 3 show graphically how the gyre varies seasonally both in location and intensity. In the winter the gyre's axis is displaced northward; in summer it is shifted approximately 100 n mi southward and intensifies by approximately 0.2 kt. Throughout the year the current speed ranges from 0.1 to 0.3 kt. Hansen (1973) hypothesized that this intensification occasionally results in the advection of large volumes of cold subarctic water into the area surrounding OWS NOVEMBER, specifically during the period of 1968 through 1970.

3. The Subarctic Transition Zone

The North Pacific Gyre, linked with the Subarctic Gyre, is also involved in the formation of a transition zone which bounds the subtropical region to the north. Roden (1970) first described the subarctic transition zone as an area characterized by large temperature and salinity gradients and low hydrostatic stabilities. This oceanic front defines the boundary between the subtropical and subarctic water masses, varying in location between 30° and 40° North. Subsequent articles by Roden (1971, 1972) further described this zone in the North Pacific.

The transition zone is largely maintained by the surface wind stress. Its northern boundary is characterized by the subarctic water mass, resulting in temperature

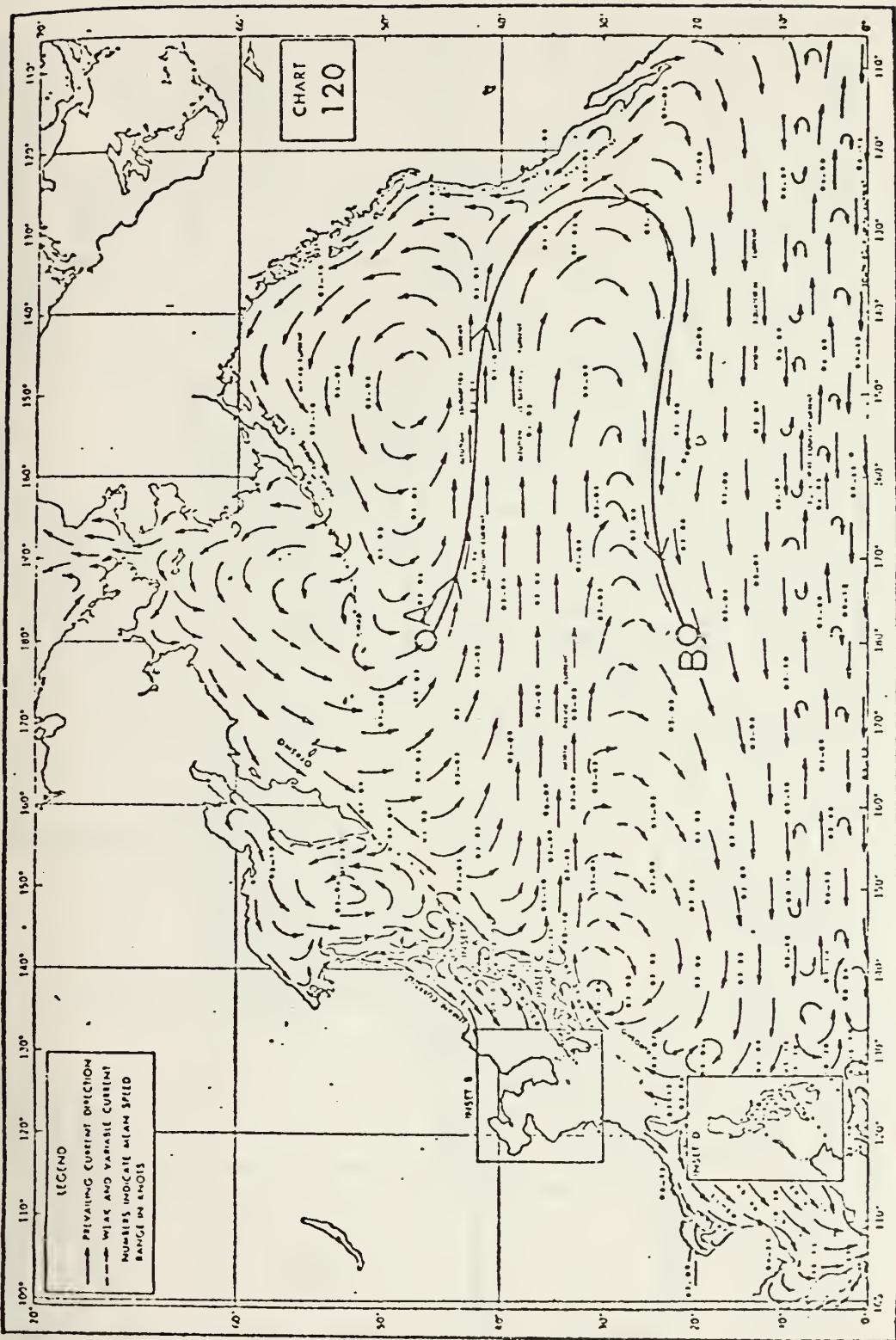


Figure 2. Winter surface circulation patterns for the North Pacific Ocean (Department of Commerce, 1961).

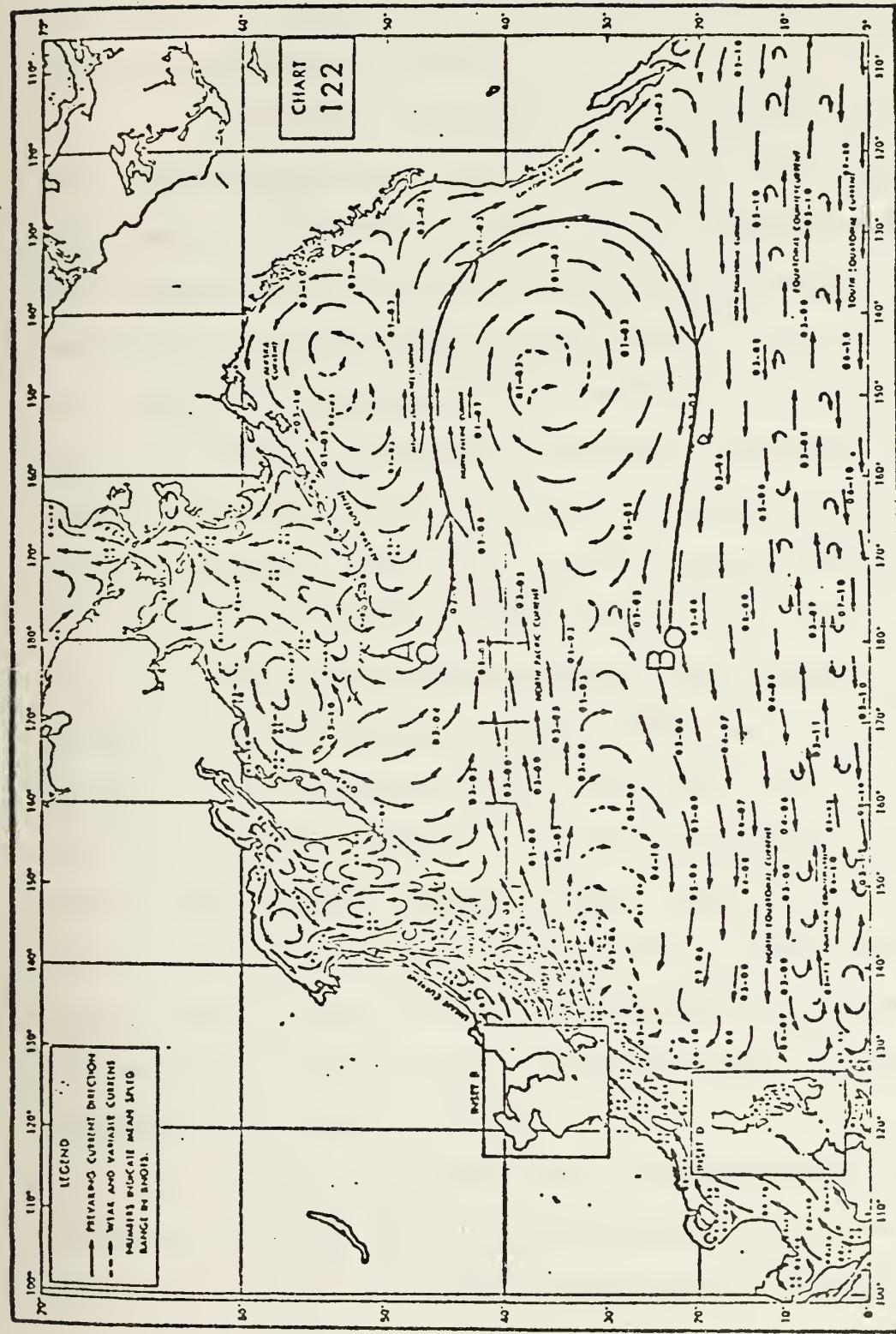


Figure 3. Summer surface circulation patterns for the North Pacific Ocean (Department of Commerce, 1961).

inversions, gravitational instabilities, a complicated sound velocity structure, and a disappearance of the subarctic halocline. It appears to lie in the area of 40N to 45N latitude throughout the year.

The southern boundary lies in the area of 30N to 32N, depending upon the season, and is displaced to the north more in the winter months. It is characterized by sharp temperature gradients in the thermocline and vertical instabilities in the mixed, upper layers. A salinity minimum is a common occurrence at this boundary. Also, around 30N there is a transition toward the higher salinities and higher temperatures of the subtropical water mass with salinities greater than 35°/oo and a strong halocline located at a depth between 50 m and 100 m.

Thus, the transition zone is defined by a deep salinity minimum, at approximately 500 m and a concurrent sound velocity minimum of less than 1480 m sec^{-1} . There is apparently both a seasonal and yearly variation in the position of the boundary (Seckel, 1968). Both the salinity minimum and sound velocity minimum rise as one goes further to the north. Roden (1970) also examined both barotropic and baroclinic motion in the transition zone and found that they rarely exceeded 5 cm sec^{-1} .

As the transition zone approaches North America, it turns to the southeast and becomes associated with the heat and salt flux divergence zones noted by Roden (1971)

in the area of OWS NOVEMBER. These are related to flux divergencies which are linked to the presence of horizontal advection of salinity as has been noted by Hansen (1973).

4. Average Annual Heat Balance

Based upon an examination of the annual heat balance of the North Pacific Ocean and the prevalent circulation patterns, Wyrtki (1965) determined that there is a small amount of heat gain in the vicinity of OWS NOVEMBER. On the average, he estimated this gain to be approximately 25 ly day^{-1} (where a langley, ly, = 1 cal cm^{-2}).

Seckel (1970) computed the heat balance for the area of OWS NOVEMBER during the Trade Wind Zone Pilot Study. His studies showed that the average net radiation computed on a daily basis was approximately 3 ly day^{-1} .

In a more recent study, Dorman (1974) studied 20 years of the heat balance at OWS NOVEMBER and separated it into two decades, 1951 through 1960 and 1961 through 1970. He found the two decades had distinctly separable heat balances, with the first decade having a deficit of 90 ly day^{-1} and the second a 40 ly day^{-1} deficit. These results were based on three-hourly meteorological observations. Dorman ascribed this difference to the change in climatic regimes described by Namias (1970, 1971), which resulted in displacement of cooler waters in the North Central Pacific by warmer water, beginning approximately in 1961.

5. Surface and Subsurface Temperature Anomalies

The sea-surface temperature and its anomalies have been widely studied in their relationships to both oceanographic and meteorological forecasting. Emphasis has also been placed on subsurface temperature anomalies. Beland (1971) illustrated their considerable importance to the flow of energy between the sea and the atmosphere. Bernstein (1974) used subsurface temperature measurements in the trade wind region northeast of Hawaii to identify mesoscale eddies drifting westward.

Namias (1971) in his study of an anomalous warming of the northeastern Pacific from May to June 1968, found the anomaly extended to a depth of 122 m. The warming, as a whole, was attributed to increased insolation and horizontal convergence of surface waters with its associated downwelling. The warming was identified with the development and maintenance of a Pacific anticyclone in June of 1968. The depth of the anomaly was attributed to bulk heating and downward advection of the warmer surface waters.

Clark (1972) concluded that both advection and non-advection processes are vital in determining the nature of surface temperature anomalies. He hypothesized that advection is the more important process in the winter and spring seasons, while the non-advection processes dominate in the summer and fall.

Beland (1971) in his study of the North Pacific, observed that about 50% of the near-surface temperature anomalies extended to at least 100 m depth. Most of these were negative anomalies in the months of March and April, apparently due to the distribution of heat throughout the upper layer by convective mixing. Positive anomalies were more common in December through February and were confined to the upper 30 m. Whereas the negative anomalies were of a magnitude of 1.0°C or more, the positive anomalies were generally less than 1.0°C.

Finally, Favorite and McClain (1973) reported that they could observe an orderly trans-Pacific movement of areas of anomalously warm and cold surface waters. These features were found to be quasi-permanent and apparently move around the North Pacific gyre, approximately 180 degrees out of phase and with periods of 5 to 6 years. These anomalies have been shown to have a great influence on the weather conditions off the west coasts of Canada and the United States, as shown by Namias and Born (1970).

From the above it can be noted that temperature anomalies, both surface and subsurface, and whether temporary or periodic, inhabit the region around OWS NOVEMBER, and in all probability, affect the location of the subarctic transition zone and the heat balance through the process of advection.

6. Recent Studies of Pacific Air-Sea Interactions

In recent years there have been three studies of the air-sea relationships at or around OWS NOVEMBER by different researchers using three different methods. Taken as a whole, they provide a logical progression, resulting in the methods utilized in this study. The first of these studies was carried out by Seckel and Yong in 1970, using harmonic analysis to study temperature and salinity fluctuations at Koko Head, Hawaii, and Christmas Island. Hansen (1973) used regression methods to study surface and subsurface temperatures and salinities at OWS NOVEMBER. Dorman (1974) used a spectral approach on time series of meteorological values at OWS NOVEMBER, dividing the time series into two ten-year decades. Since each of these studies utilized a different approach, each will be described in some detail here.

Seckel and Yong (1970) utilized temperatures and salinities measured at two points in the North Central Pacific; Koko Head, Hawaii, beginning in 1956 and running through 1969, and Christmas Island, beginning in 1954 and continuing through 1969. Their hypothesis was that fluctuations in the temperature and salinity time series reflected air-sea interaction processes (heat exchange, E-P) and oceanographic processes (advection, diffusion, etc.). Their approach used a Fourier series to provide a least squares fit to the observed values, which was done for each year of the available data. This approach expressed the

temperature and salinity as a function of time. Phase angles and coefficients for each year of salinity and temperature data from both locations were computed and tabulated.

The results at Koko Head showed seasonal variability in the salinity, but not the temperature series. At both locations the annual cycle ($n=1$) was the dominant cycle. Also of note were the 60- and 52-day cycles, ascribed by the authors to the presence of large geostrophic eddies. The Christmas Island results were quite similar except that they lacked evidence of the eddy-induced seasonal variations.

The next major method was a regression analysis used by Hansen (1973) in his study of temperature and salinity fluctuations at the surface, 50 m, and 200 m at OWS NOVEMBER. His data were for 1968 through 1970 and consisted of daily hydrographic records. Deterministic approaches showed that the temperature system was seasonal so the data were divided into four seasonal records, February to April, etc., in three month segments corresponding to Spring, Summer, etc. Regressions were then computed with the surface salinity, and 50 m and 200 m temperatures being regressed against the surface temperatures.

Most of the results showed little correlation, except occasionally during the spring and fall transition periods. Only a very poor correlation was found between the surface salinity and the surface temperature. From this, it

was concluded that these seasonal correlations were due to the advection of water by the shifting of the subarctic transition zone.

The final approach was that utilized by Dorman (1974) who computed spectra for surface wind speed, air pressure, sea temperature, air temperature and dew-point temperature for the 10-year period 1961 through 1970. Spectra were also computed for each of the four seasons. From this it was found that the energy peaked significantly at the annual, diurnal and semi-diurnal periods. Variances were largest in the winter season and smallest in the summer season, except for the sea-surface temperature which had its largest variance in the spring and the smallest in the winter. The maximum spectral density, for all variables, was at the annual period. From the annual period the spectral densities followed a decrease to the shorter periods, where significant spikes at the diurnal and semi-diurnal periods were noted in many of the spectra.

II. TREATMENT OF THE DATA

A. SOURCES

The climatological data used in this study were taken from the NWRC, Asheville data file of ocean weather station observations. These observations were begun at OWS NOVEMBER in 1947 and have continued to the present. The observations were taken every three hours, except for some short periods when all or only some of the observations were not taken for some parameters. The data extracted included the sea-surface temperature, air temperature, wind speed, surface pressure, cloud amount and the present weather codes for 1947 through 1970.

The salinity data used in this study came from the NODC hydrocast data. Since July 1966, the U.S. Coast Guard has conducted a continuing program of collecting Nansen cast data at OWS NOVEMBER. The surface salinity data for 1968 through 1970 was utilized in one segment of this study.

B. TREATMENT OF THE DATA TIME SERIES

Daily surface salinity values from 1968 though 1970 were extracted from the NODC data set and any gaps in the series were linearly interpolated. The daily salinity values were then included with surface pressure, evaporation, solar radiation, heat of evaporation, and wind speed, all taken directly from or computed from the Asheville data set. To this was added the calculated values of precipitation and

the evaporation minus precipitation terms to form a three-year time series which was then used in the regression analysis segment of this study.

The Asheville meteorological data were used primarily to create the 24-year (1947-1970) time series used in the spectral analysis segment of this study. The values of SST, air temperature, cloud amount, air-sea temperature difference, the saturated vapor pressure of the air, the vapor pressure over the water, wind speed, surface pressure and wind speed squared were extracted or computed for every three hours and then converted into daily average values and these were the values utilized in all the spectral analyses. For the cloud amount parameter, only daylight observations were used. From these data sets the heat budget terms were computed on a daily basis. For the occasional gaps in the data, linear interpolations were made.

C. THE HEAT BUDGET EQUATIONS

Empirically derived formulas were used to create the heat flux data for this study. The heat budget equation utilized excludes the effects of heat gained or lost by ocean currents:

$$\begin{aligned} \text{HEAT CHANGE} &= \text{HEAT INCOMING} - \text{HEAT LOSSES} \\ QN &= QS - QE - QB - QC \end{aligned} \tag{1}$$

where QN is the measured heat change, i.e., the heat gained or lost by the water due to the exchange processes. On the right-hand side of the equation, QS is the incident solar radiation, QB is the effective back radiation, QE is the heat of evaporation and QC is the conduction of sensible heat. Each of these terms will be described in detail below as to their formulation.

The equation used for QS is an improved version of that described by Seckel (1970). The basic form is:

$$QS = Q_0 K \quad (2)$$

where

$$Q_0 = A_0 + A_1 \cos\phi + B_1 \sin\phi + A_2 \cos 2\phi + B_2 \sin 2\phi \quad (3)$$

From the Smithsonian Tables (List, 1951), utilizing a transmission coefficient of 0.7, the above terms can be computed as follows for latitude 30N:

$$A_0 = -32.65 + 674.76 \cos \lambda = 551.7 \quad (4)$$

$$A_1 = 19.88 + 397.26 \cos(\lambda+90) = -178.7 \quad (5)$$

$$B_1 = -6.75 + 224.38 \sin \lambda = 105.4 \quad (6)$$

$$A_2 = -1.32 + 16.10 \sin 2(\lambda-45) = -9.4 \quad (7)$$

$$B_2 = -1.04 + 29.76 \cos 2(\lambda-5) = 18.1 \quad (8)$$

where:

λ is the latitude

t is the time of the year in days

C is the cloudiness in tenths, ranging from 0 to 1, and

$$\phi = \frac{2\pi}{365} (t-21) \quad (9)$$

The cloud factor correction utilized was one for cumulus type clouds (Seckel, personal communication):

$$K = 1.0 - 0.012C - 0.363C^2 \quad (10)$$

An albedo factor of 0.95 was utilized.

The new equation results in values approximately one-fourth to one-third greater than previously estimated.

The effective back radiation, QB, is a Berlyand expression utilized by Seckel (1970):

$$QB = 1.14 (10)^{-7} (273.16 + T_w)^4 \quad (11)$$

$$(0.39 - 0.05 \sqrt{e_a}) (1 - 0.6C^2)$$

$$+ 4.58 (10)^{-7} (273.16 + T_w)^3$$

$$(T_w - T_a)$$

where T_w is the sea-surface temperature in degrees C. It utilizes a cloud coefficient of 0.6 and neglects latitudinal variation.

QE, the heat of evaporation depends upon the difference between the saturation vapor pressure over the sea at the sea-surface temperature and the vapor pressure of the air above the sea surface:

$$QE = 3767 C_D (0.98 e_w - e_a) W \quad (12)$$

where:

e_a is the vapor pressure of the air (mb)

e_w is the saturation vapor pressure over water
at sea-surface temperature (mb)

C_D is the drag coefficient

w is the wind speed ($m sec^{-1}$)

Whereas Seckel utilized a varying drag coefficient, this study utilized a constant drag coefficient of 1.3×10^{-3} . As in Seckel's study, the vapor pressure difference was multiplied by 0.98 to account for the salt effect.

The final heat balance term is the conduction of sensible heat, QC:

$$QC = 4 (T_w - T_a) w \quad (13)$$

where:

T_w is the sea-surface temperature (C)

T_a is the air temperature (C)

w is the wind speed ($m sec^{-1}$)

It can be seen readily that QC is dependent upon two factors, the air-sea temperature difference and the wind speed. It is usually the smallest of all the heat exchange terms. As with Seckel's study, a conduction factor of 4.0 is used.

The above equations were those used to create the five heat flux time series used in this study. QS was computed using only the daylight cloudiness observations, night

values being ignored. A daily value for each day was calculated for the span of 1947 through 1970. They were compared with Seckel's (1970) results to guarantee their accuracy, where possible. The utilization of the newer QS formula resulted in higher values for this term than he reported and correspondingly more positive values of QN.

D. COMPUTATION OF PRECIPITATION

In relation to the study of the surface salinity at OWS NOVEMBER, it was vital to have some measurement of E-P, the evaporation-precipitation term. The evaporation was readily available, having been calculated from the data already available. However, no direct measurement of precipitation is made routinely at ocean weather stations. To gain a measure of the precipitation, Tucker's (1961) method was used. Elliot and Reed (1973) applied this method along the Pacific Northwest Coast with good results. This approach is based on the current weather code which gives a numerical value to the type of weather being observed at three hour intervals; this yields an estimation of the amount of rainfall. The amount of rainfall was based on the amount of precipitation observed for a similar weather code at a land station where actual rainfall measurements were available. Based on a frequency analysis for all the codes and the measured rainfall, Tucker was able to fit a numerical system of equations to the prediction of precipitation, using only the weather code as an input.

In Tucker's system all codes less than 50 are rejected since they contribute no appreciable amount of precipitation. The codes used had values ranging from 50 to 99. Each of these codes is associated with a term made up of some combination of three rainfall parameters. These are x, associated with light, continuous rainfall; y, associated with moderate, continuous rainfall; and z, which is a measure of heavy, continuous rain. These three parameters satisfy the following criteria:

$$y = nx \text{ mm where } 1 < n < 5$$

$$z = my \text{ mm where } 1 < m < 3$$

$$z = 10 \text{ mm (in 3 hr)}$$

For this study x was set equal to 1.85, $y = 5.66$, and $z = 8.13$, the same values as derived by Tucker from his frequency distributions of the codes recorded and the actual rainfall at land stations. Thus, the amount of rainfall, in millimeters, associated with each current weather code can be determined from one of the above, or a combination or ratio of these terms. For example, weather code 64, which is "rain, not freezing, intermittent, but heavy at the time of observation", is described by the term $z/2$. Likewise, code 84, "showers of rain and snow mixed, moderate to heavy", is represented by the formula, $\frac{y+z}{4}$.

For the purposes of this study, the precipitation was calculated every three hours and then summed to obtain the total precipitation for that particular day. This was then subtracted from the evaporation term for that day to yield the E-P time series for 1968 through 1970.

E. COMPUTATION OF THE ENERGY SPECTRUM

The covariance function, $R_{xx}(\tau)$, of a time series of random data is a measure of the general dependence of the values of the data at one time, t , with the values at another time, $t+\tau$. The covariance equation in integral form is:

$$R_{xx}(\tau) = \lim_{T \rightarrow \infty} \int_0^T x(t) x(t+\tau) dt \quad (14)$$

This yields the average of the product of the value of a data point, $x(t)$, with that of $x(t+\tau)$, a point at some lagged time increment, τ . As T approaches infinity, the resulting average product approaches the exact autocovariance function.

Since this data forms a finite data set, a lag window is applied to the covariance function. The Parzen window used in this analysis is defined as:

$$P(\tau) = 1 - 6(\tau/T_m)^2 + 6(\tau/T_m)^3 \quad \tau = 0, 1, 2, \dots, T_m/2 \quad (15)$$

$$= (1 - (\tau/T_m))^3 \quad \tau = T_m/2 + 1, \dots, T_m \quad (16)$$

$$= 0 \quad \tau = T_m \quad (17)$$

where τ is the lag and T_m is the total lag time. This window has the desirable effect of minimizing any numerical instability due to side lobe resonance (Bendat and Piersol, 1971).

The energy-density spectrum is obtained by Fourier transforming the covariance function. The form of the energy-density spectrum, corrected by the Parzen window is:

$$S_{xx}(f) = \int_{-\infty}^{\infty} P(\tau) R_{xx}(\tau) e^{-i2\pi f\tau} d\tau \quad (18)$$

with f being a specific frequency. The energy-density spectrum is the statistical estimate of the energy-density distribution for the frequency bands forming the total spectrum.

F. COMPUTATION OF THE CROSS-SPECTRAL DENSITY

When comparing two-time series of the same length and time interval, the cross-covariance function similar to the covariance function, is used. For a given lag time interval, τ , it has the form:

$$R_{xy}(\tau) = \lim_{T \rightarrow \infty} \frac{1}{T} \int_0^T x(t)y(t+\tau) dt. \quad (19)$$

Because the data sets are records of finite length, the Parzen lag window is applied. Thus, similar to the covariance function and the spectral density, the cross-spectral density is the Fourier transform of the Parzen window lagged, cross-covariance function.

The cross-spectral density is a non-zero function, having real and imaginary parts of the form:

$$S_{xy}(f) = \int_{-\infty}^{\infty} P(\tau) S_{xy}(\tau) e^{-i2\pi f\tau} d\tau \quad (20)$$

which can be expressed in terms of its real and imaginary parts:

$$S_{xy}(f) = C_{xy}(f) - iQ_{xy}(f) \quad (21)$$

where $C_{xy}(f)$, the real co-spectrum, is given by:

$$C_{xy}(f) = 2 \int_0^{\infty} [S_{xy}(\tau) + S_{xy}(-\tau)] \cos(2\pi f\tau) d\tau \quad (22)$$

and the imaginary quadrature spectrum, $Q_{xy}(f)$:

$$Q_{xy}(f) = 2 \int_0^{\infty} [S_{xy}(\tau) - S_{xy}(-\tau)] \sin(2\pi f\tau) d\tau \quad (23)$$

G. COMPUTATION OF THE COHERENCE FUNCTION

The coherence function, for two records x and y , is a measure of the statistical dependence of two series. The coherence is defined as:

$$\gamma_{xy}^2(f) = \frac{[S_{xy}(f)]^2}{S_{xx}(f) S_{yy}(f)} \quad (24)$$

where:

$$0 \leq \gamma_{xy}^2(f) \leq 1.$$

If $\gamma_{xy}^2(f) = 1$, for all frequencies, then the two records are totally coherent and statistically dependent. If $\gamma_{xy}^2(f) = 0$ for all frequencies, then the x and y records are totally independent. For intermediate range coherences three conclusions may be inferred concerning the two records (Bendat and Piersol, 1966); (1) the systems are non-linear, (2) noise is present in the measurements, or (3) the output response is due to more than one input function. It should be noted that this is the coherence squared, but will hereafter be referred to as the coherence.

H. COMPUTATION OF THE PHASE ANGLE

The average angular difference by which the cross-correlated points in $y(t)$ lead their corresponding components of $x(t)$ is called the cross-spectral phase angle, $\epsilon_{xy}(f)$. It is computed from the co- and quad-spectra by:

$$\epsilon_{xy}(f) = \arctan [Q_{xy}(f) / C_{xy}(f)] \quad (25)$$

Since this study deals with a time step of one day, the phase angles can be converted to a lag in days from:

$$L(\text{days}) = \frac{\epsilon_{xy}(f) * 0.01745}{f * 2\pi} \quad (26)$$

When either the value of the phase angle (measured in degrees) or the lag time (in days) is positive, then $y(t)$ lags behind $x(t)$ by that amount; if the value is negative, then $x(t)$ lags behind $y(t)$.

In the computations involved in this study there were 8766 data points. A standard lag of 876, 10% of the total time series, was used. The 90% confidence interval for the energy densities was computed by the formula:

$$r = \frac{N}{M} 3.7 \quad (27)$$

when $N = 8716$, $M = 876$.

The limits were found from tables to be:

$$a = 24.075$$

$$b = 52.192$$

where a and b are the values which when subtracted and added, respectively, to the value of the spectral energy density result in the lower and upper 90% confidence limits.

III. RESULTS

A. REGRESSION ANALYSIS

The regression analysis phase of this study was limited to the three-year time series of the NODC hydrocast data due to the desire to utilize the salinity data present in this data set alone. The period covered was from 1968 through 1970, the same time span of data analyzed by Hansen (1973). At this point it is appropriate to visually inspect the surface temperature and salinity data sets in a deterministic approach.

Figure 4 shows the monthly means of the sea-surface temperature for 1968 through 1970 compared to Robinson's (1971) twenty-year long-term monthly means. A characteristic annual cycle is evident in all three years as well as the long-term mean. The only major difference from the long-term mean is a persistent cold anomaly which is present in the spring of all three years.

The year 1968 was the most anomalous, being both colder in March and April and warmer for the balance of the year. Namias (1971) attributed this to bulk heating, advection and mixing. The year 1969 was more normal, except for an anomalous cooling in the summer, attributed by Roden (1970) to a southward displacement of the transition zone. This has been supported by Wickett and Thompson (1971a, 1971b) whose meridional transport data, Figure 5, shows a transport

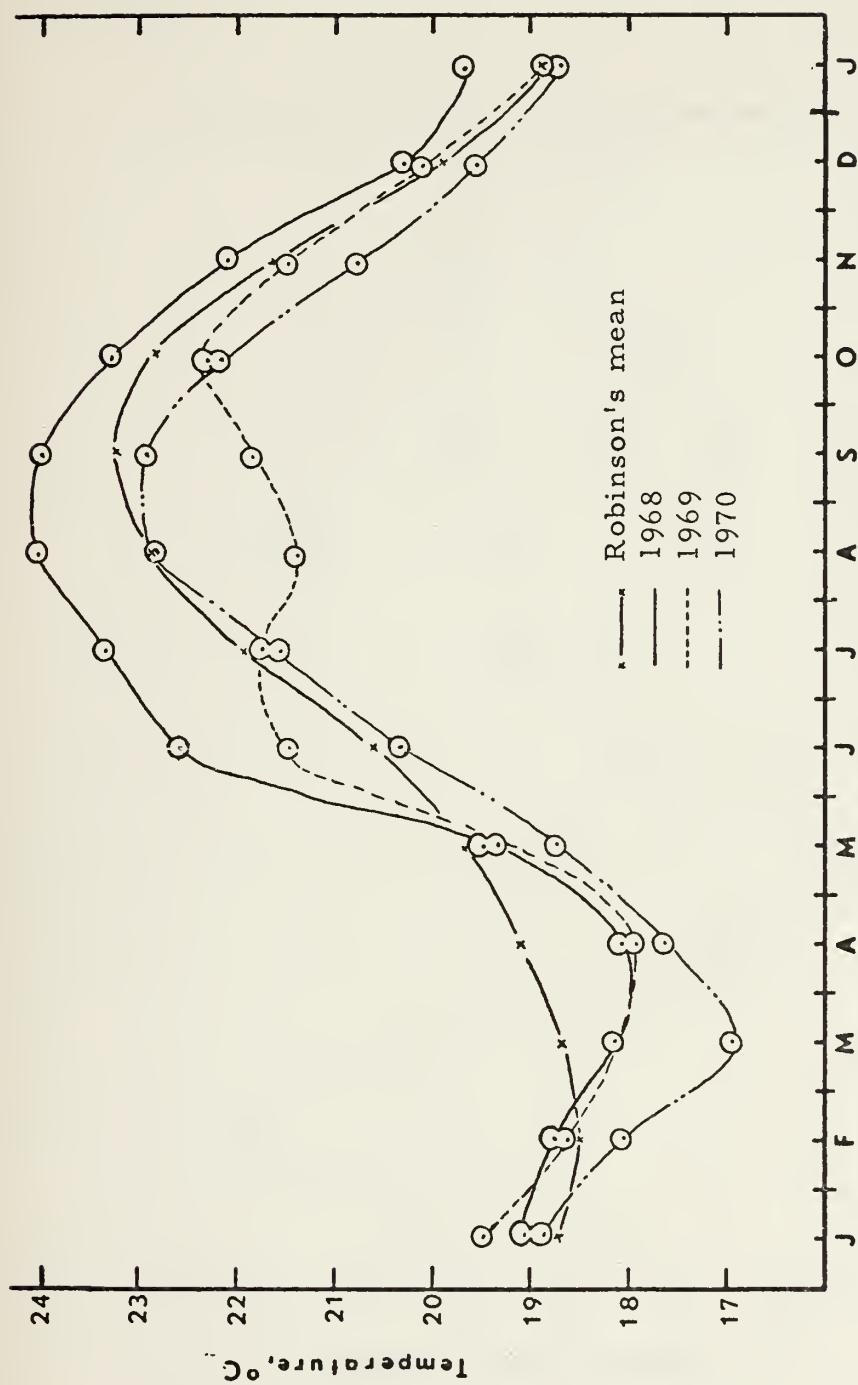


Figure 4. Monthly means of the Observed Surface Temperatures, 1968 through 1970 and Robinson's (1971) Long-term Monthly Means.

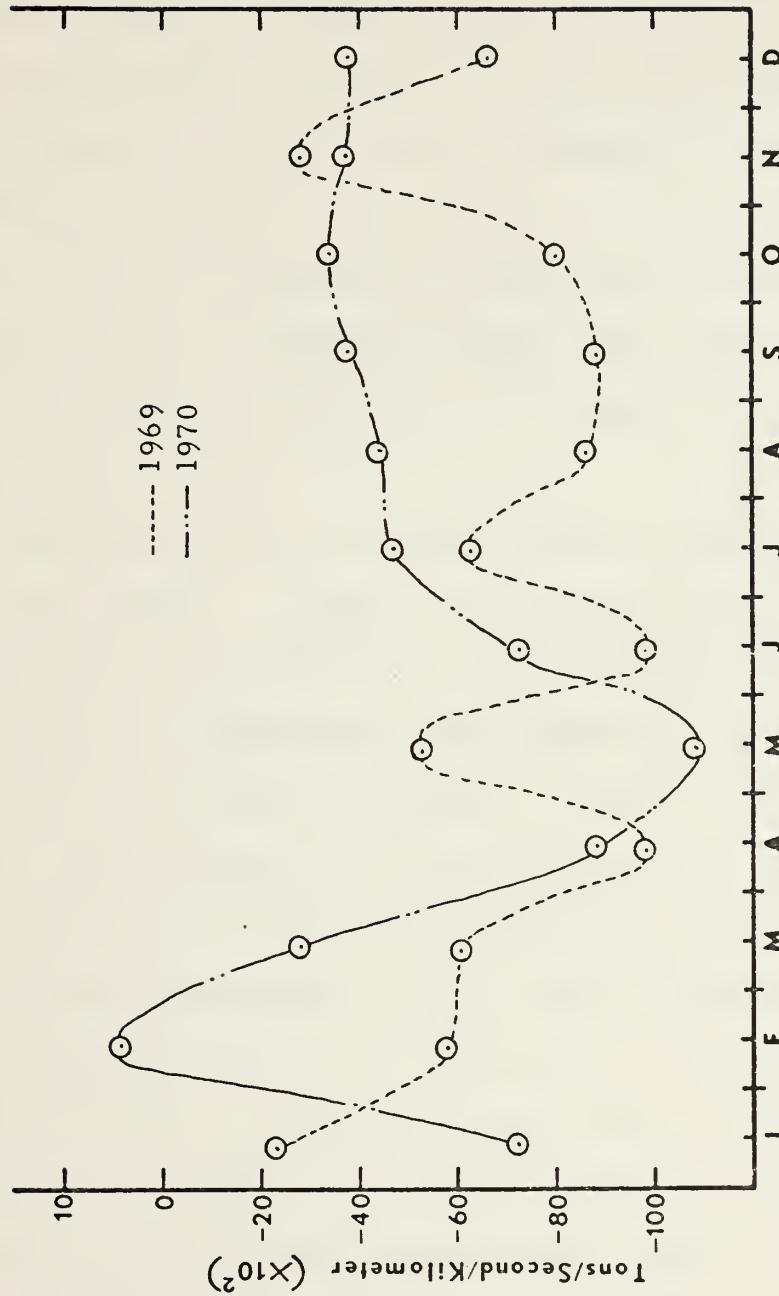


Figure 5. Total monthly Meridional Transport Between 135° and 145° West at 30° North during 1969 and 1970.
Positive Northward, Negative Southward (after Wickett and Thomson, 1971a, 1971b).

of cooler water equatorward during April through October, 1969. Except for the spring months, 1970 was more typical to the long-term mean than the preceding two years.

Figure 6 shows the monthly means of surface salinity for 1968 through 1970. The most obvious fact from this figure is that the salinity shows evidence of having a semi-annual periodicity. In addition to this, there is the feature in 1970 of a reduction in the salinity, possibly due to the advection of abnormally fresh, cool subarctic water, indicating an equatorward shift in the subarctic transition zone.

As stated previously, Hansen (1973) did much work on the same data series. Figure 7 is the result of his bivariate analysis of the monthly mean surface temperatures and salinities. It shows that the temperature follows four phases, namely the seasonal phases of summer, autumn, winter and spring, with autumn and spring appearing as transitional phases. Hansen also performed a regression test of the seasonal surface temperature versus the surface salinity and also the temperature at 50 m. The correlation coefficients are given in Table I. For most cases the correlations were poor, especially between the surface salinity and the surface temperature. The only good correlations occurred during the transition periods when both the surface temperature and salinity were responding to the changes of the annual cycle.

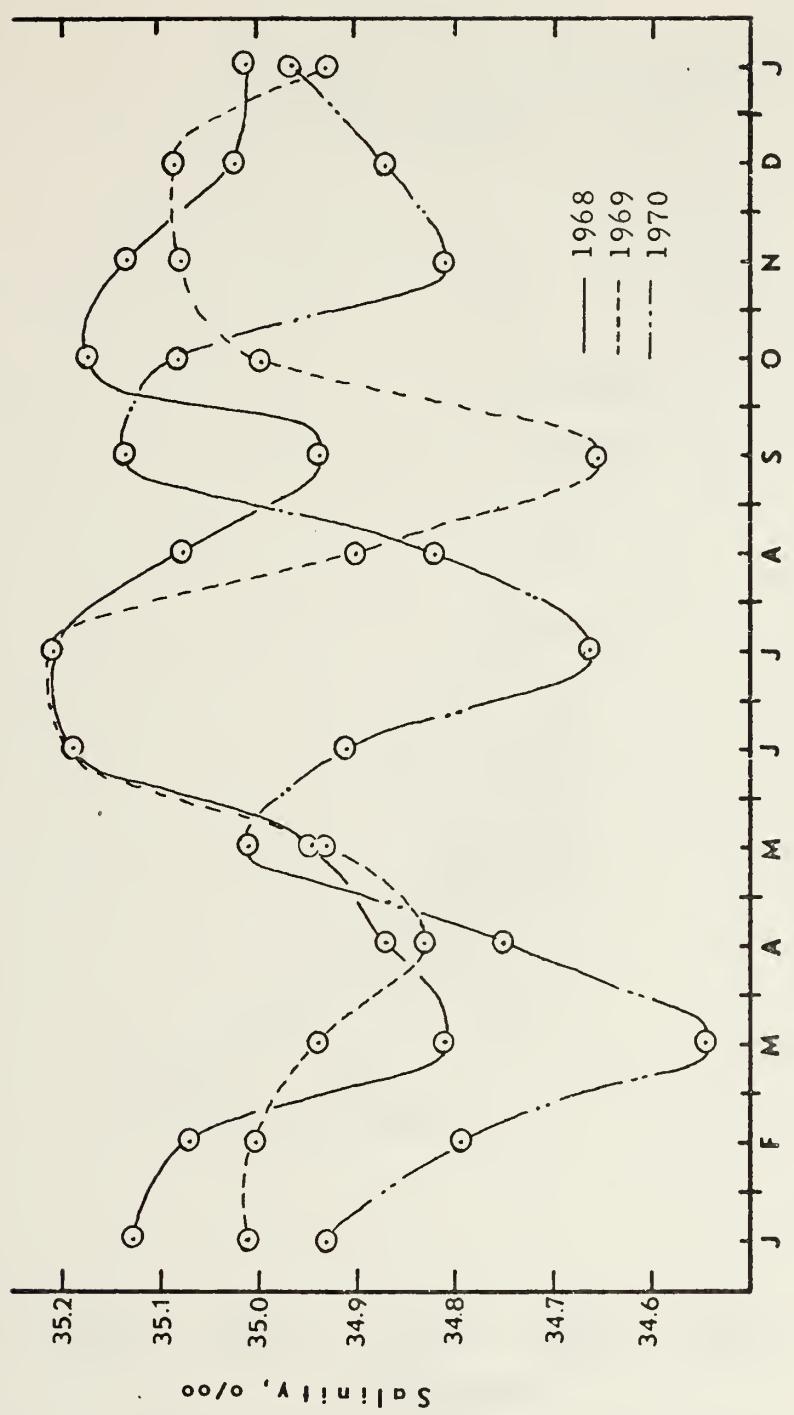


Figure 6. Monthly Means of Observed Surface Salinity,
1968 through 1970 (from Hansen, 1973).

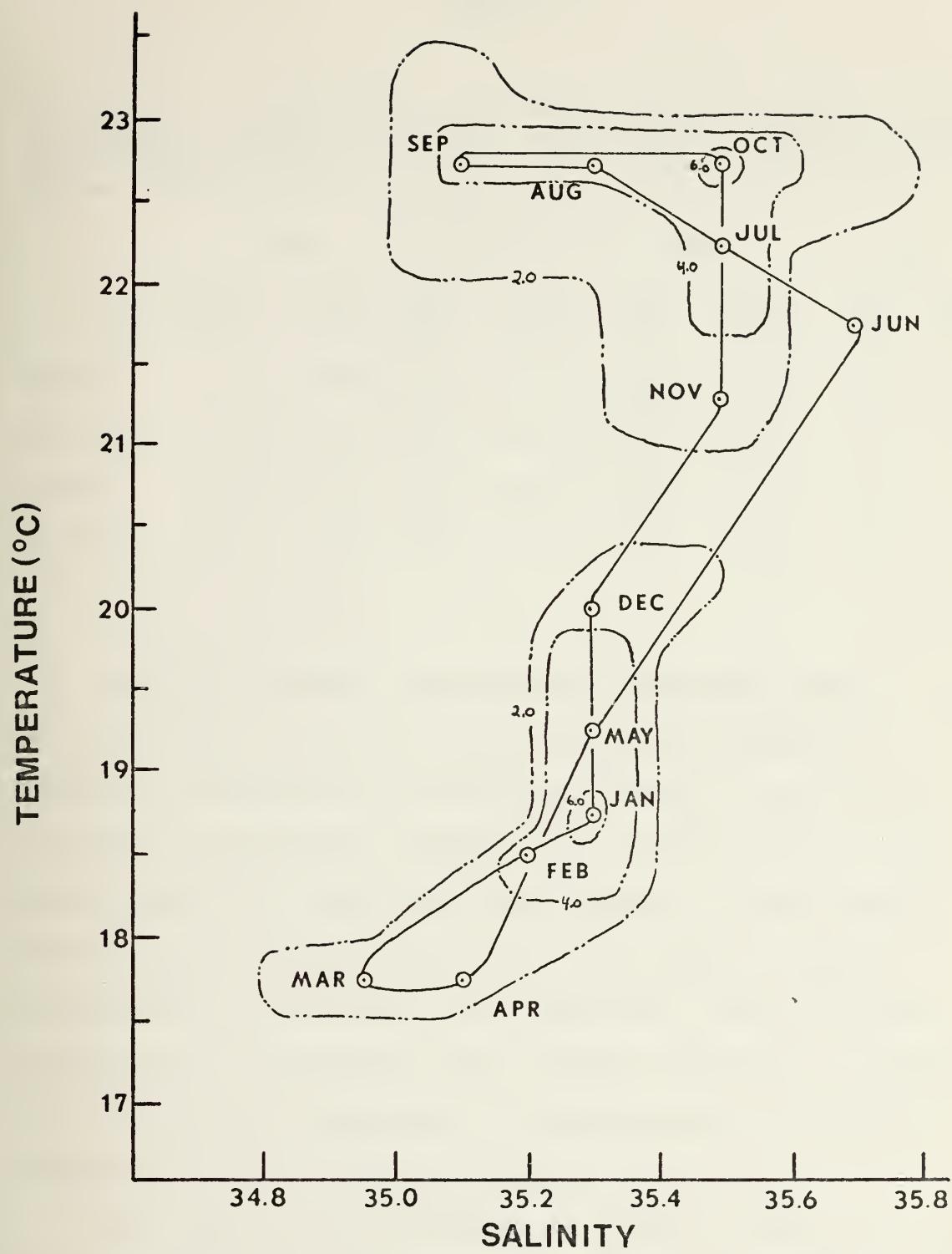


Figure 7. Monthly Temperature and Salinity Modal Points Resulting from the Bivariate Analysis of the Data, 1968 through 1970 (from Hansen, 1973).

Table I

Correlation coefficients, by season, resulting from a linear regression analysis of surface temperature, surface salinity, and 50 meter temperature data, 1968 through 1970 (from Hansen, 1973).

	1968		1969		1970		Mean	
	Sfc Sal	50 m Temp	Sfc Sal	50 m Temp	Sfc Sal	50 m Temp	Sfc Sal	50m Temp
Winter	0.61	0.58	0.61	0.56	0.77	0.42	0.66	0.52
Warming	0.59	0.90	0.77	0.70	-0.11	0.67	0.42	0.75
Summer	-0.28	0.31	0.16	0.40	-0.05	-0.29	-0.05	0.14
Cooling	0.57	0.88	0.84	0.98	-0.01	0.96	0.47	0.94

Since the surface salinity and temperature were obviously responding to different causal parameters, a further regression analysis was performed. The surface salinity was regressed against many of the classical parameters associated with its fluctuations. In this instance, there were seven variables chosen; the vapor pressure difference, E-P, the heat of evaporation, the wind speed, evaporation, precipitation and surface pressure. No separation into seasons was made, all computations being made as a three-year series of daily average observations.

The results of the regression study are shown in Table II. Two major observations are apparent from this study. First, the standard parameters often associated with a change in salinity, especially E-P, showed very little correlation with the surface salinity. However, this is quite reasonable

Table II

Correlation coefficients, all parameters vs surface salinity.

PARAMETER	1968	1969	1970
$e_w - e_a$.094	.043	.097
E-P	.045	.151	-.016
QE	.028	.161	-.001
W	-.030	.163	-.043
PS	.830	.803	.803
E	.034	.166	.004
P	-.030	-.019	.028

considering the location of OWS NOVEMBER with its year round excess of evaporation over precipitation. The other major observation is that only one parameter is consistently highly correlated, i.e., the surface pressure.

This latter fact could mean that the changes in the surface salinity are largely due to an advective process, as is supported by the fact that the other non-conservative terms associated with the air-sea interface have such poor correlation coefficients. This suggests that the changes in the surface salinity are related to some geostrophic process, most likely an advection of water from the subarctic region, corresponding to a southward shift of the transition zone.

B. SPECTRAL ANALYSIS

The major thrust of this study was to use spectral analysis as a means of numerically describing the interactions at the air-sea interface at OWS NOVEMBER. Specifically, it was to identify those meteorological parameters most closely associated with fluctuations in the sea-surface temperature. Since this latter term has a dominant annual cycle and most of the other parameters were compared with this term, the results presented here will primarily be those for the yearly cycle. The logic of this is further demonstrated by comparing the percentages of the variance explained by the yearly cycle for the different terms. The percentages shown in Table III were calculated from the energy-density functions which were produced for each time series. This function is a measure of the variance contained at that frequency. To compute the variance contributed by the yearly cycle a bandwidth was specified by those frequencies for which the value of the energy-density function decreased to less than half the value at the yearly cycle.

From Table III it can be seen that 86% of the total variance of the sea-surface temperature is explained by the yearly cycle alone. For this reason, even though the spectral results were computed for frequencies as short as the monthly cycle, the emphasis of this study was placed on the yearly cycle. Also, since the major areas of interest were the correlations with the heat budget terms, each of

these is discussed describing its influence on the sea-surface temperature. In addition, the contributing effects of the parameters involved in the computation of that heat flux term will be discussed.

Table III

Contribution of the yearly cycle to the total variance explained

<u>Parameter</u>	<u>% of Total Variance</u>
SST	86
T _a	80
T _w - T _a	20
e _w	86
e _a	46
e _w -e _a	14
E (evaporation)	34
W	8
W ²	9
C1 (cloud cover)	7
PS	8
QS	70
QB	87
QE	11
QC	17
QN	42

The first heat budget term to be examined is the solar radiation, QS, which like the surface temperature exhibits a strong annual periodicity. From Tabe III, it can be seen that 70% of the total variance of this term is explained by the annual cycle. The term which primarily controls the variation of QS is the cloud cover, C1, and in Table IV the spectral results for the yearly cycle are given for the interactions of the SST and C1 with QS. Along with a correlation of the cloud cover with the SST, the complete spectral results are shown in Figures 8, 9, and 10. Figures A-1, A-2, and A-3 of Appendix A show the respective auto-correlation functions and Figures B-1 and B-2 of Appendix B illustrate 15-year segments of the 24-year time series of the SST and QS.

As would be expected, Table IV shows that the solar radiation is strongly coherent with the SST at the yearly cycle. However, there is a lag of almost 60 days, which cannot be explained by the cloud cover correlation with its low coherence and therefore dubious lag time. Also, there is a strong coherence at the semi-annual cycle, which appears to reflect the effect of the cloud cover as seen in Figure 10.

Table IV
Spectral results for the solar radiation,
SST and cloud cover at the yearly cycle.

Parameter Codes	Coherence	Phase (degrees)	Lag (days)
QS vs SST	.94567	- 60.0	-58.5
QS vs C1	.85236	126.2	122.0

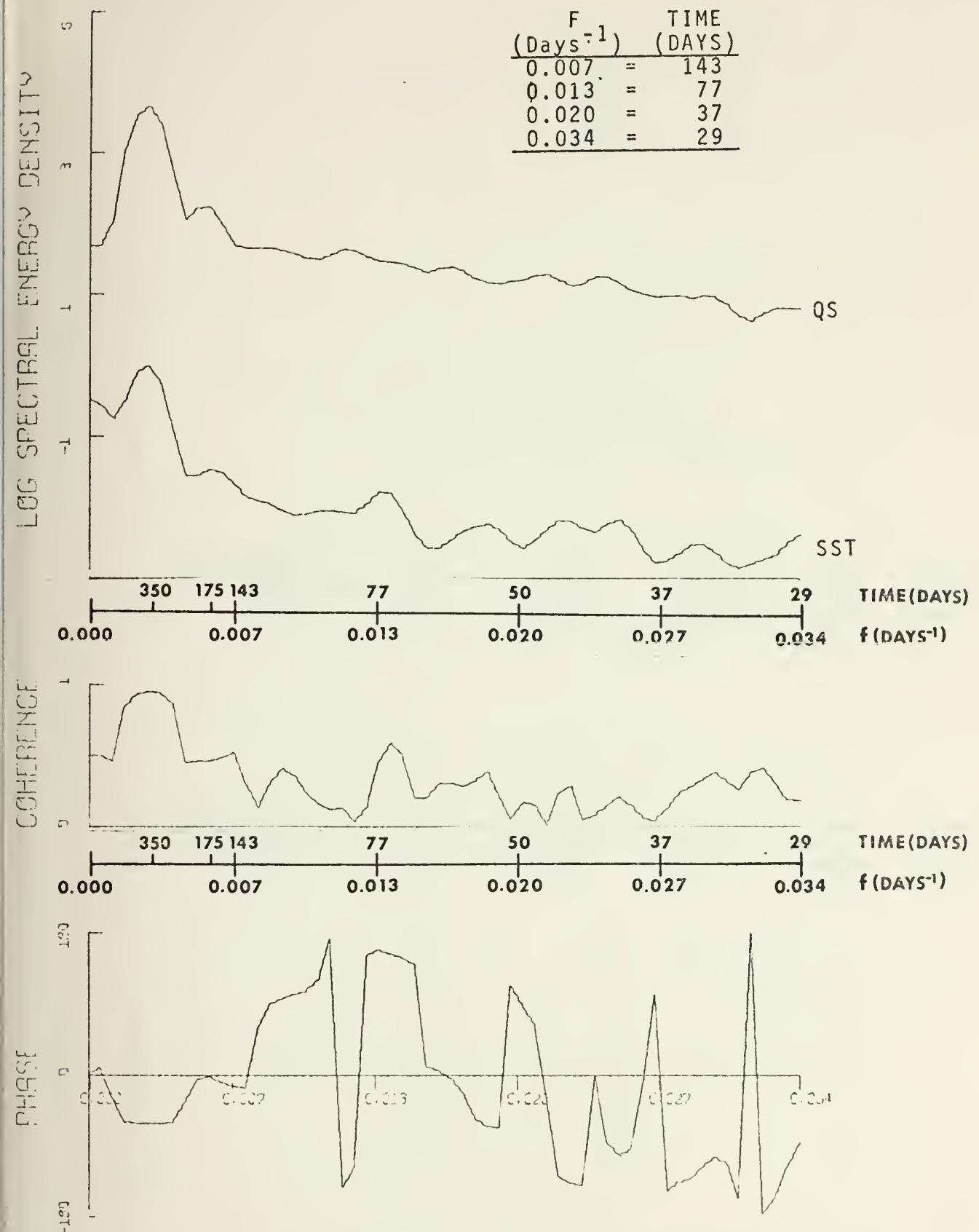


Figure 8. Spectral density, coherence and phase for the SST vs the solar radiation.

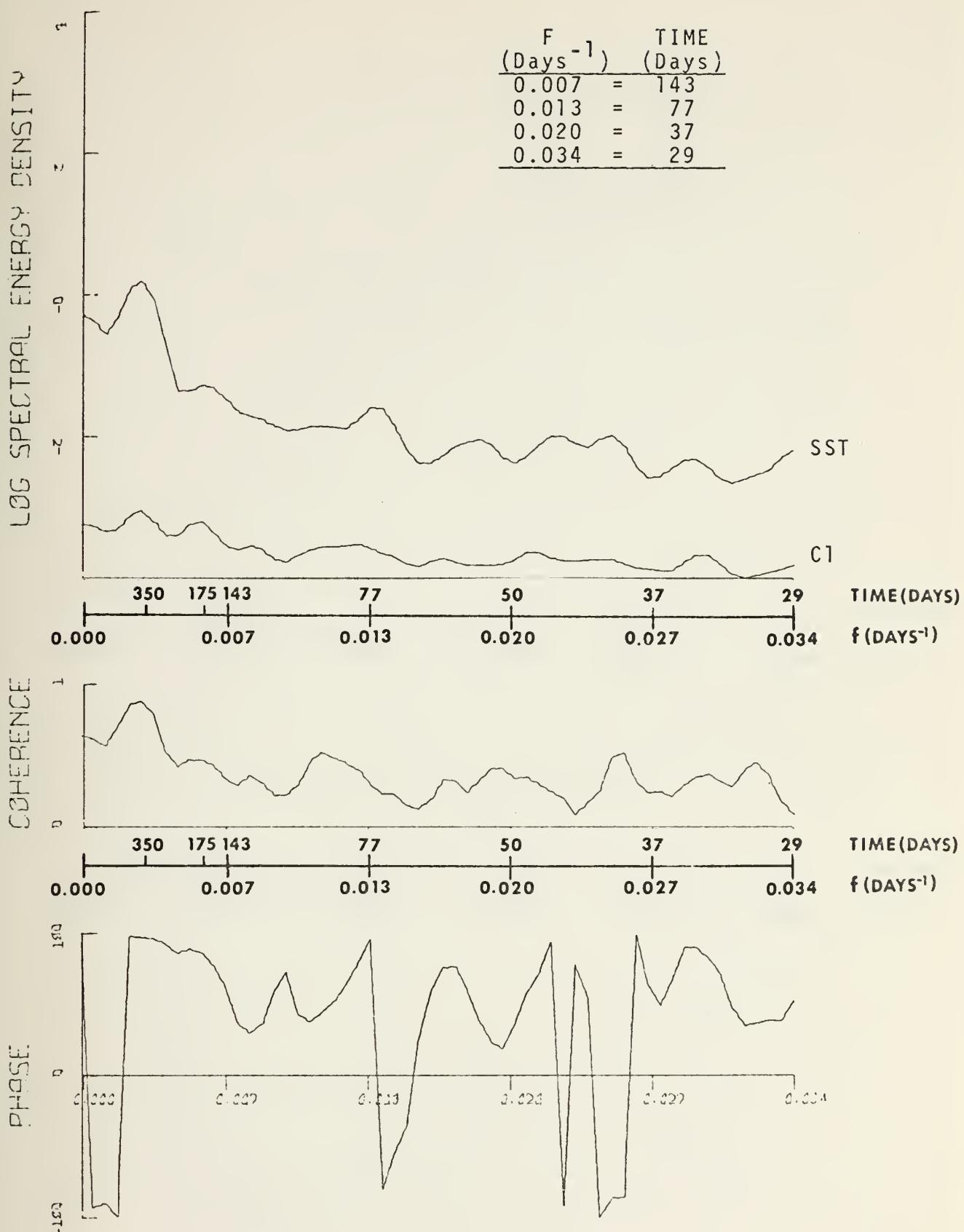


Figure 9. Spectral density, coherence and phase for the SST vs cloud cover.

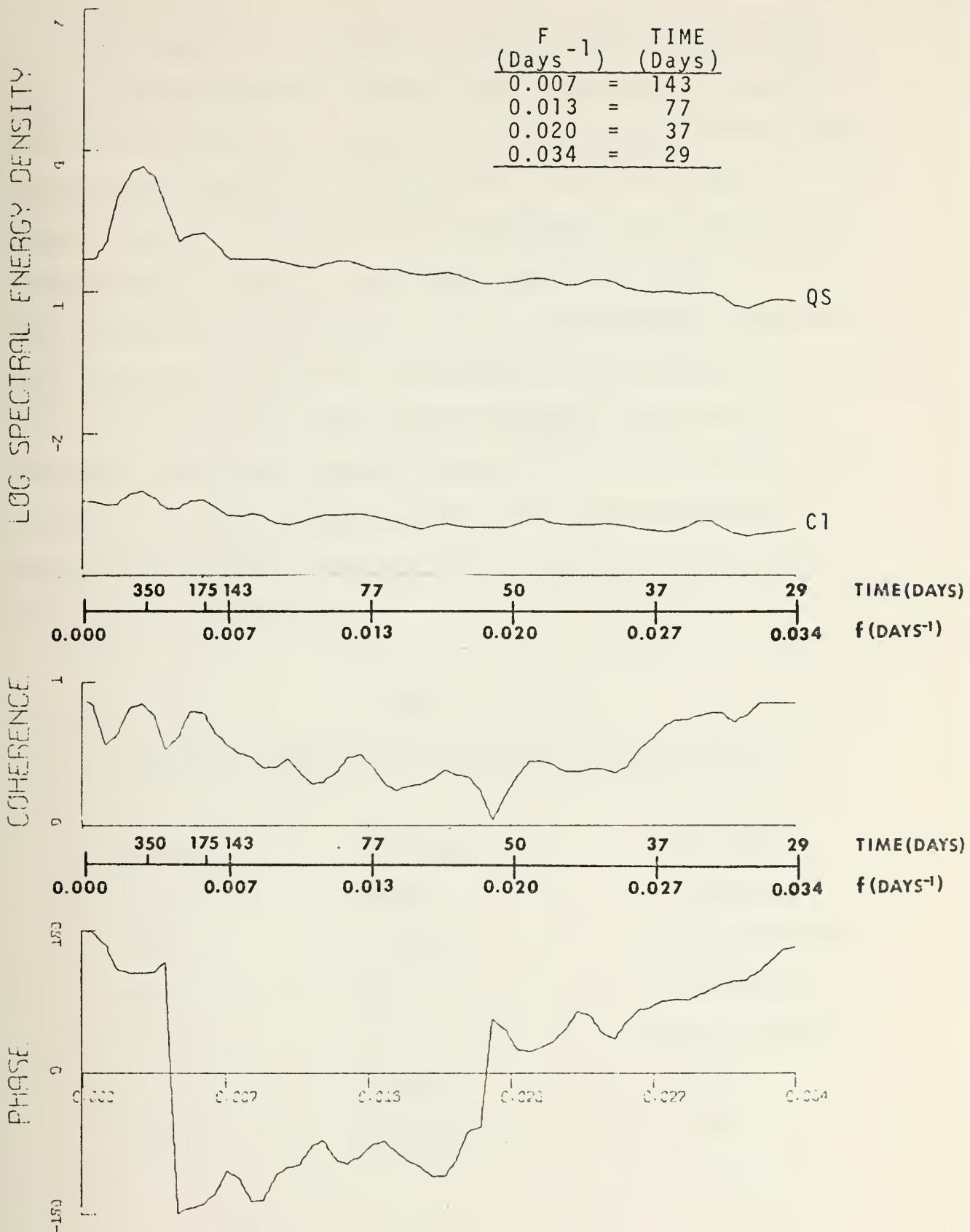


Figure 10. Spectral density, coherence and phase for the solar radiation vs cloud cover.

This rather large lag is equally evident in Table V which shows the months in which SST and the heat transfer terms achieve their maximum and minimum values. Along with Figures B-1 and B-2, it can be seen that the surface temperature is a maximum during September and at a minimum during March. However, the solar radiation peaks in August and is a minimum in January. This is approximately the 58 day lag shown in Table IV. This lag can be explained by the fact that the solar radiation is absorbed and transferred throughout the upper mixed layers of the water column. Thus, a lag of two months exists before the total effects of the solar radiation are represented by a warming or cooling of the surface temperature.

Table V

Months of maximum and minimum value for the surface temperature and the heat budget term.

<u>Parameter</u>	<u>Minimum</u>	<u>Maximum</u>
SST	March	September
QS	January	August/September
QE	May	November
QB	July/January	March/November
QC	August	March
QN	December	June

In Figures 8, 9, and 10 there are also noticeable peaks at the six month cycle and three month cycle. These results are thought to be, for the most part, artifacts of the yearly cycle and are relatively small when compared to the yearly value; yet, as mentioned before, several show evidence of being real.

The major result of this portion of the study is that the use of the new formula for QS results in significantly, higher estimates of this term than previously calculated. Fluctuations in solar radiation precede those of the surface temperature by approximately two months and thus the solar radiation is a good predictor of the surface temperature, as would be expected.

The back radiation, QB, represents the loss of heat energy as long-wave radiation. It is strongly dependent upon the surface temperature and the vapor pressure of the atmosphere. Table VI summarizes the results for the yearly cycle of QB and associated terms with the surface temperature.

Table VI

Spectral results for the back radiation and related parameters at the yearly cycle.

QB vs SST	.59755	121.8	119.0
T _a vs SST	.98800	- 6.0	- 5.9
T _w -T _a vs SST	.67194	-49.1	-47.9
e _a vs SST	.98815	- 6.4	- 6.2
QB vs C1	.37565	-82.9	-78.1

From Table III it was found that about 87% of the total variance of the back radiation was contained in the yearly cycle; yet, Table VI shows only a slight coherence between the SST and QB for the yearly cycle and a large dubious lag value. This can be explained by examining the interactions involved. The back radiation term decreases as the SST increases due to the effect of the humidity in the atmosphere directly above the surface. Thus, over the yearly cycle, QB would ideally be 180 degrees out of phase with the surface temperature. However, the back radiation is also decreased by an increase in the cloud coverage. Thus, the condition illustrated by Figure B-3 and Table V may be seen where QB is shown to have two maxima and minima on roughly a semi-annual basis. From Figure A-4 it may be seen that this feature is superimposed over a larger annual cycle which explains the high percentage of the variance contained by the yearly cycle.

The complete spectral results for the above correlations are given in Figures 11 through 15. From Figure 11 it is seen that QB has two peaks, one at the yearly cycle and another lesser one at one-half the yearly cycle. Of the related parameters, only the air temperature and the vapor pressure of the air are highly correlated and then at only the yearly cycle. The air-sea temperature difference is only mildly correlated at the yearly cycle and its associated lag is highly suspect for that reason. The air-sea temperature

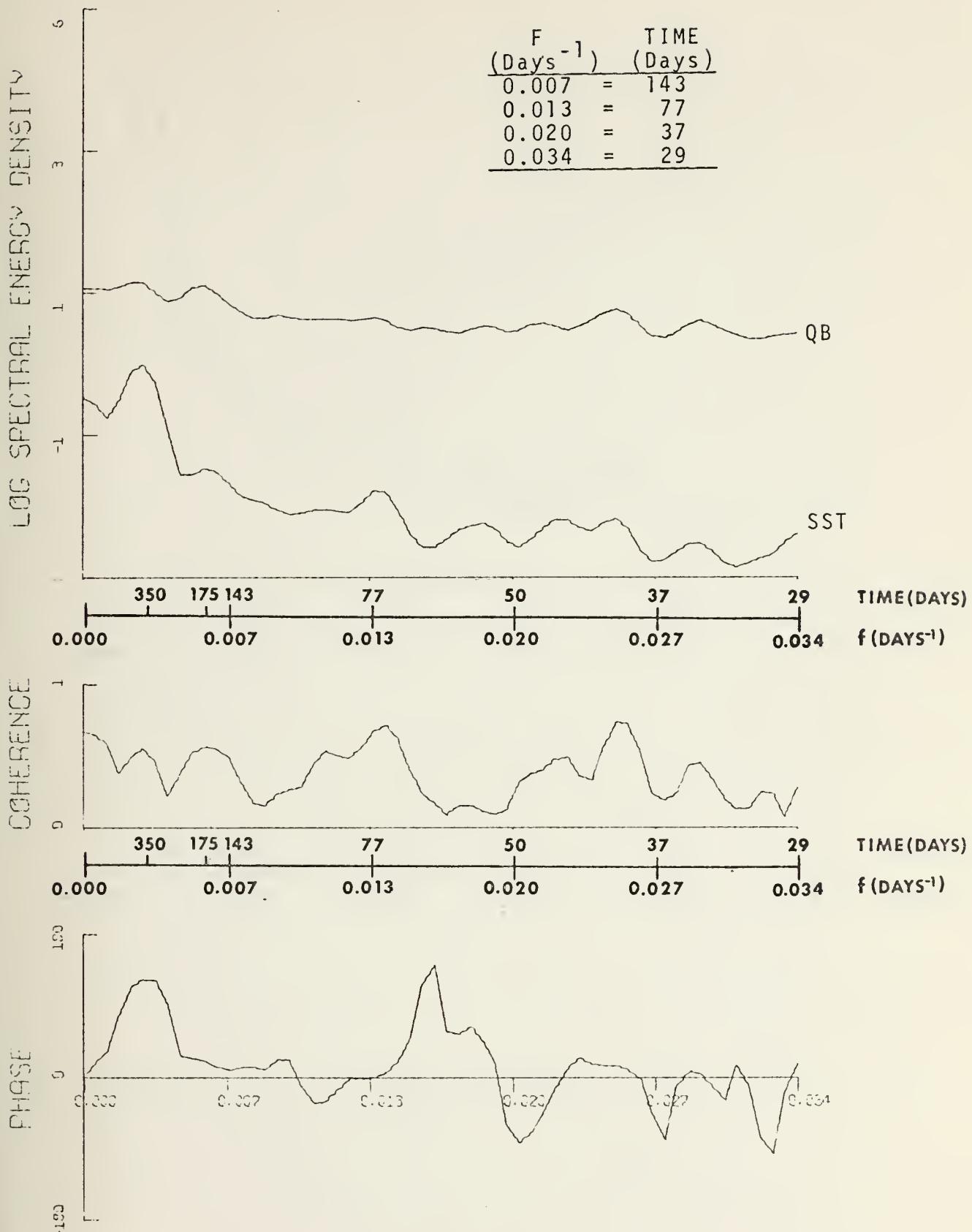


Figure 11. Spectral density, coherence and phase for the SST vs the back radiation.

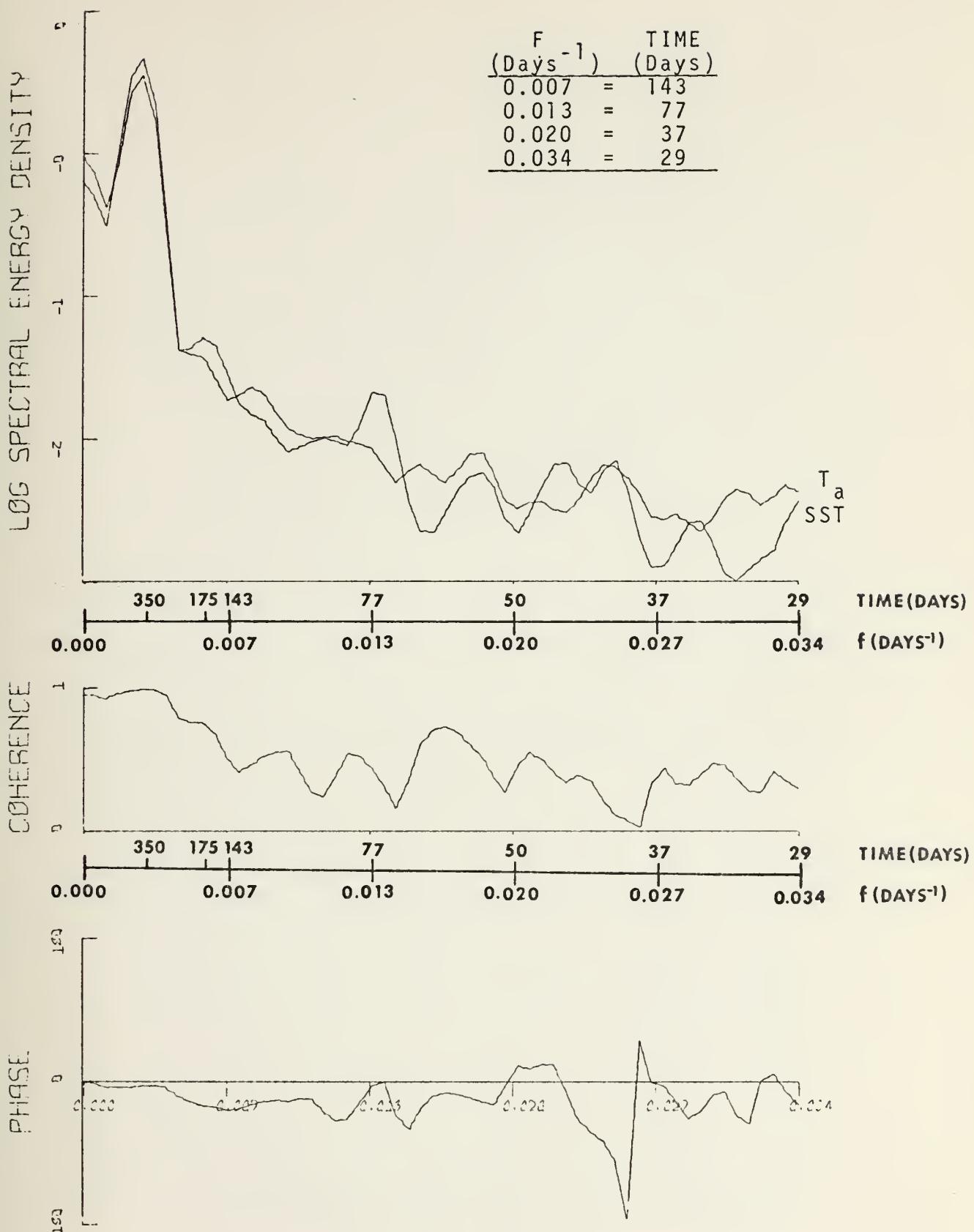


Figure 12. Spectral density, coherence and phase for the SST vs air temperature.

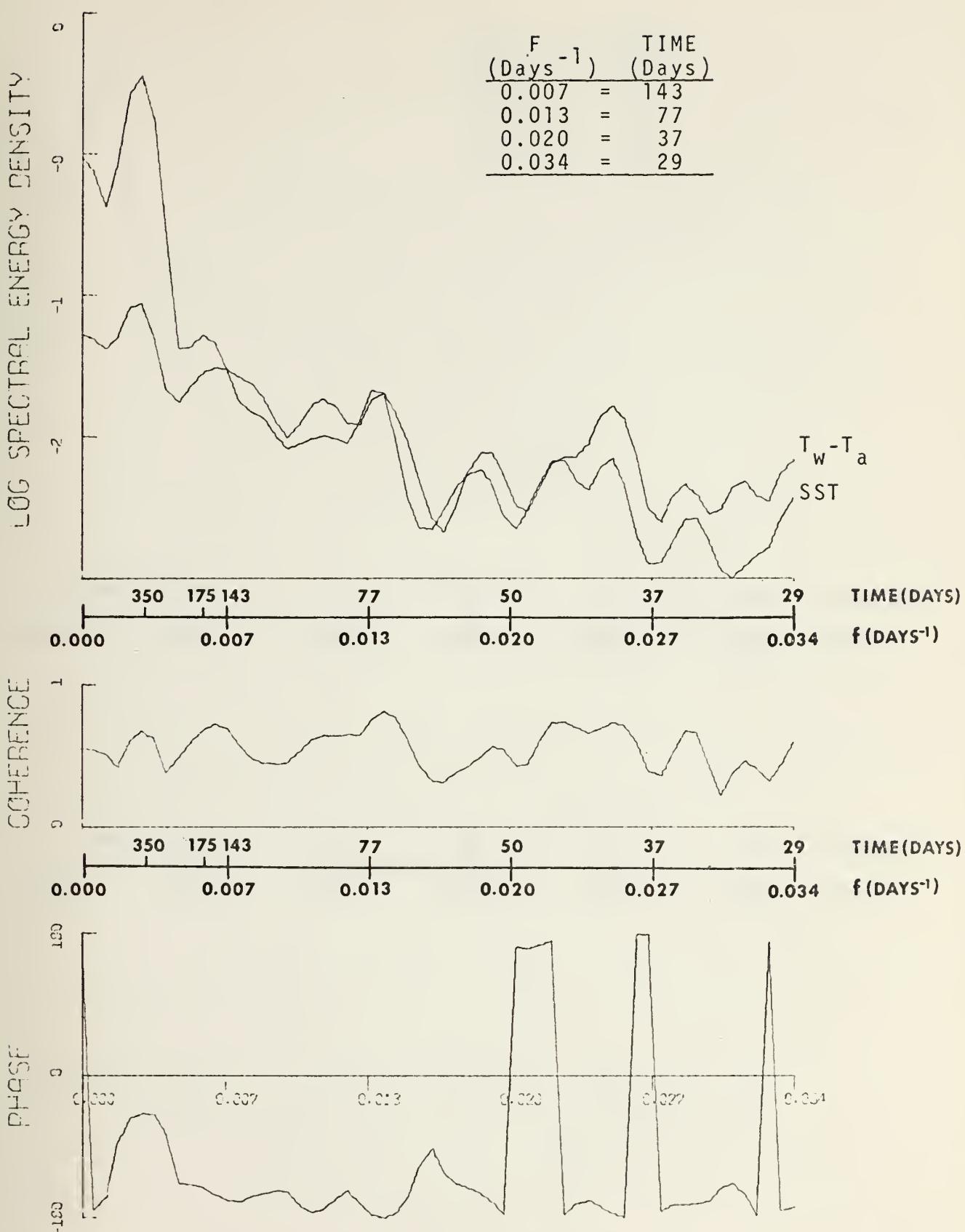


Figure 13. Spectral density, coherence and phase for the SST vs the air-sea temperature difference.

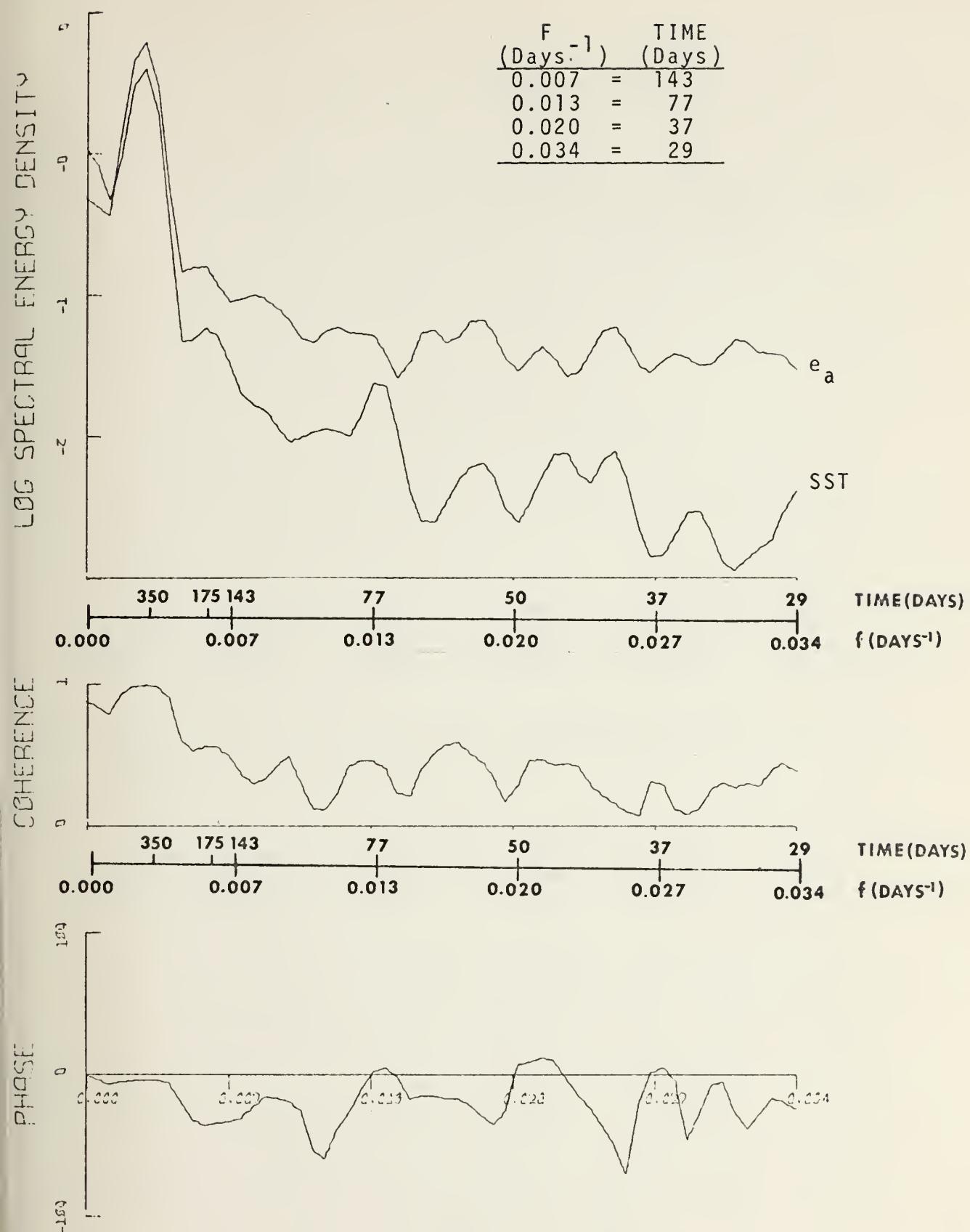


Figure 14. Spectral density, coherence and phase for the SST vs the vapor pressure of the air.

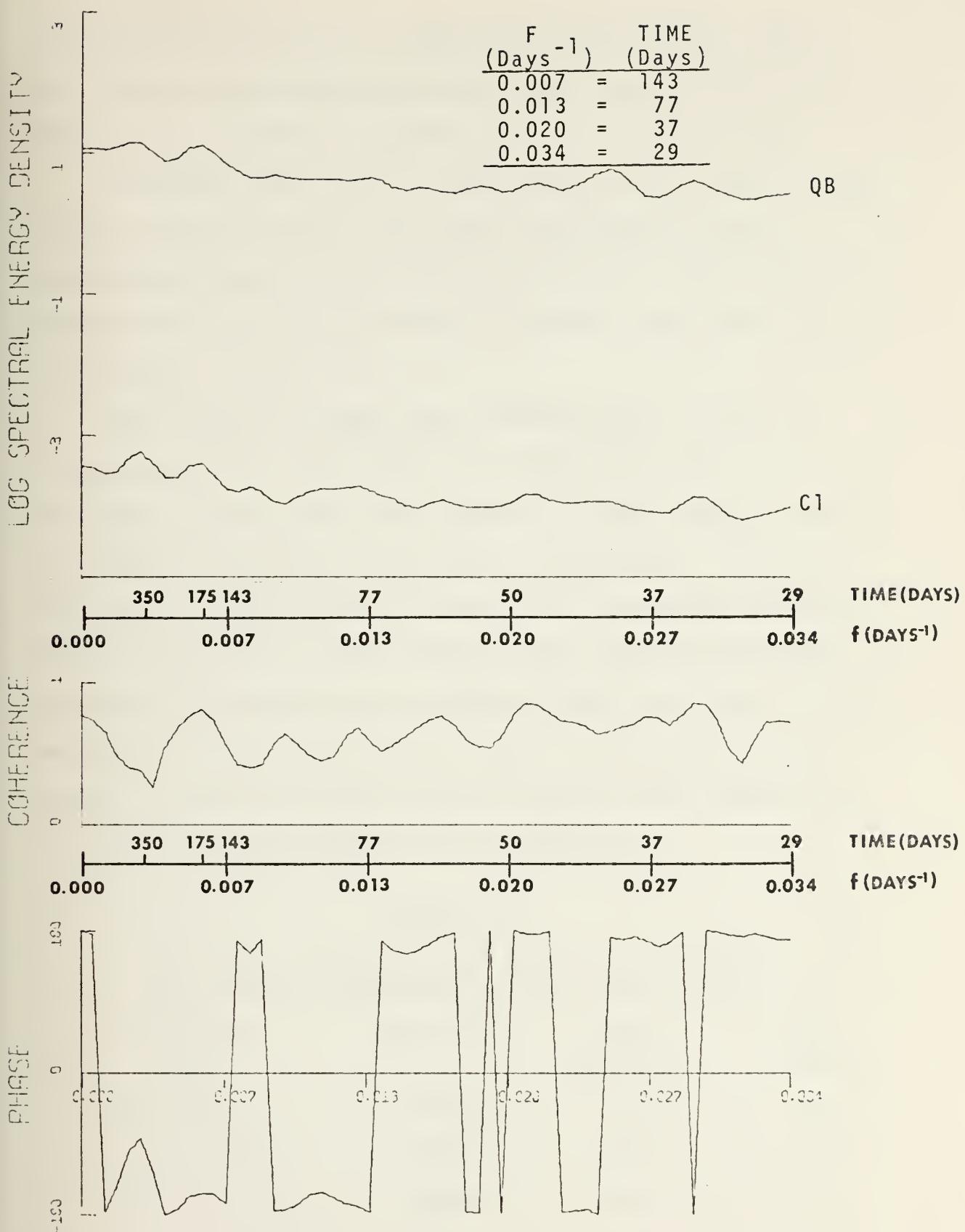


Figure 15. Spectral density, coherence and phase for the back radiation vs cloud cover.

difference demonstrates a coherence at the 70-day cycle. Also the back radiation and cloud cover show evidence of this possible seasonal coherence.

The major conclusion to be drawn from this part of the study is that the back radiation is semi-annual with an underlying yearly cycle. As a predictor for the surface temperature, the back radiation is poorly associated with fluctuations in SST.

The next heat budget term examined was the heat of evaporation, QE. Only 11% of the total variance of QE is explained by the yearly cycle which in itself means that QE is a poor predictor of the surface temperature. This is verified by Table VII which shows only a moderate coherence for QE and the SST at the yearly cycle. This lessens the credibility of the large lag shown; also from Table V it can be seen that the actual lag is more of the nature of two months. Additionally the lag is positive which means that fluctuations of QE follow those of the SST.

Table VII

Spectral results for the heat of evaporation and related parameters at the yearly cycle.

Parameter Codes	Coherence	Phase (degrees)	Lag (days)
QE vs SST	.84783	97.5	95.0
$e_w - e_a$ vs SST	.86425	22.8	22.3
W vs SST	.95593	138.1	135.4
W^2 vs SST	.95512	135.4	132.0
E vs SST	.81722	76.9	74.9

Whereas QE is only weakly correlated with the SST at the annual cycle, the wind speed and wind speed squared are highly correlated to the SST with large positive lags (approximately 4 months). These would seem to explain much of the lag in QE when they are coupled with the vapor pressure difference. QE is a transfer mechanism of the form:

$$QE = - A_e \frac{\partial f}{\partial z} \quad (28)$$

where:

A_e is the eddy coefficient for water vapor

$\frac{\partial f}{\partial z}$ is the humidity gradient above the sea surface.

Both of these are highly dependent upon the wind speed. Even though the vapor pressure difference and evaporation are weakly correlated to the SST, they combine with the wind speed to produce the low coherence and large lag for QE.

Figures 16 through 20 give the complete spectral results for the correlations shown. Being computed from the surface temperature, e_w and SST are naturally perfectly correlated and is not included. The yearly cycle dominates in all except for QE, the vapor pressure difference and the evaporation, as would be expected. Both the wind speed and wind speed squared show a definite coherence at about the 80-day cycle. The autocorrelation functions are found in Figures A-9 through A-13 and a 15-year plot of QE and its anomaly is shown in Figure B-4.

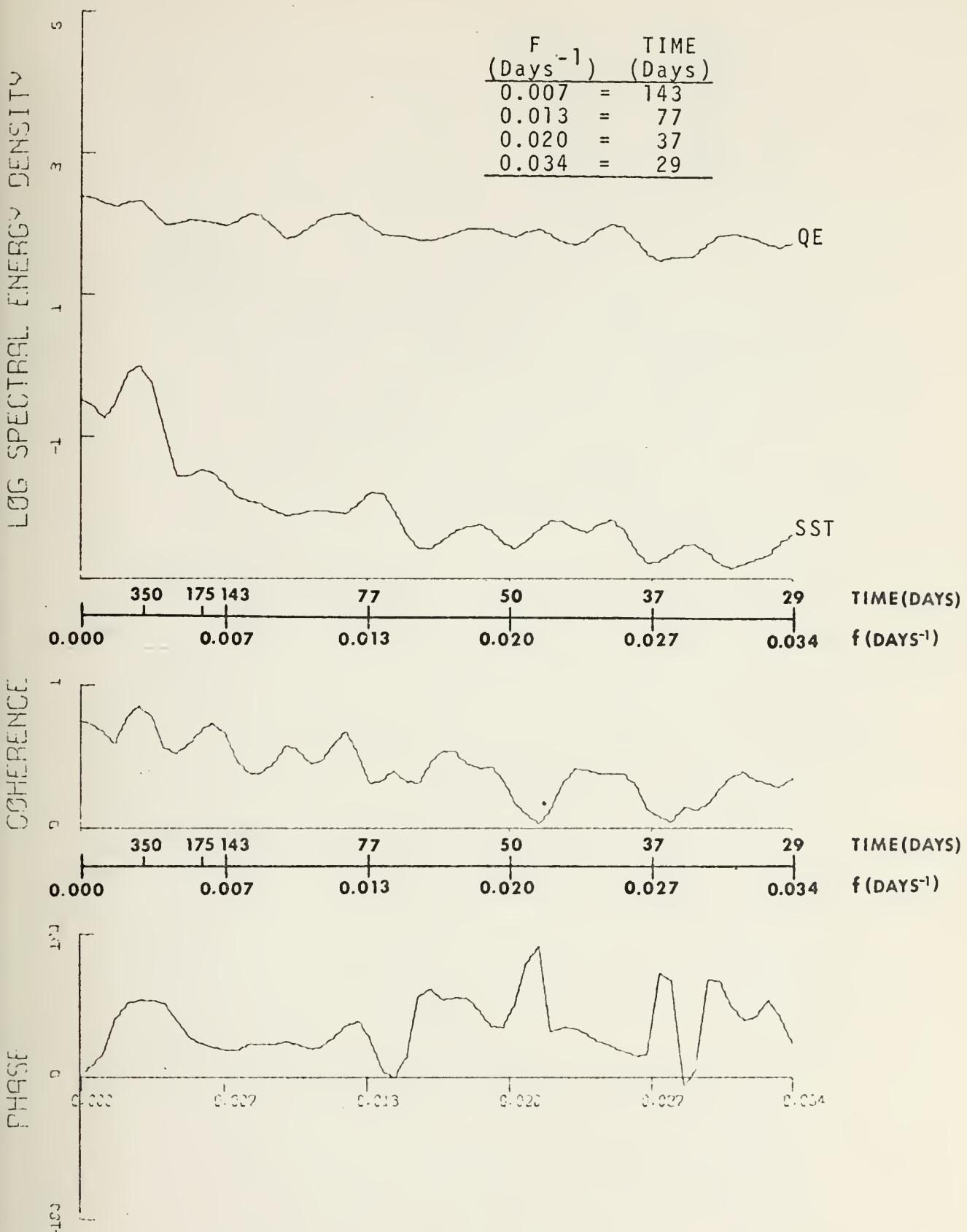


Figure 16. Spectral density, coherence and phase for the SST vs the heat of evaporation.

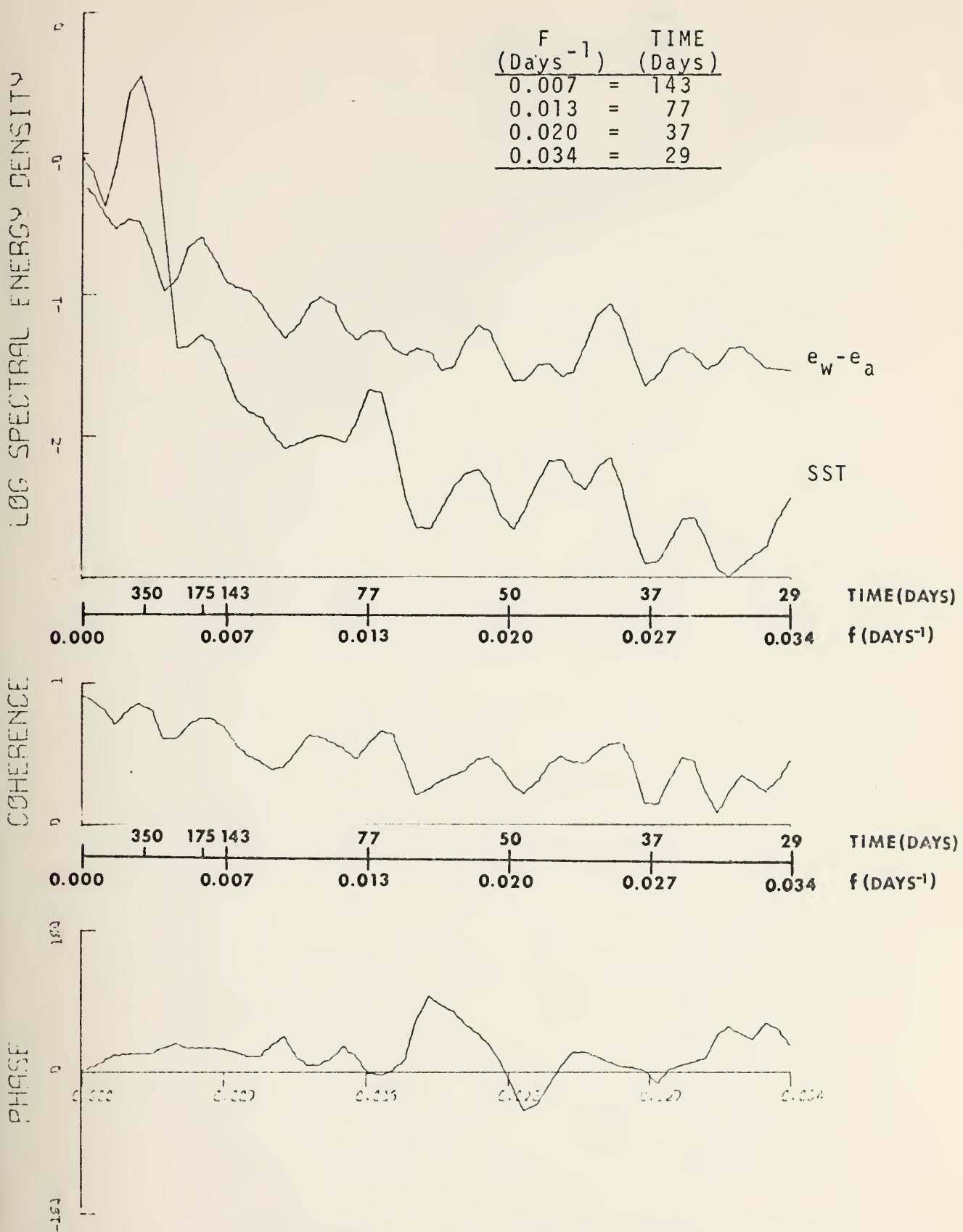


Figure 17. Spectral density, coherence and phase for the SST vs the vapor pressure difference.

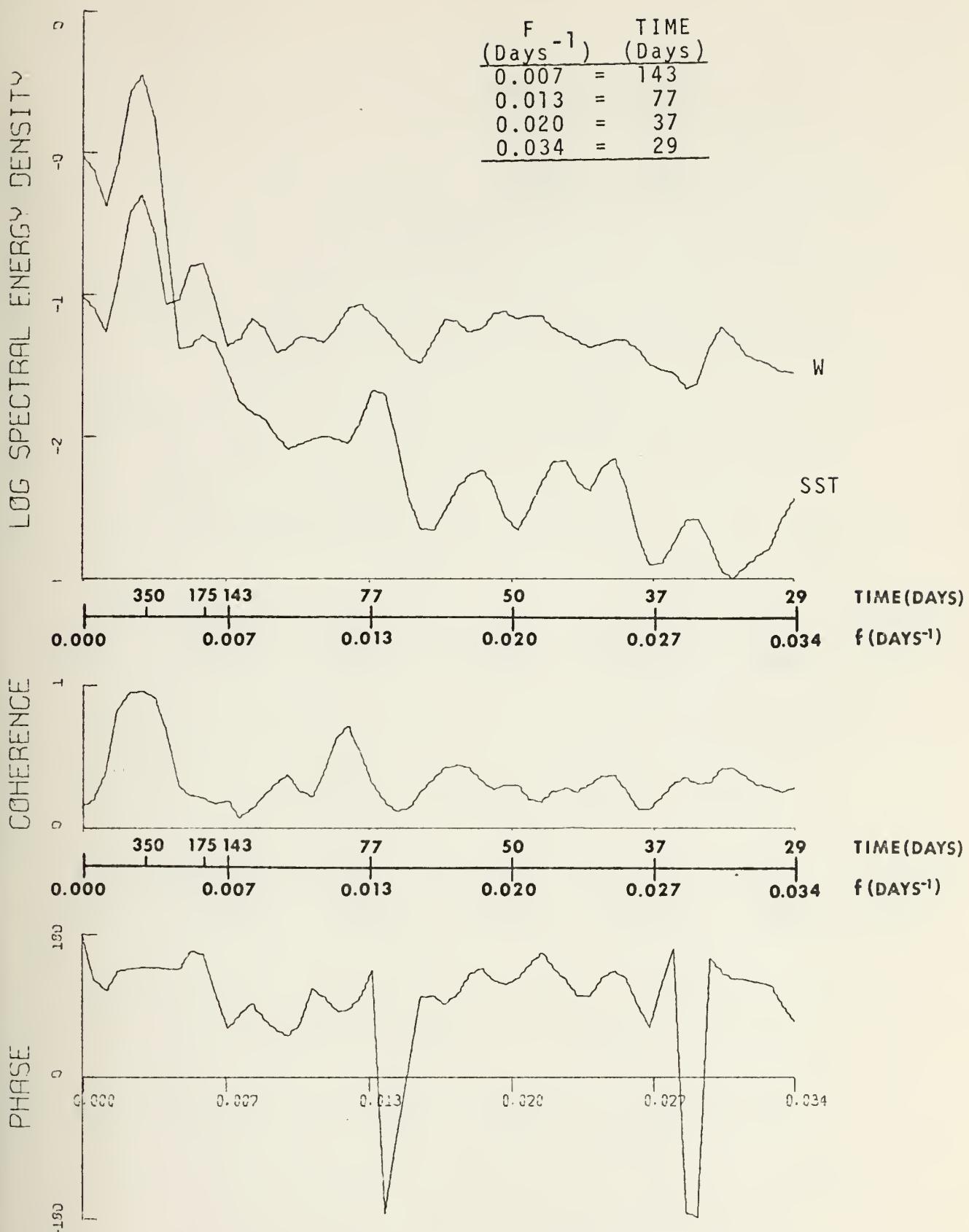


Figure 18. Spectral density, coherence and phase for the SST vs the wind speed.

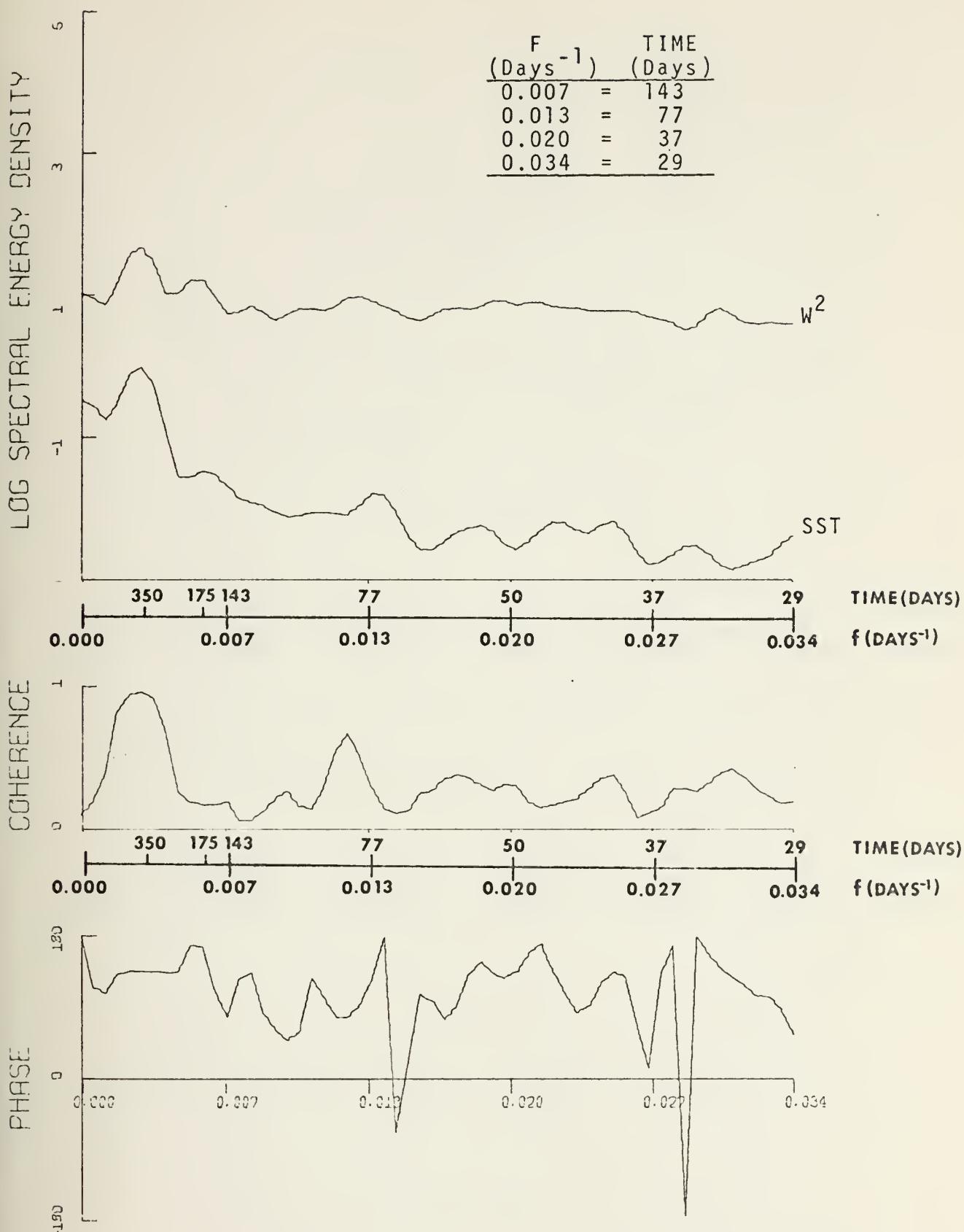


Figure 19. Spectral density, coherence and phase for the SST vs the wind speed squared.

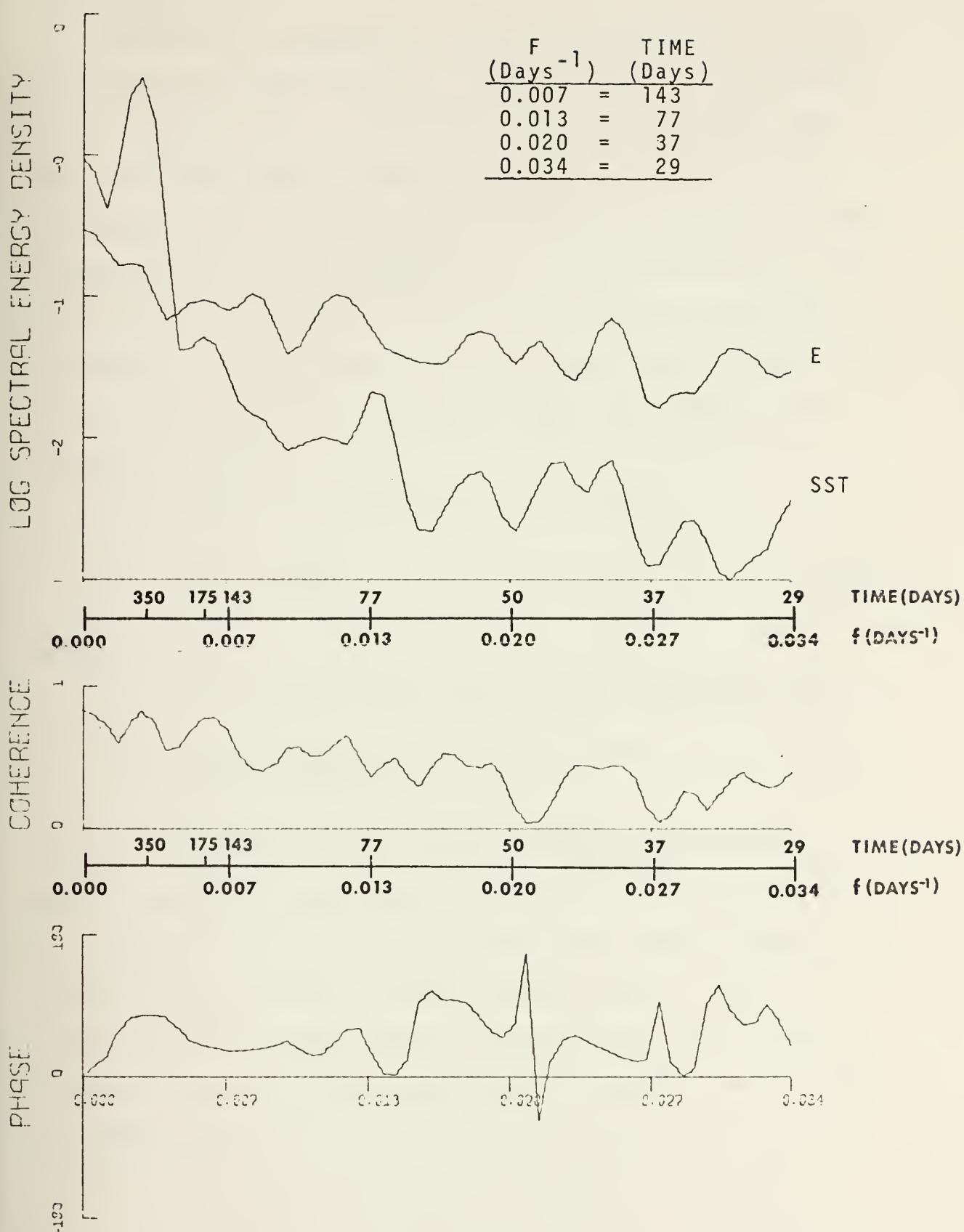


Figure 20. Spectral density, coherence and phase for the SST vs the evaporation.

The heat of conduction, QC, is the smallest of the heat flux terms and only 17% of its total variance is explained by the yearly cycle. From Table VIII it can be seen that this term also has an unusually large lag, approximately 130 days, and the lowest coherence of any of the heat budget terms with the surface temperature. From Figure 21 it can be seen that QC is nearly of the form of white noise with no distinct periodicities. This is confirmed by Figure B-5 which shows that QC is constantly small and contains few anomalies of note.

Like QE, QC is a flux term of the form

$$QC = -A_h \frac{\partial T}{\partial z} \quad (29)$$

where

A_h is the eddy coefficient of heat conduction, and
 $\frac{\partial T}{\partial z}$ is the temperature gradient of the air overlying the sea surface.

For an area of the nature of OWS NOVEMBER, dominated by fairly constant temperature patterns for both the air and the ocean, it can be seen that this term would be expected to be fairly constant. It does contain quite a deal of energy in cycles other than the annual and shows evidence of good coherence at the semi-annual, seasonal and monthly periodicities.

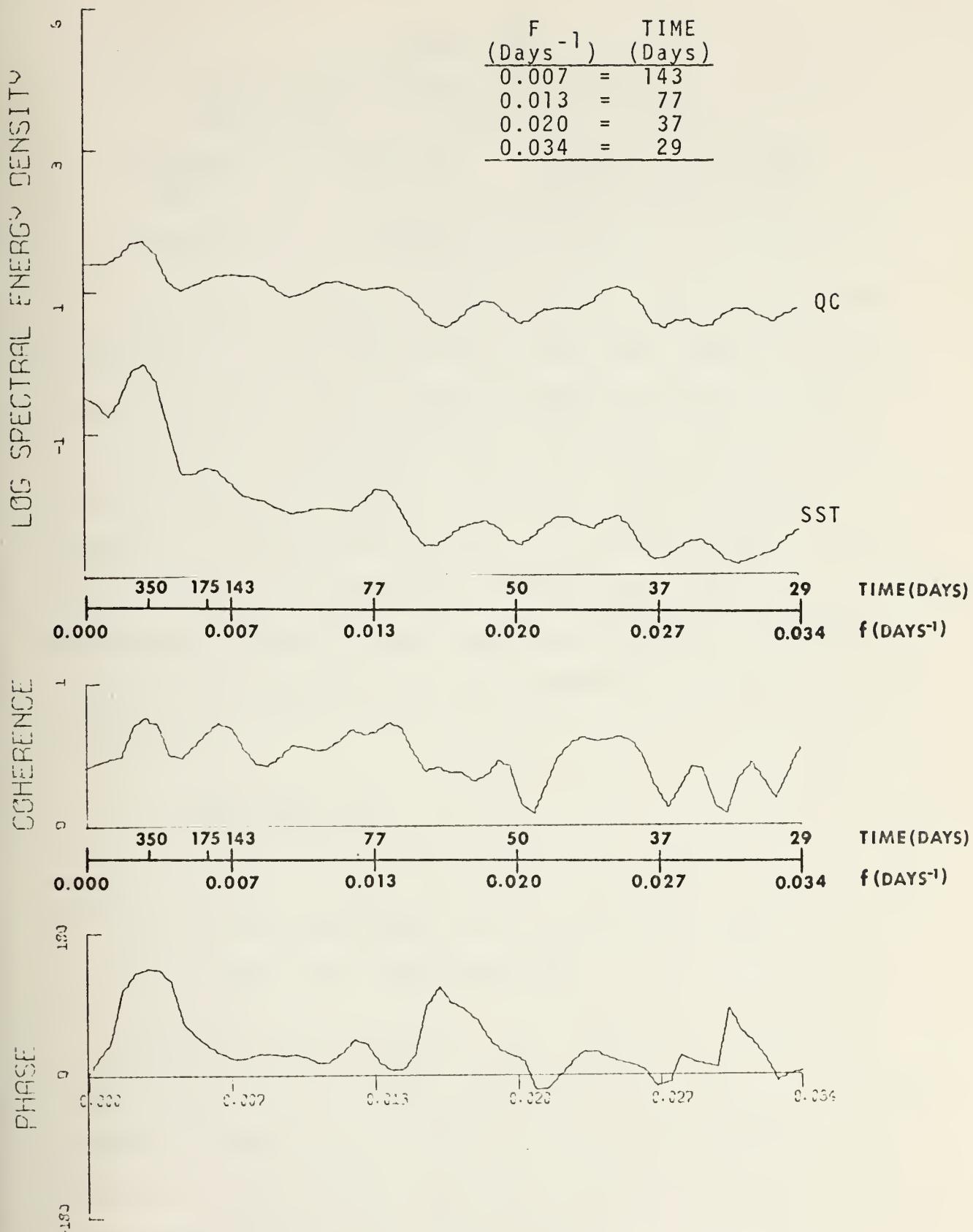


Figure 21. Spectral density, coherence and phase for the SST vs the convective heat.

TABLE VIII.

Spectral results for the convective heat vs the SST at the yearly cycle

Parameter Codes	Coherence	Phase (degrees)	Lag (days)
QC vs SST	.75694	134.1	130.7

The final heat budget term, the net heat transfer, QN , is perhaps the most interesting of the heat equation terms, representing the gain or loss of heat by the water column. To this point, the advective and local heat exchange terms have been ignored. However, since the values for the net radiation computed from Equation (1) did not result in zero values, then it is evident that these terms have to be considered. If QN is then viewed as a combination of the advective heat term and the local generation of heat it takes on the form:

$$QN = Q_v + Q_\theta$$

where:

Q_v is the advection of heat by currents, and

Q_θ is the local heat generated.

Over the annual cycle $Q_\theta = 0$; therefore, QN must be representative of only the advection term. This is supported by the fact that for the entire 24-year period QN was found to have a slope of only 0.003 which signifies that no long term generation of heat has occurred.

From Tables III and IX it can be seen that the net heat exchange is dominated by the annual cycle and is highly coherent with the surface temperature at this period. Forty-two percent of the total variance is explained by the yearly cycle, which allows for some seasonal effects and Figure 22 shows some evidence of high coherence at the semi-annual and seasonal cycles. The lag is confirmed by both Table V and Figure B-6. Much of this lag is undoubtedly due to the effect of the 58-day lag of the solar radiation and the balance by the other heat budget terms. In addition, the lag of QN and the surface pressure, considered here because of the previously noted high correlation of pressure with the surface salinity, may also contribute to the lag of QN and SST.

TABLE IX

Spectral results for the net radiation and related parameters at the yearly cycle

Parameter Codes	Coherence Squared	Phase (degrees)	Lag (days)
QN vs SST	.97954	- 76.7	- 74.8
PS vs SST	.79993	-120.6	-118.6
QN vs PS	.79250	- 46.9	- 45.3

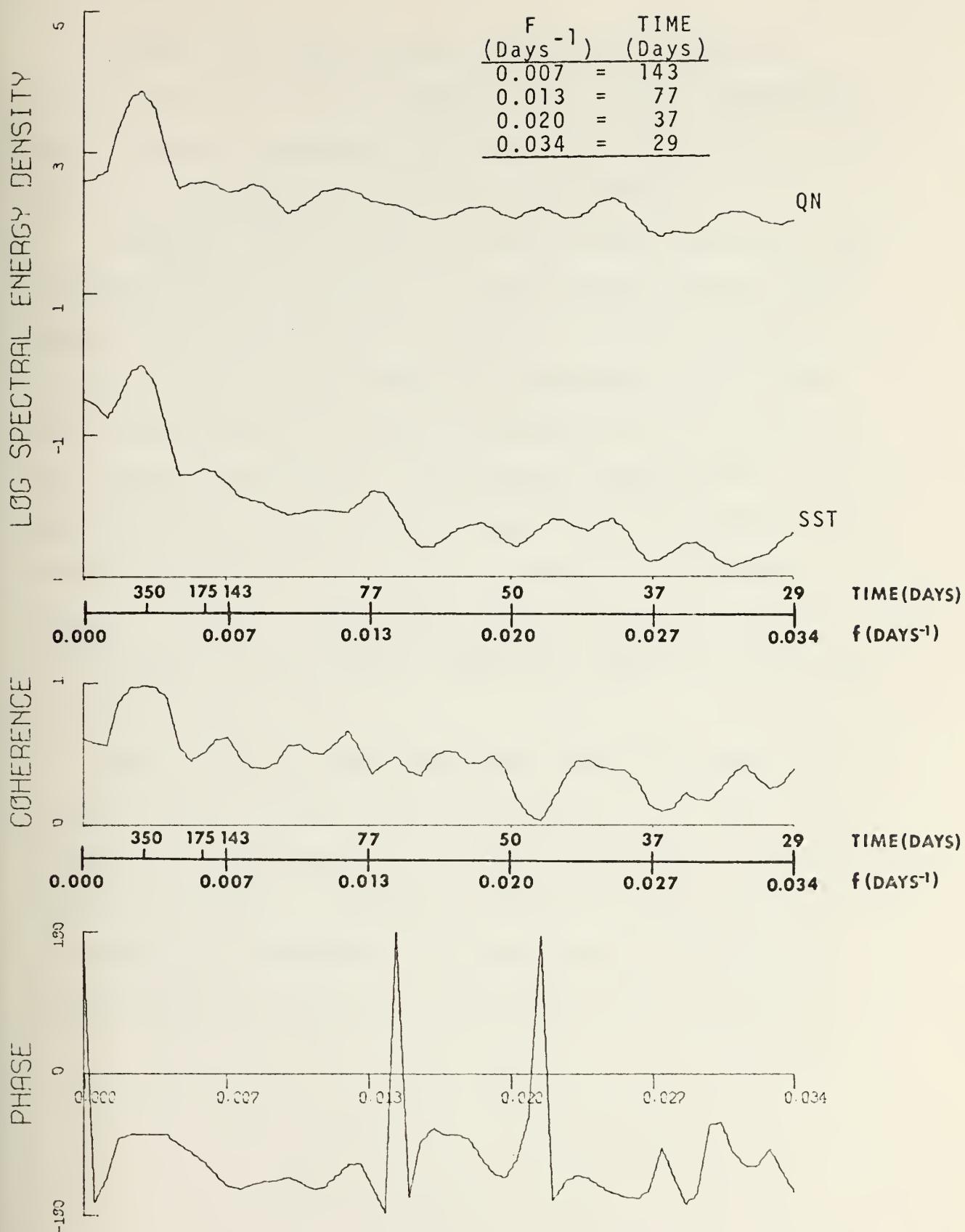


Figure 22. Spectral density, coherence and phase for the SST vs the net radiation.

Figures 22 through 24 give the complete spectral results for the net heat exchange, the surface temperature, and the surface pressure correlated versus each other. In all three instances the yearly peak dominates.

The major result of this segment is that the new solar radiation formula results in larger positive values of QN. Secondly, for the yearly cycle QN is a measure of heat advection and at OWS NOVEMBER this advection is of nearly equal importance as the solar radiation as an influence on the surface temperature. The negative lag makes QN a fairly good predictor of the surface temperature, even though much of the variance of the net heat exchange is contained in the seasonal cycles. The surface pressure seems to be less correlated with the surface temperature than it is with the salinity.

Very little has been mentioned about the autocorrelations; however, White (1974) has used this function in the study of SST and baroclinic transport anomalies. He found an apparent 5-year periodicity in the SST anomalies in the vicinity of OWS NOVEMBER. This cycle was attributed to deeper anomalous features in the baroclinic structure.

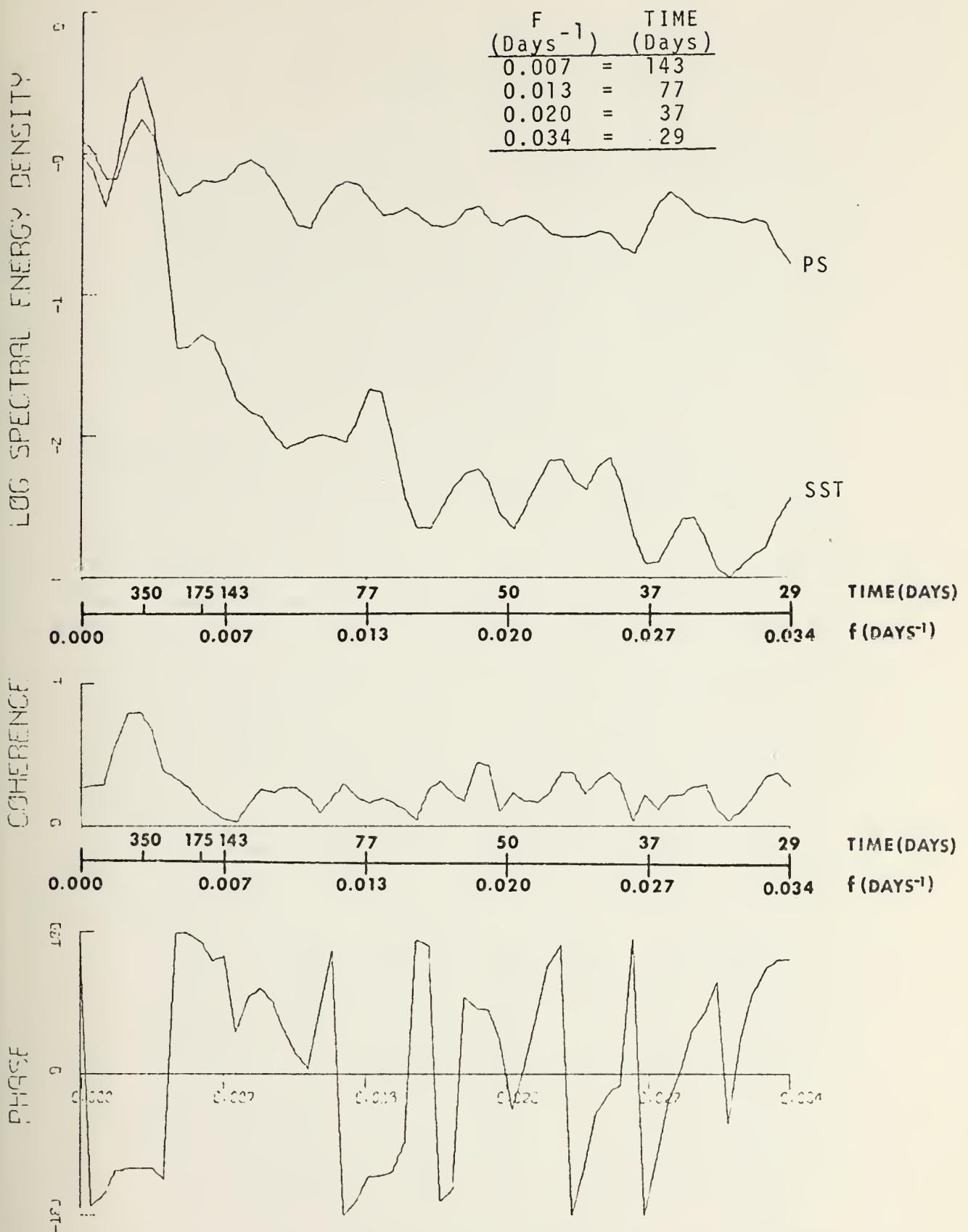


Figure 23. Spectral density, coherence and phase for the SST vs the surface pressure.

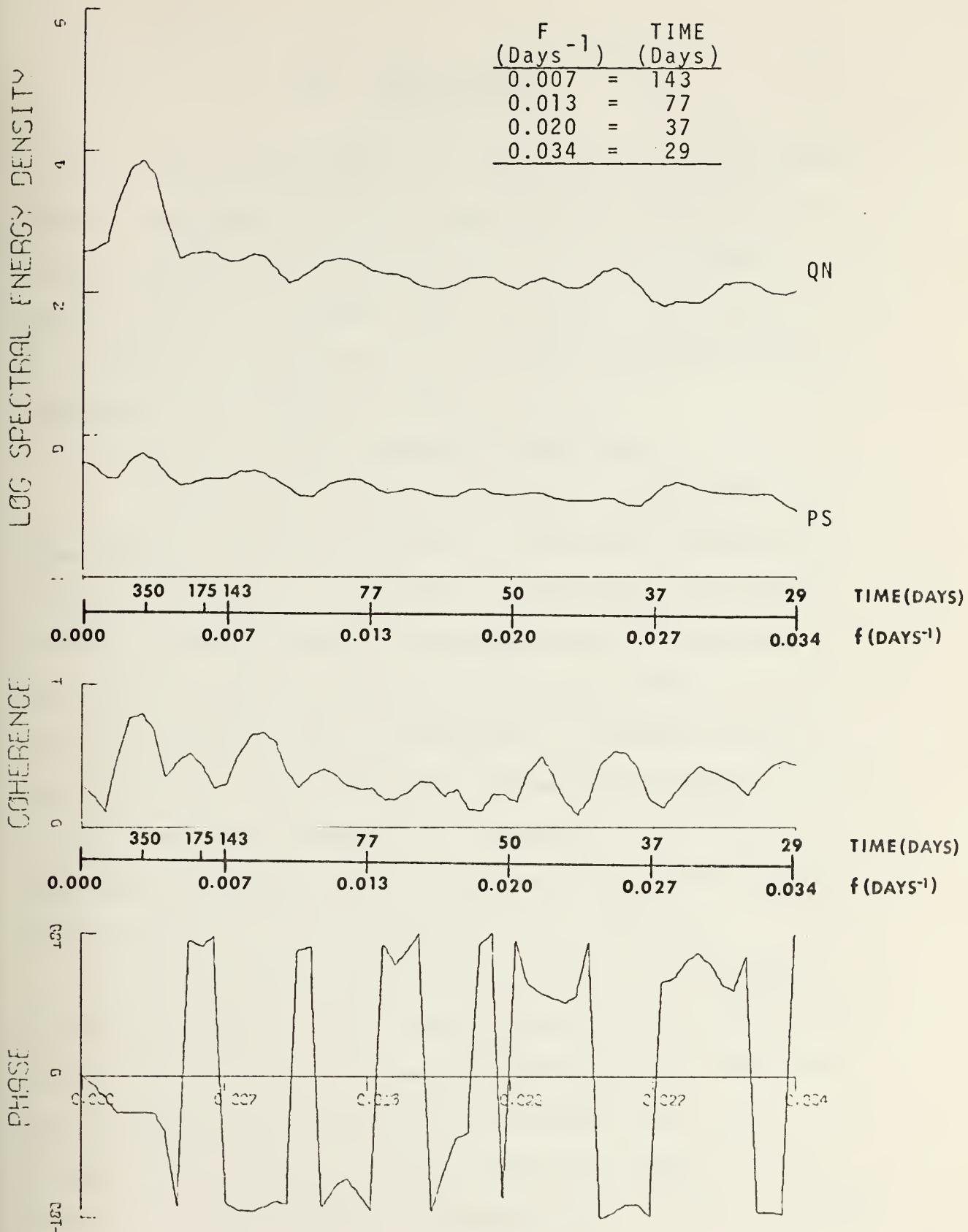


Figure 24. Spectral density, coherence and phase for the net radiation vs the surface pressure.

IV. CONCLUSIONS

The aim of this study was to use the tools of spectral analysis and other statistical approaches to describe the interactions of the air-sea interface at OWS NOVEMBER, specifically as to finding those parameters most closely associated with fluctuations of the daily average surface temperature and surface salinity.

The results of this research showed that:

1. The surface salinity from 1968 through 1970 was shown through regression analysis to be poorly correlated with the parameters traditionally associated with changes in the observed salinity. The evaporation, evaporation minus precipitation, and precipitation were found to have little correlation with fluctuations in surface salinity. These terms are non-conservative terms associated with the air-sea interface and their low correlation coefficients would appear to negate their importance in controlling the surface salinity at this location in the Pacific.

2. The surface pressure was found to be the most strongly correlated of all the parameters regressed against the surface salinity. It was therefore hypothesized that this represents the presence of a movement of a baroclinic system of cooler water from the subarctic region into the vicinity of OWS NOVEMBER, corresponding to a shift southward of the transition zone.

3. For the yearly cycle, it was shown that the net radiation is a measure of the advection present and that this advection is of nearly equal magnitude as the solar radiation.

4. The new formula for the solar radiation leads to values for this term approximately 20 to 30 percent higher than previously computed. In turn, this leads to much larger positive values of the net heat exchange term and larger positive anomalies for both terms.

5. From the spectral analyses it was found that all the 24-year time series were dominated by the yearly cycle; except for the back radiation, cloud cover, wind speed and wind speed squared series which showed evidence of semi-annual and seasonal cycles as well.

6. As predictors of the surface temperature, the solar radiation and advection in the form of the net heat exchange term were the only terms of suitable phase and coherence at the yearly cycle to be valid. The other heat transfer terms seems to have less effect upon fluctuations of the surface temperature. Several of the observed parameters, such as the air temperature, were also highly coherent and, on the whole, had smaller lags associated with them.

APPENDIX A
THE AUTOCORRELATION FUNCTIONS

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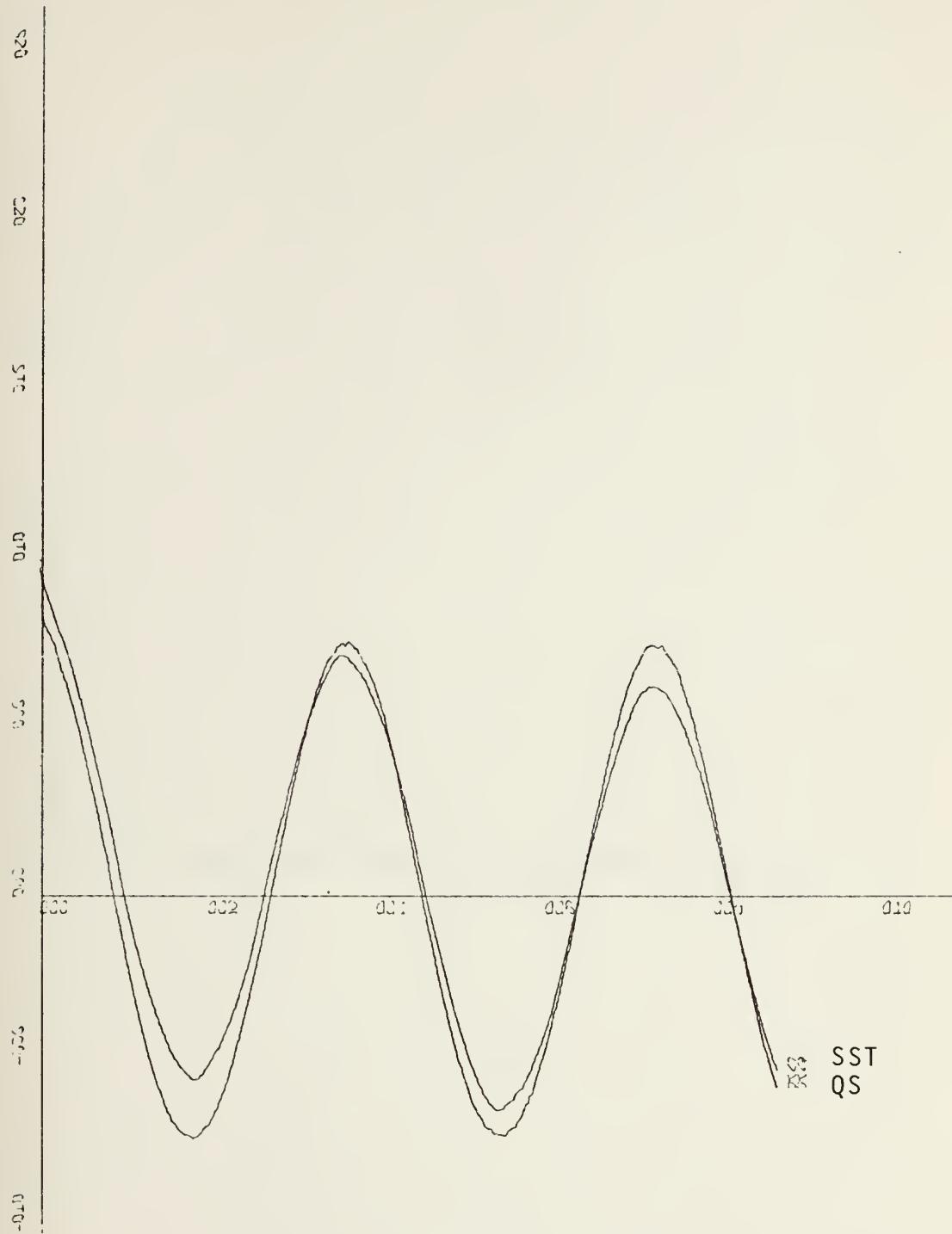


Figure A-1. Autocorrelation curves for the SST and the solar radiation, QS (abscissa is given in hundreds of days).

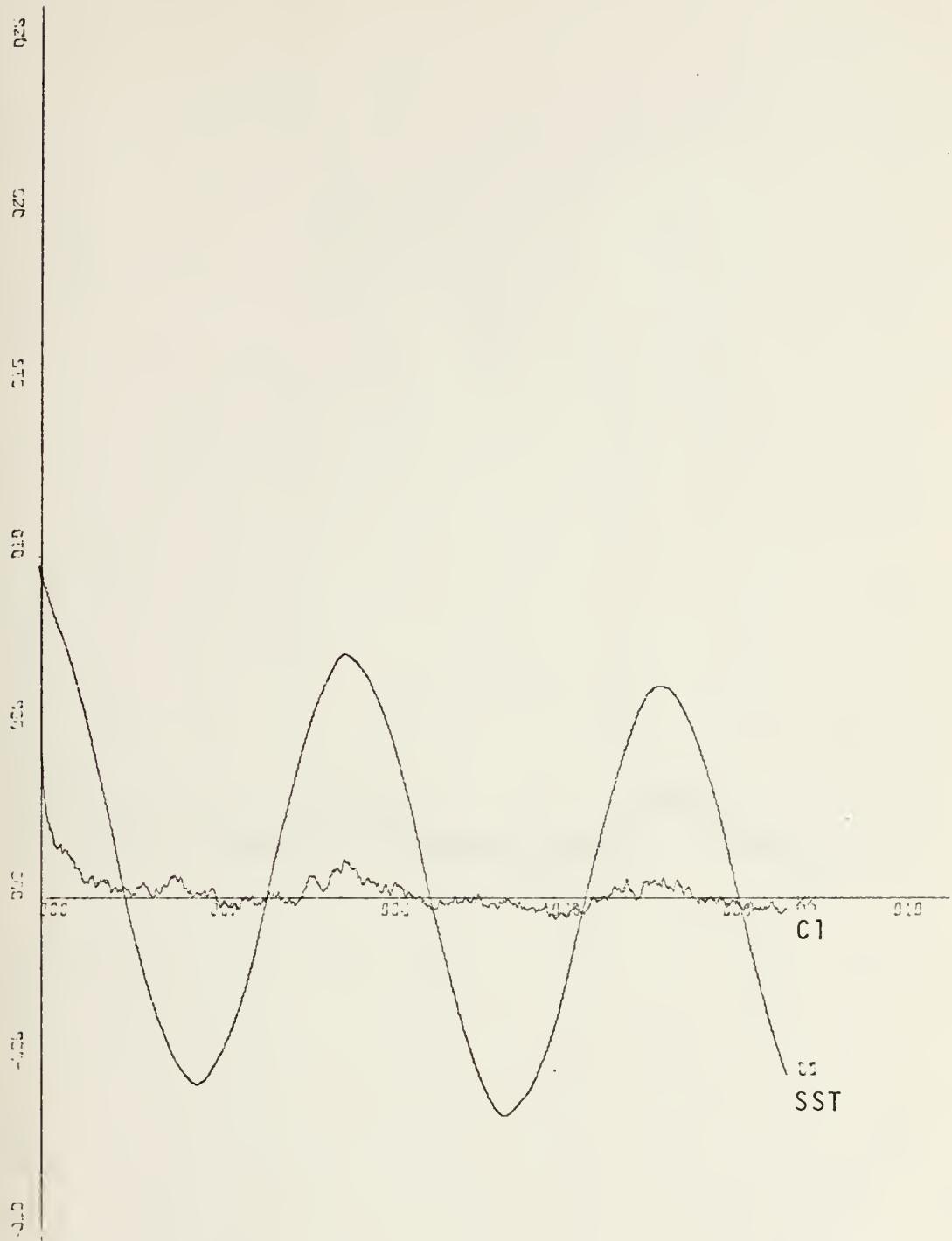


Figure A-2. Autocorrelation curves for the SST and cloud coverage (abscissa is given in hundreds of days).

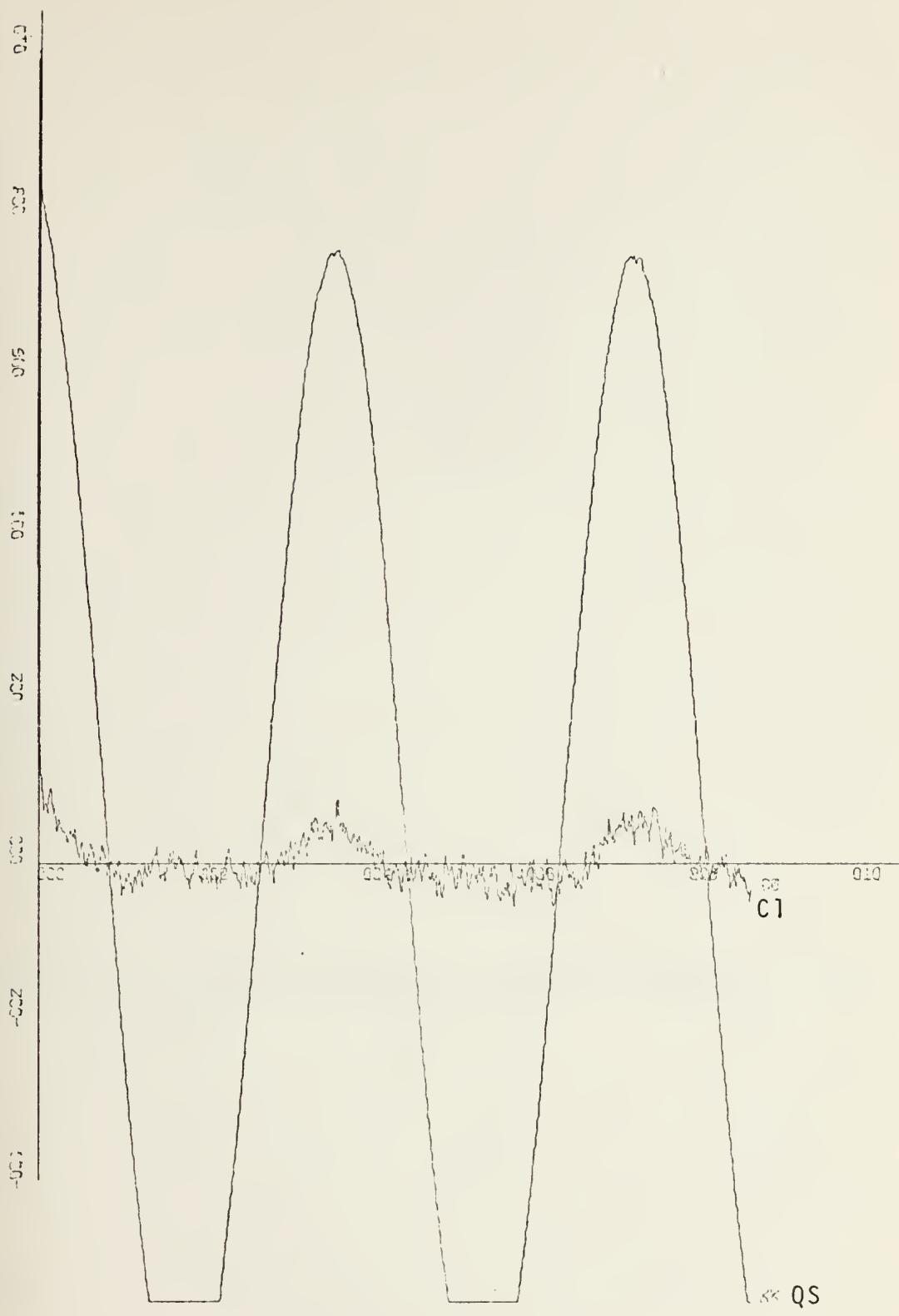


Figure A-3. Autocorrelation curves for the cloud coverage and the solar radiation (abscissa is given in hundreds of days).

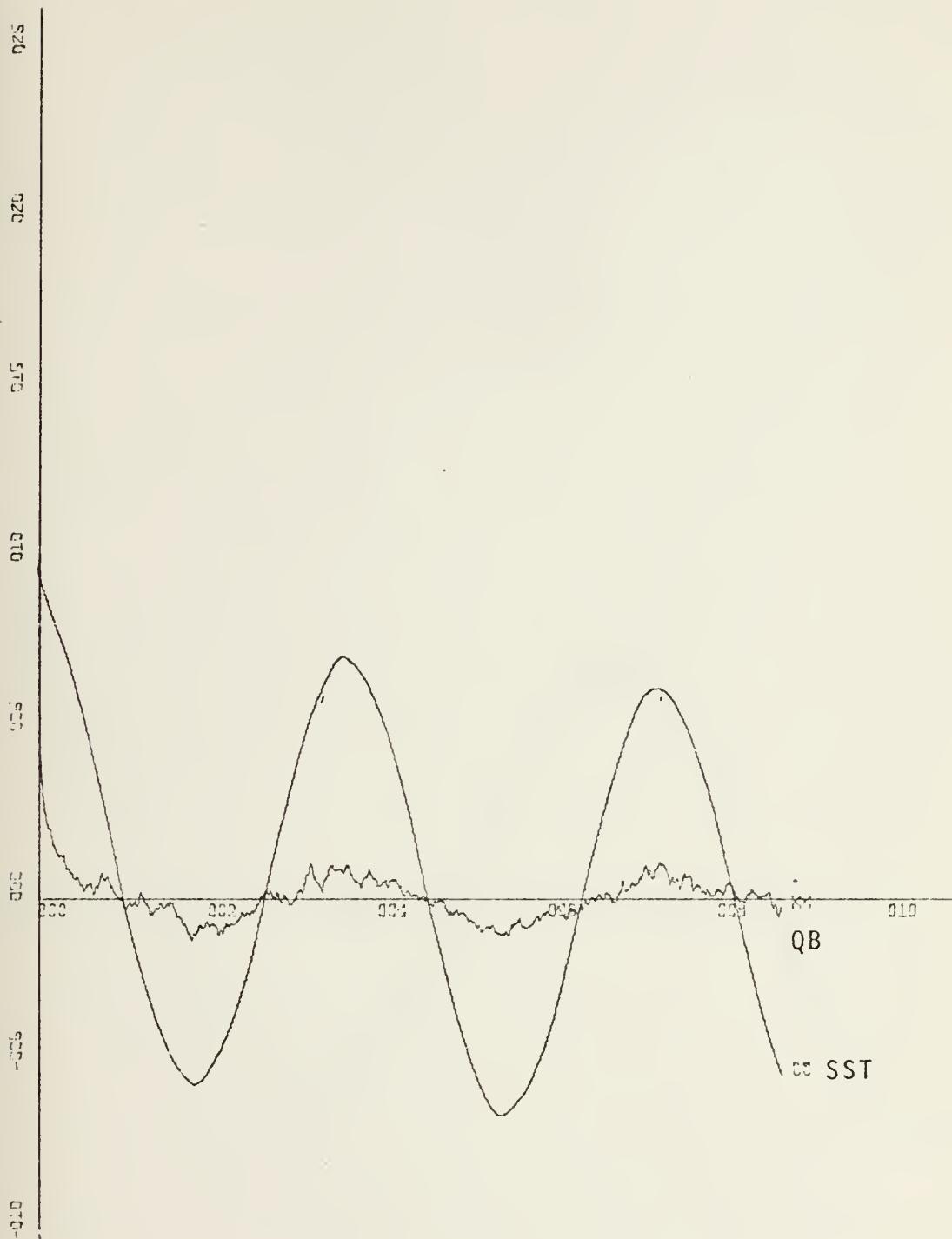


Figure A-4. Autocorrelation curves for the SST and the back radiation (abscissa is given in hundreds of days).

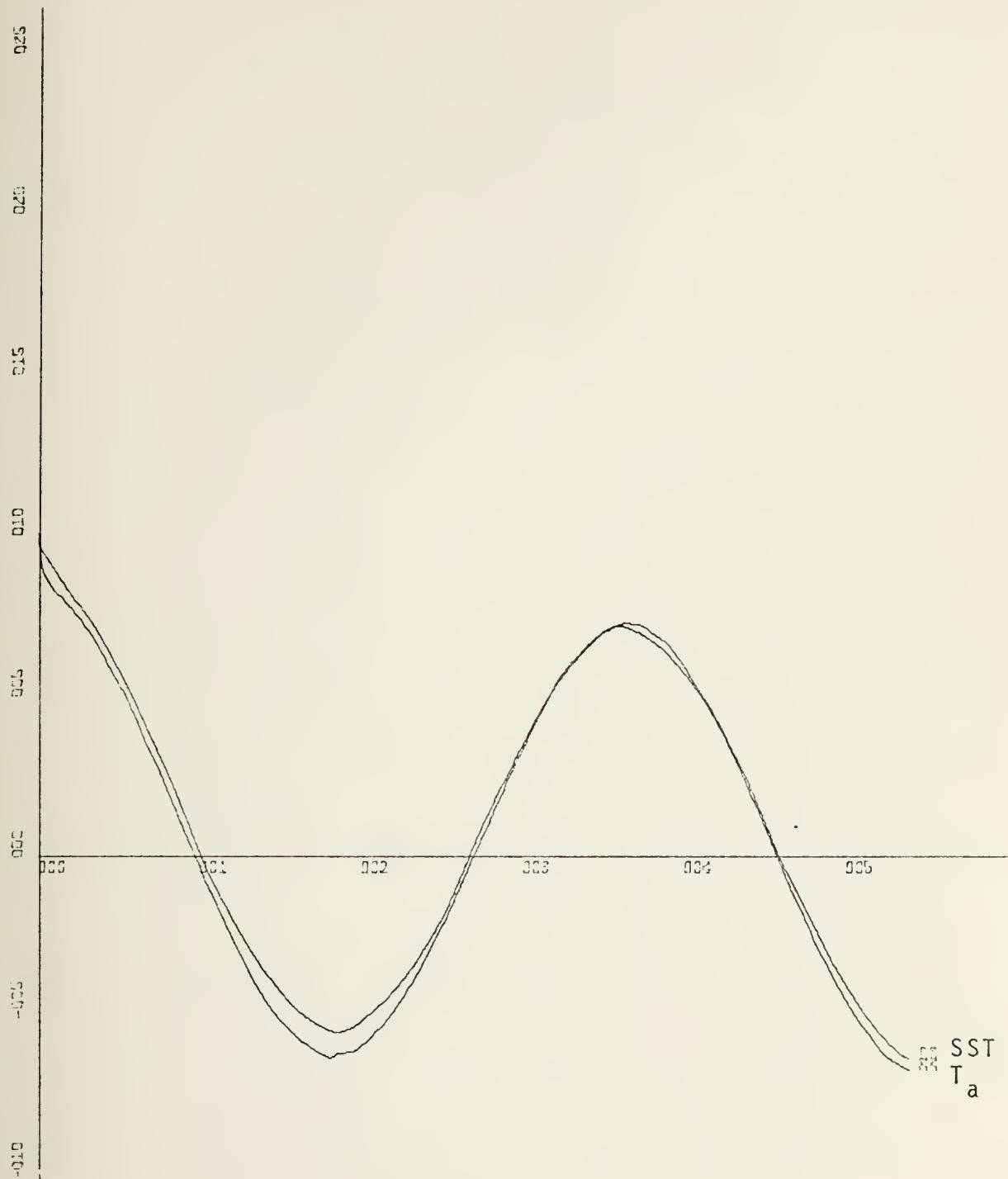


Figure A-5. Autocorrelation curves for the SST and the air temperature (abscissa is given in hundreds of days).

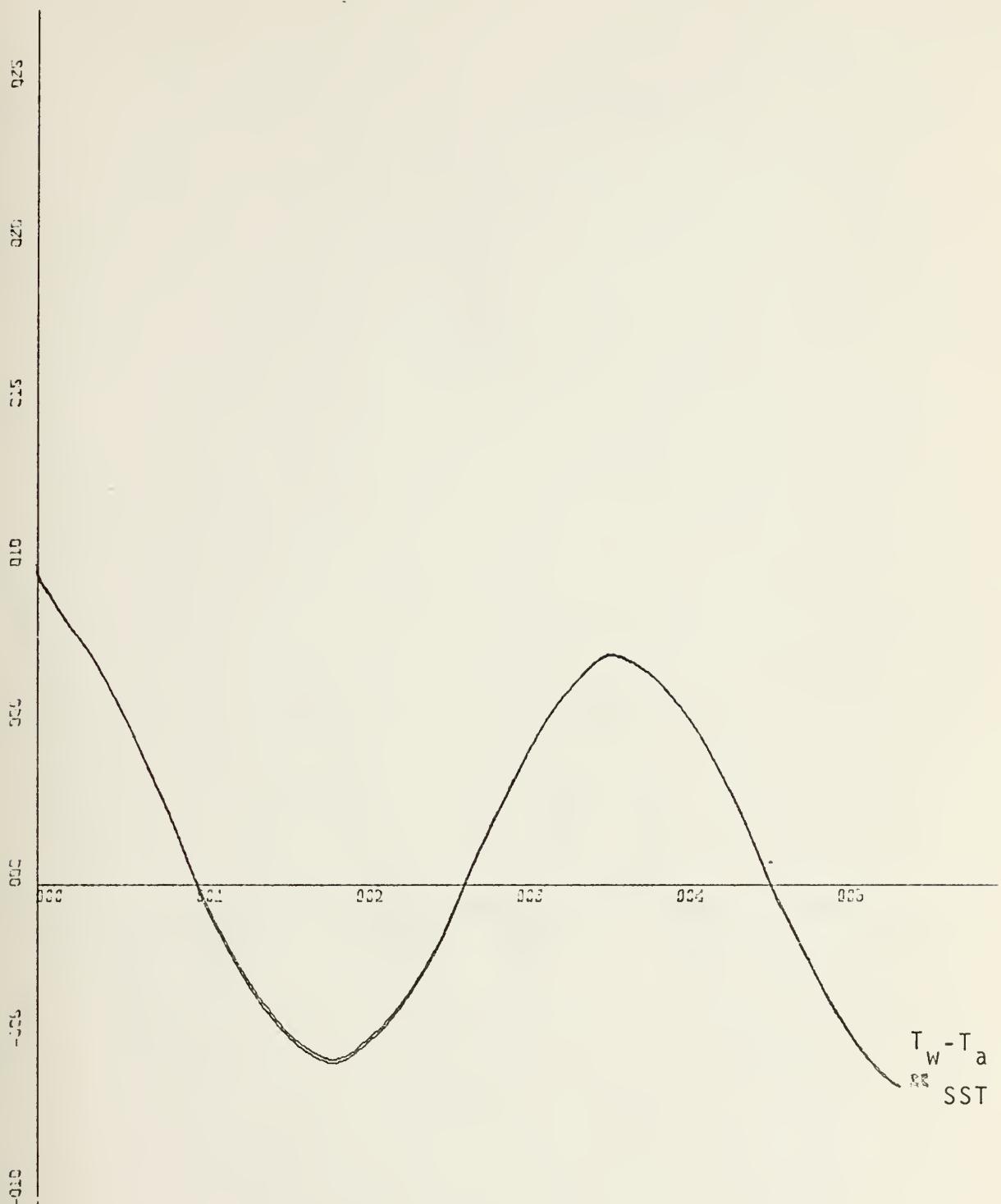


Figure A-6. Autocorrelation curves for the SST and the air-sea temperature difference (abscissa is given in hundreds of days).

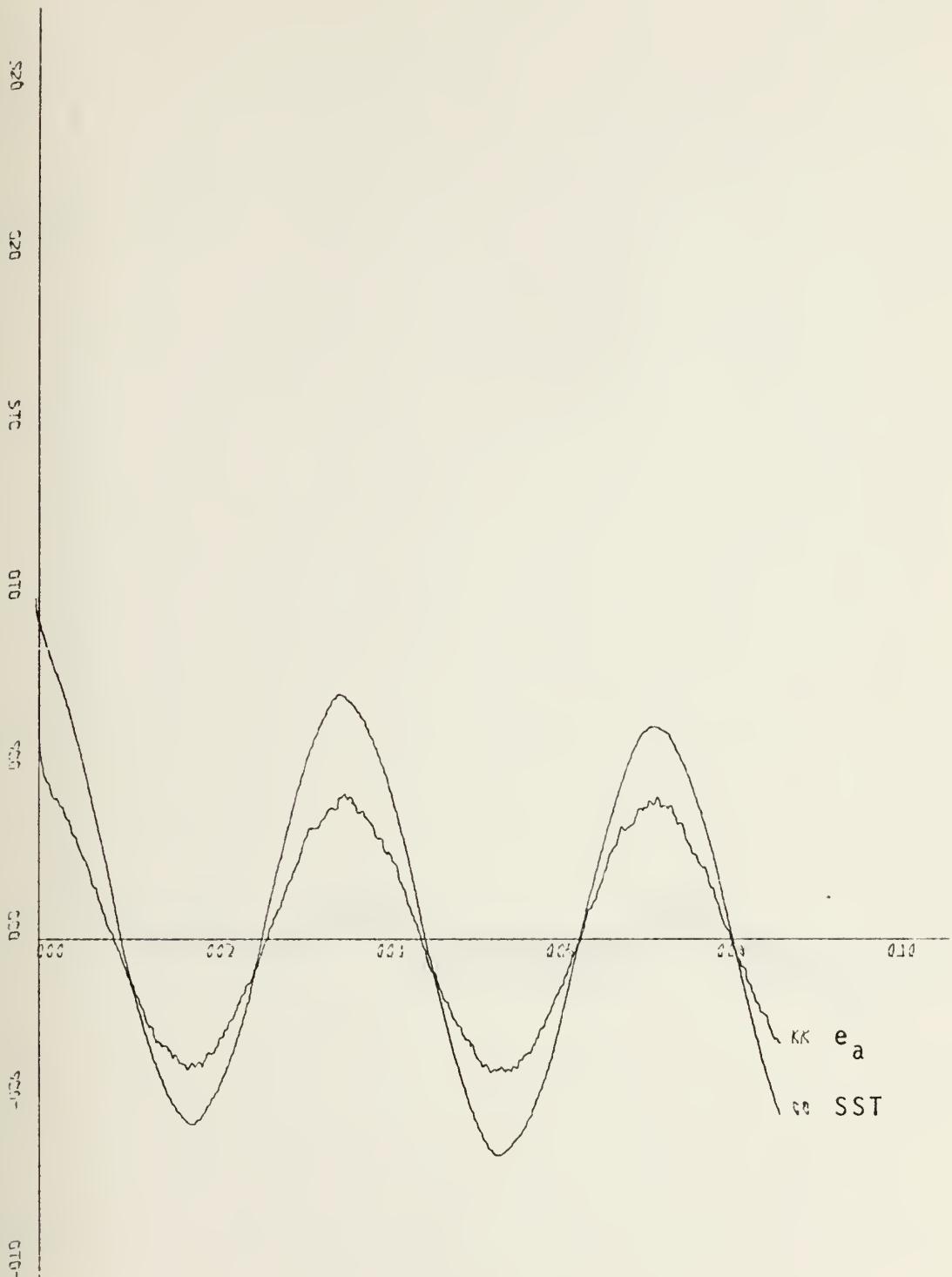


Figure A-7. Autocorrelation curves for the SST and the vapor pressure of the air (abscissa is given in hundreds of days).

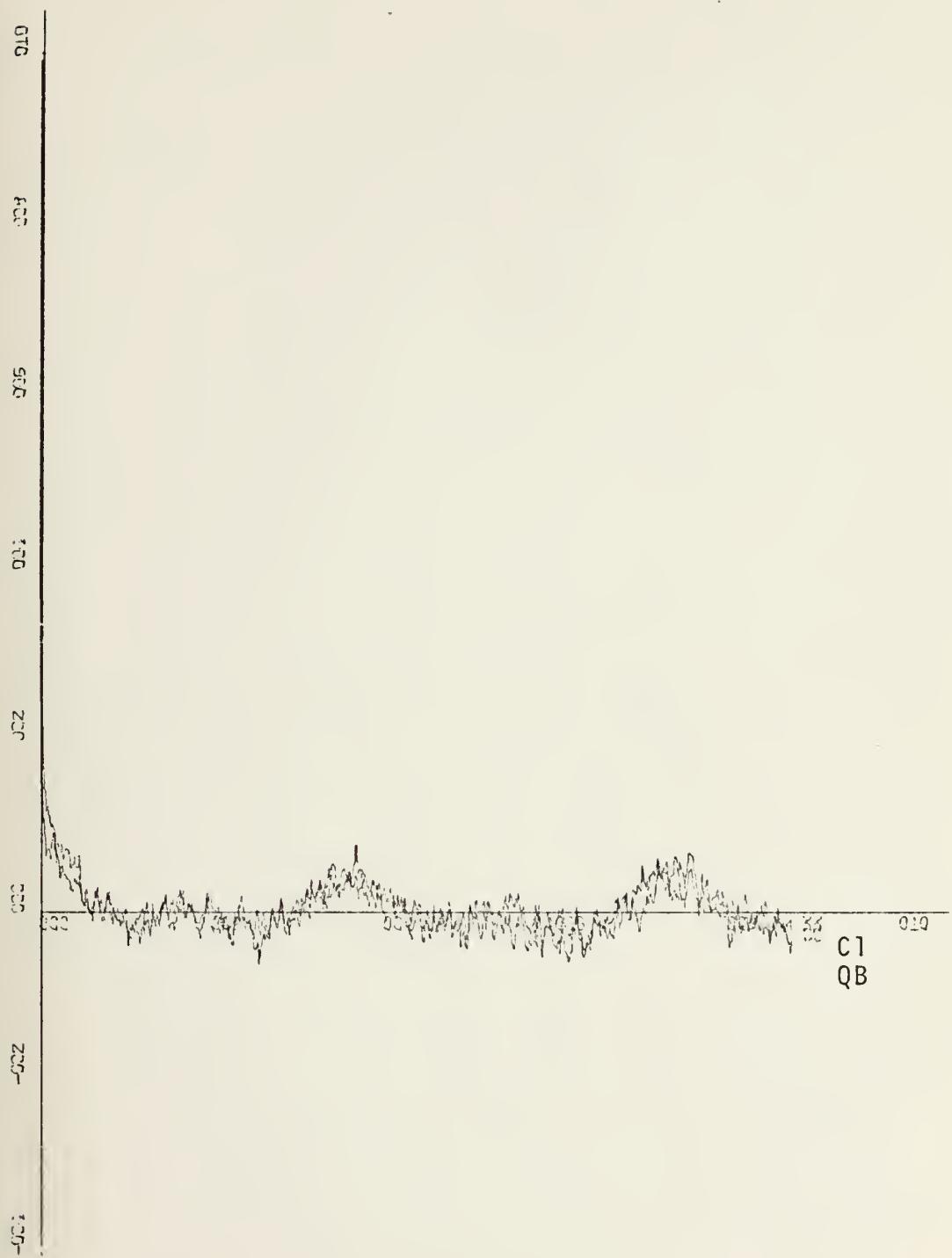


Figure A-8. Autocorrelation curves for the back radiation and cloud cover (abscissa is given in hundreds of days).

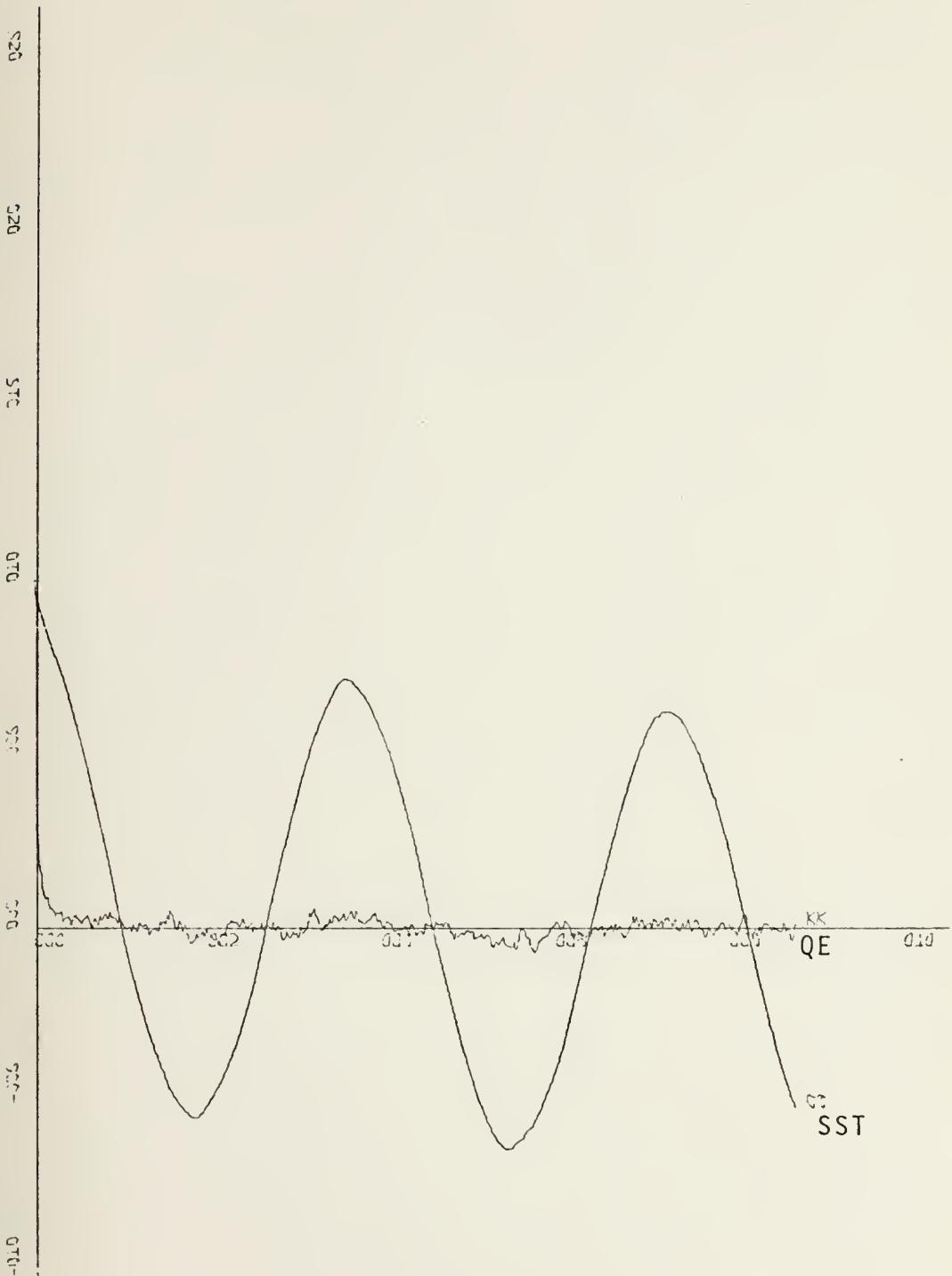


Figure A-9. Autocorrelation curves for the SST and the heat of evaporation (abscissa is given in hundreds of days).

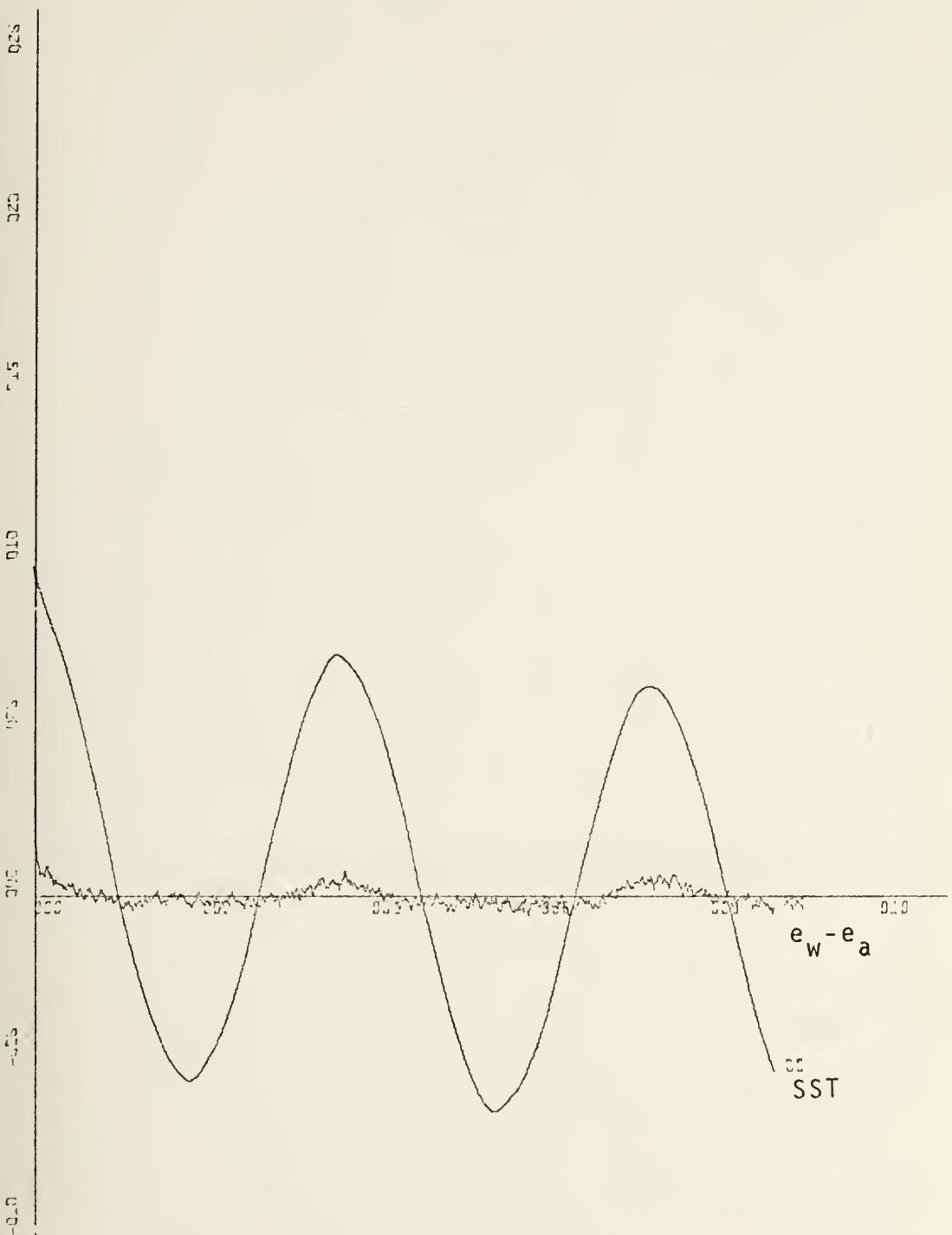


Figure A-10. Autocorrelation curves for the SST and the vapor pressure difference (abscissa is given in hundreds of days).

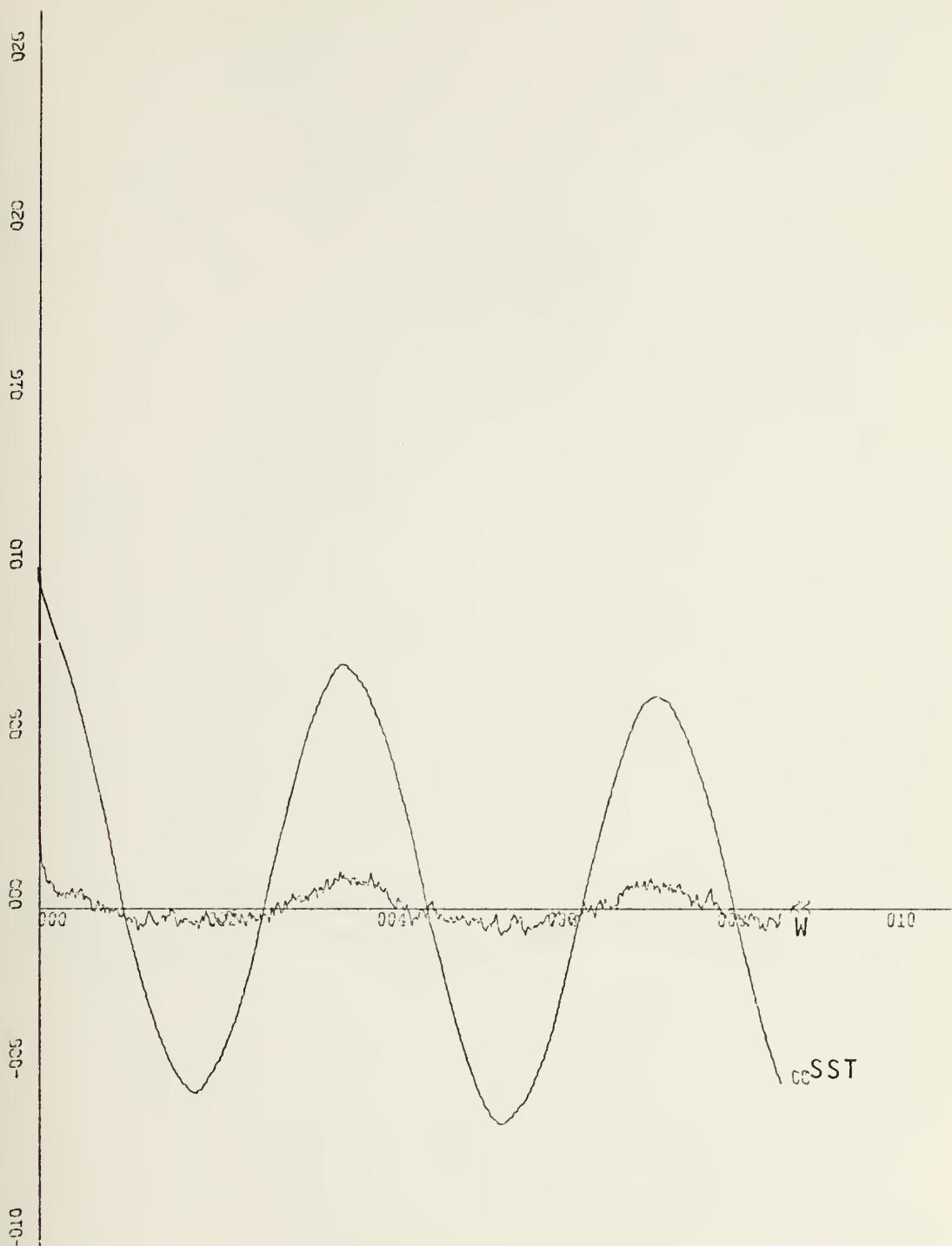


Figure A-11. Autocorrelation curves for the SST and the wind speed (abscissa is given in hundreds of days).

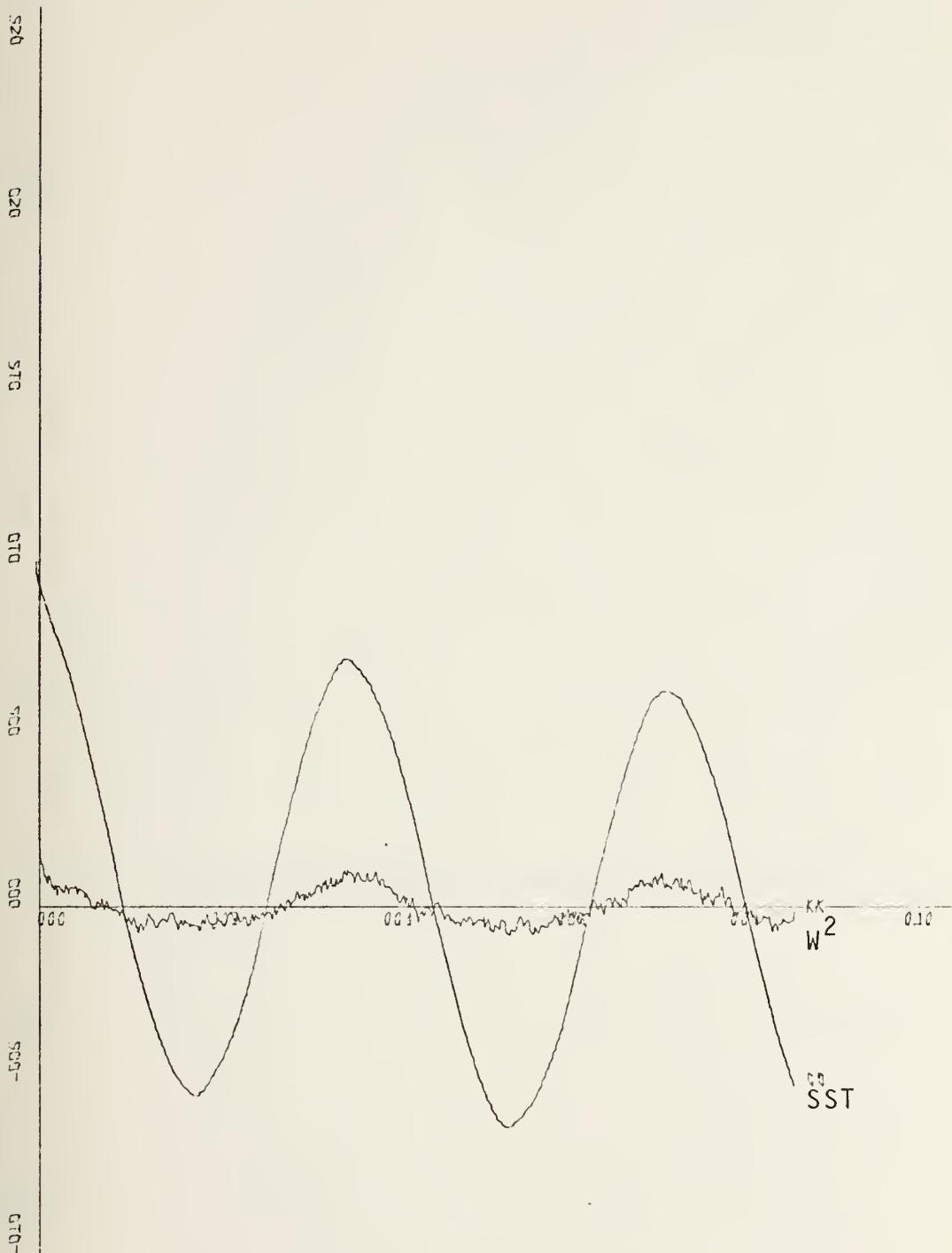


Figure A-12. Autocorrelation curves for the SST and the wind speed squared (abscissa is given in hundreds of days).

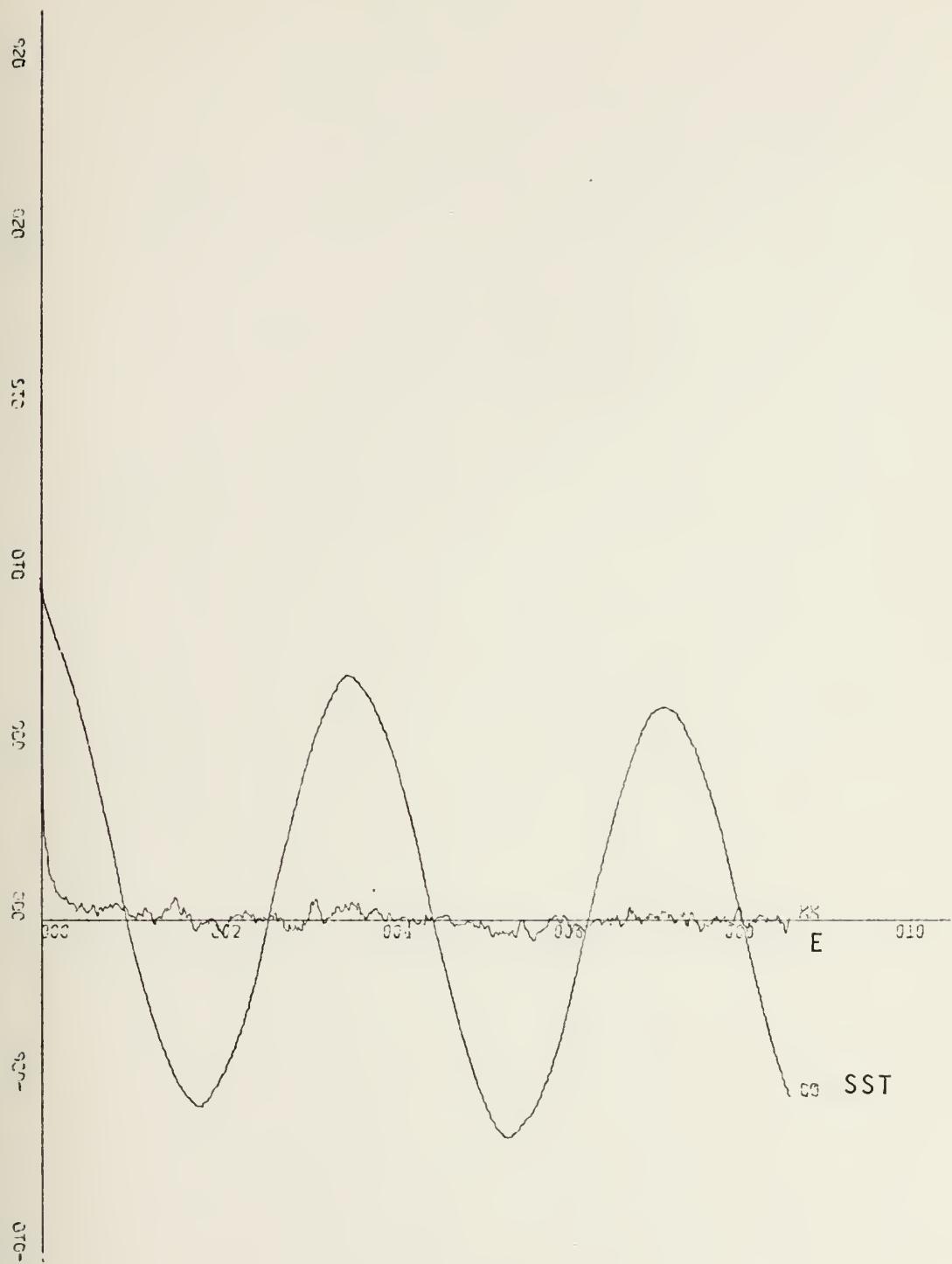


Figure A-13. Autocorrelation curves for the SST and evaporation (abscissa is given in hundreds of days).

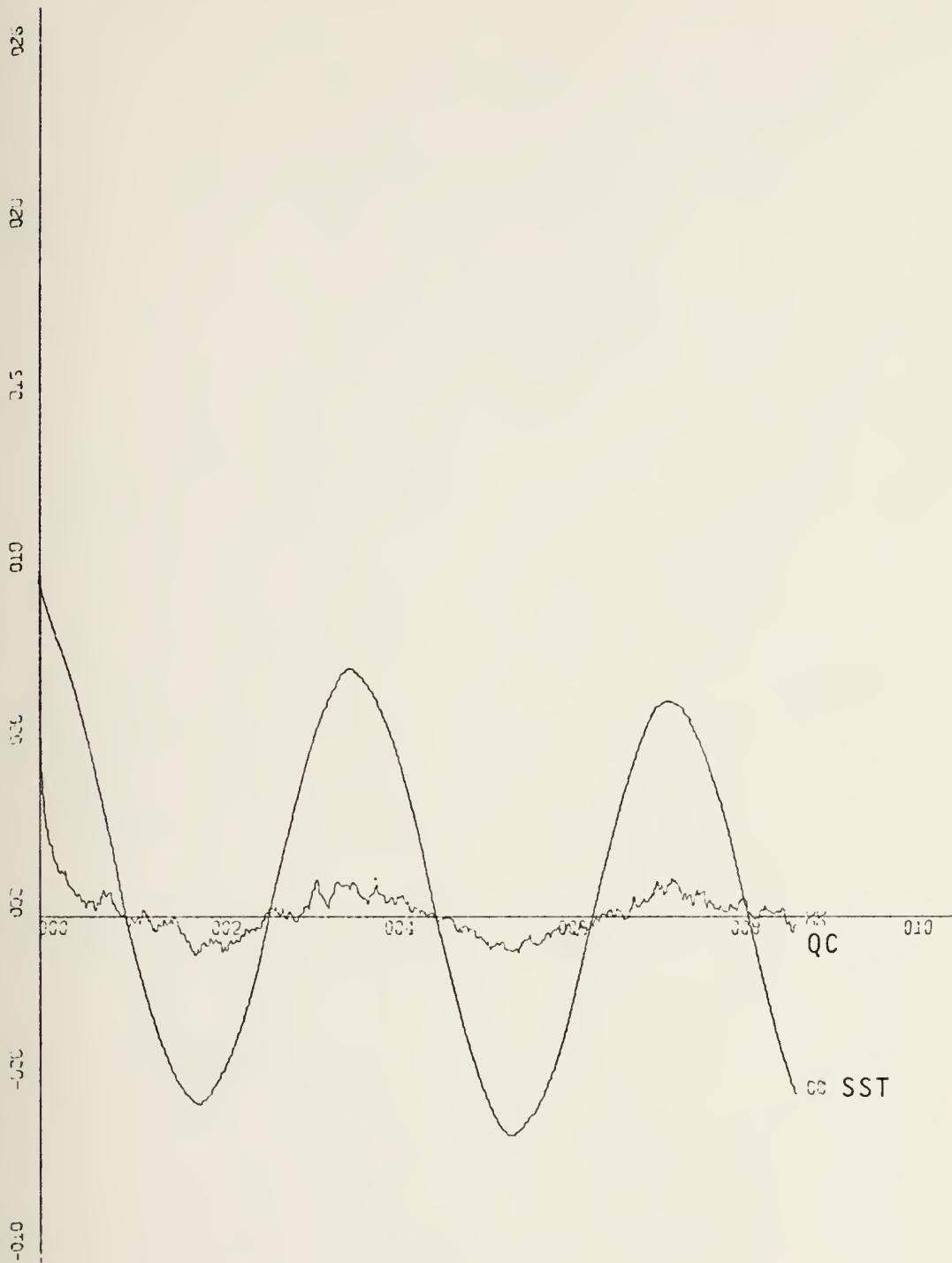


Figure A-14. Autocorrelation curves for the SST and the convective heat (abscissa is given in hundreds of days).

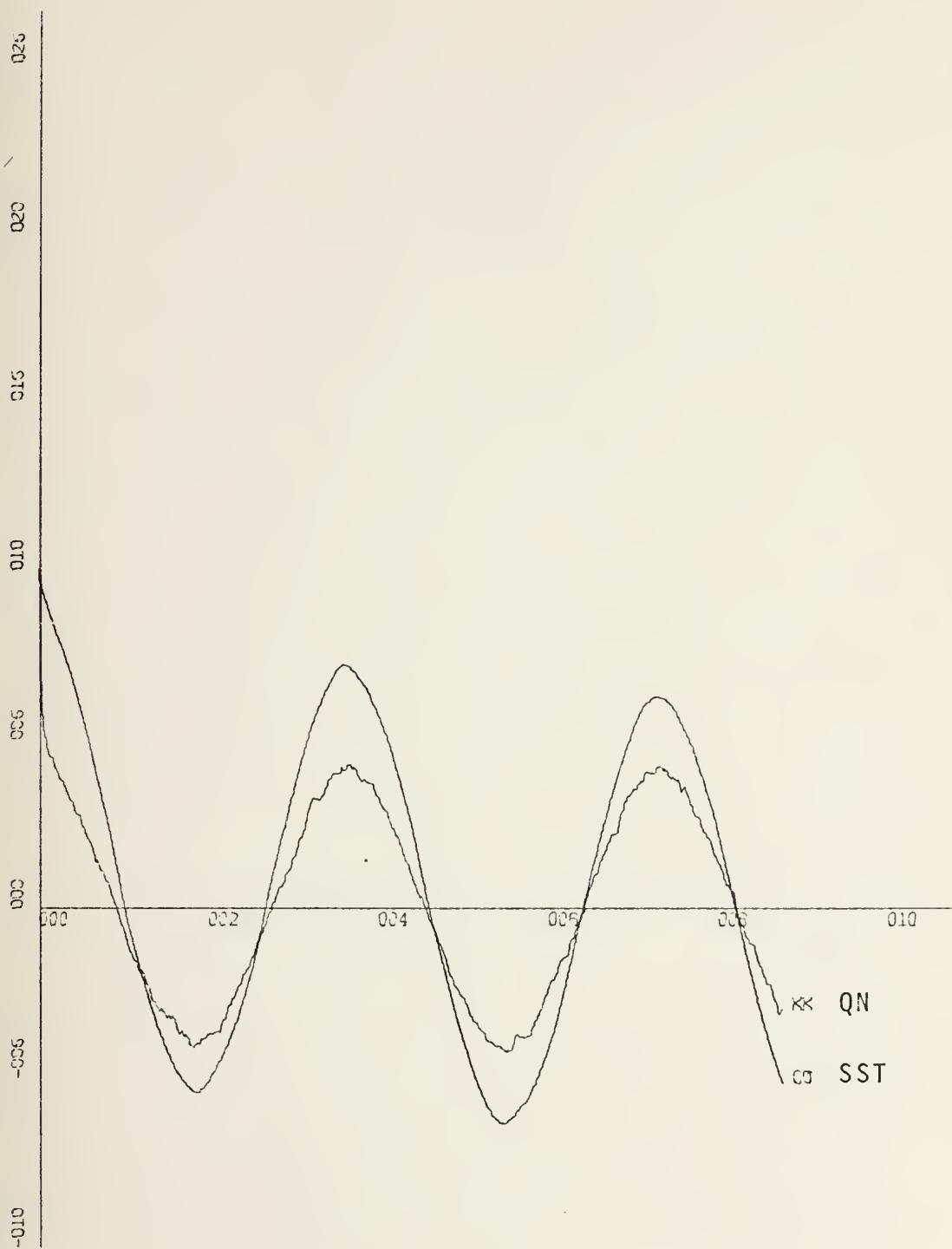


Figure A-15. Autocorrelation curves for the SST and the net radiation (abscissa is given in hundreds of days).

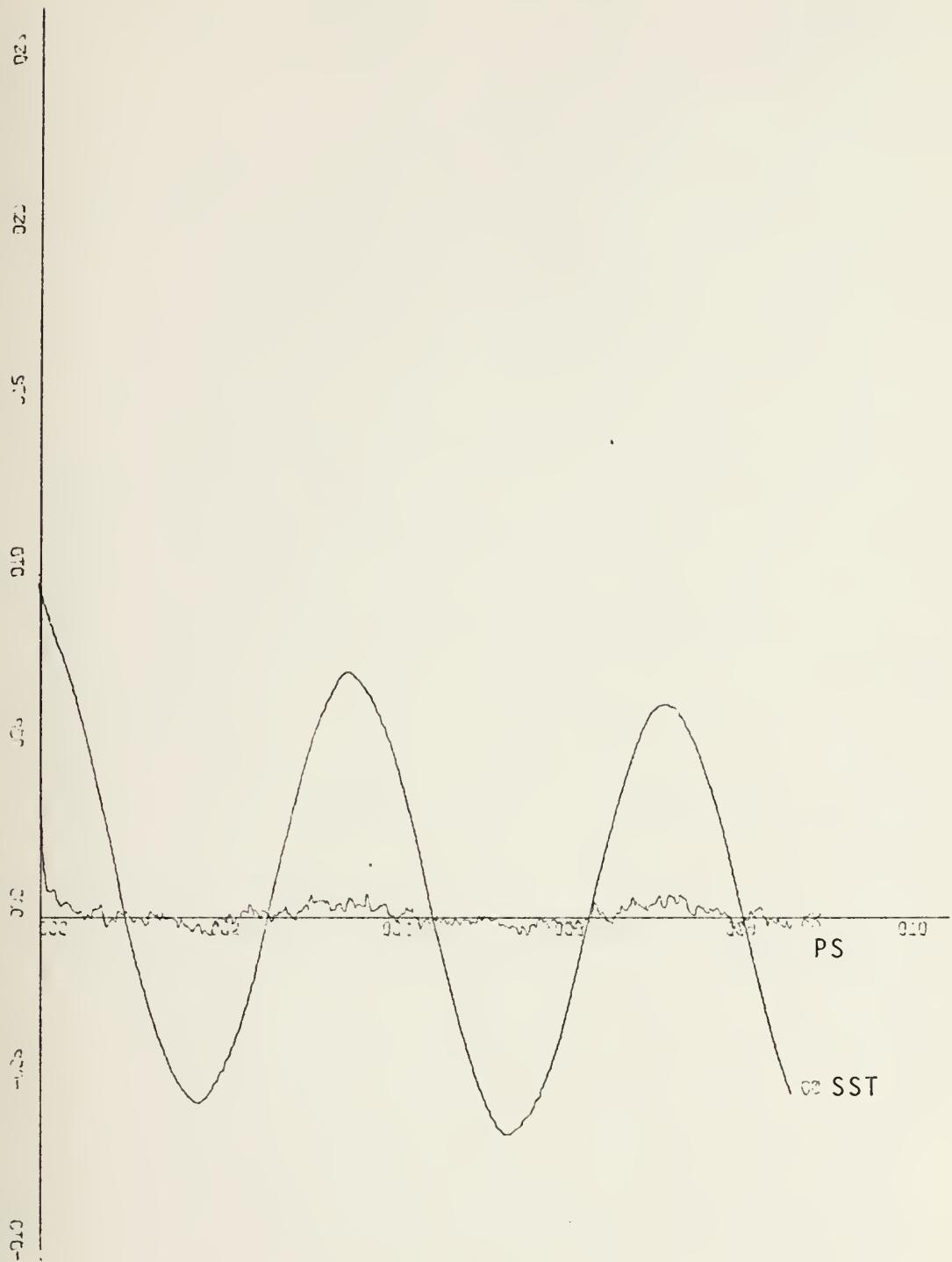


Figure A-16. Autocorrelation curves for the SST and surface pressure (abscissa is given in hundreds of days).

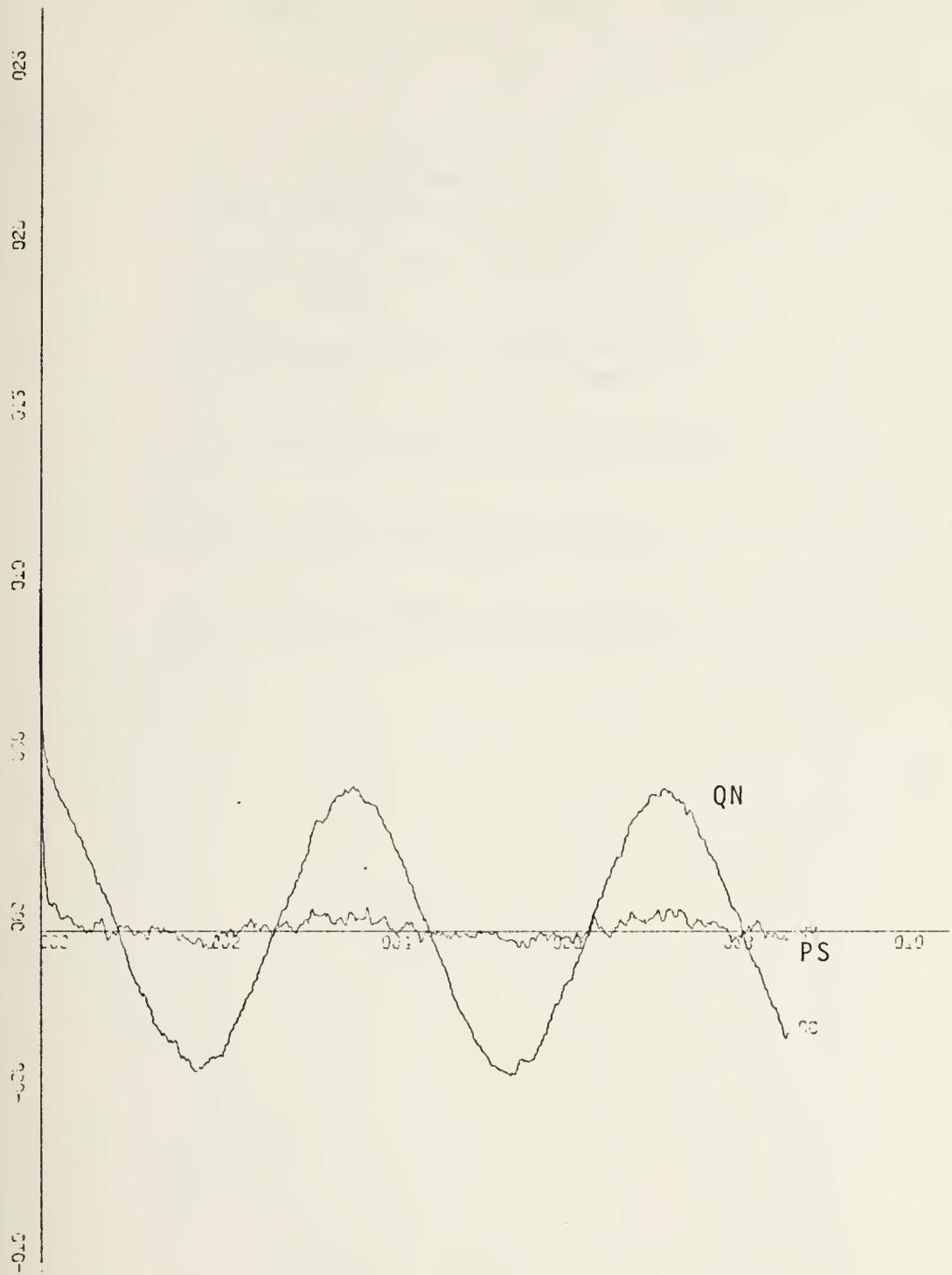


Figure A-17. Autocorrelation curves for the net radiation and the surface pressure (abscissa is given in hundreds of days).

APPENDIX B
FIFTEEN-YEAR TIME SERIES SEGMENTS

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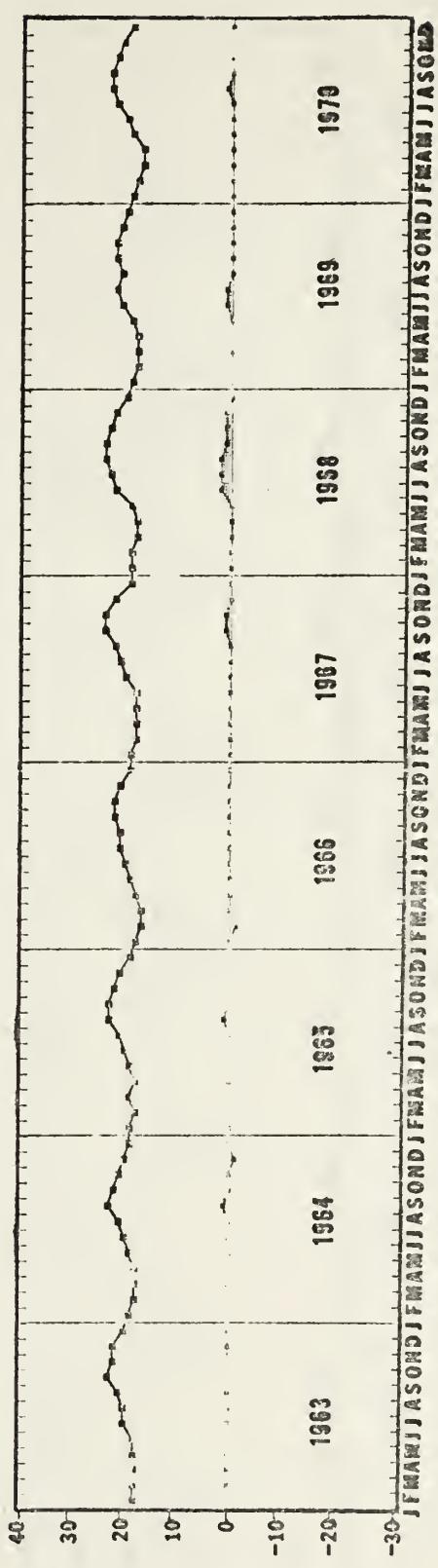
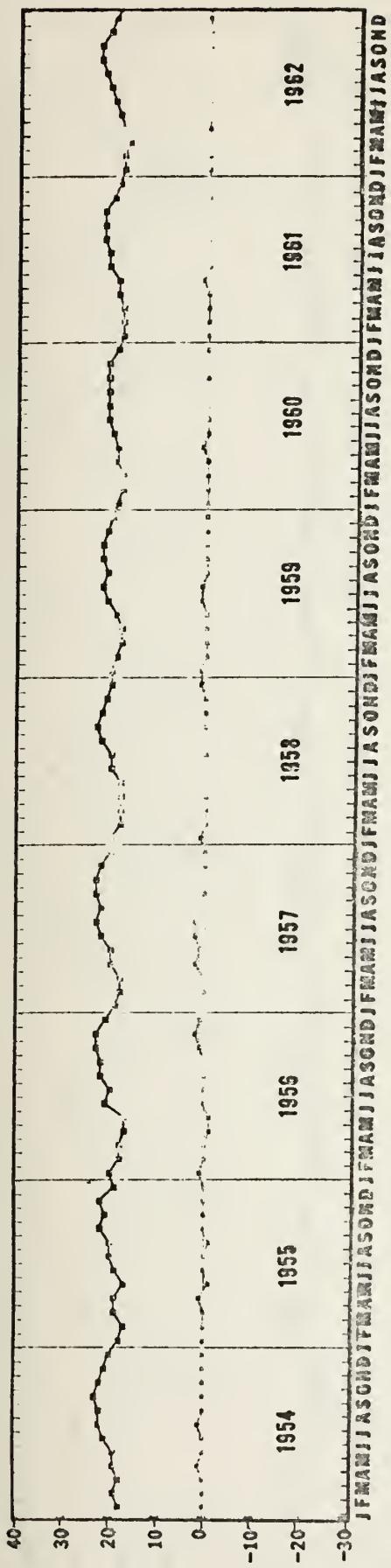


Figure B-1. Fifteen-year time series of the surface temperature and its anomaly ($^{\circ}\text{C}$).

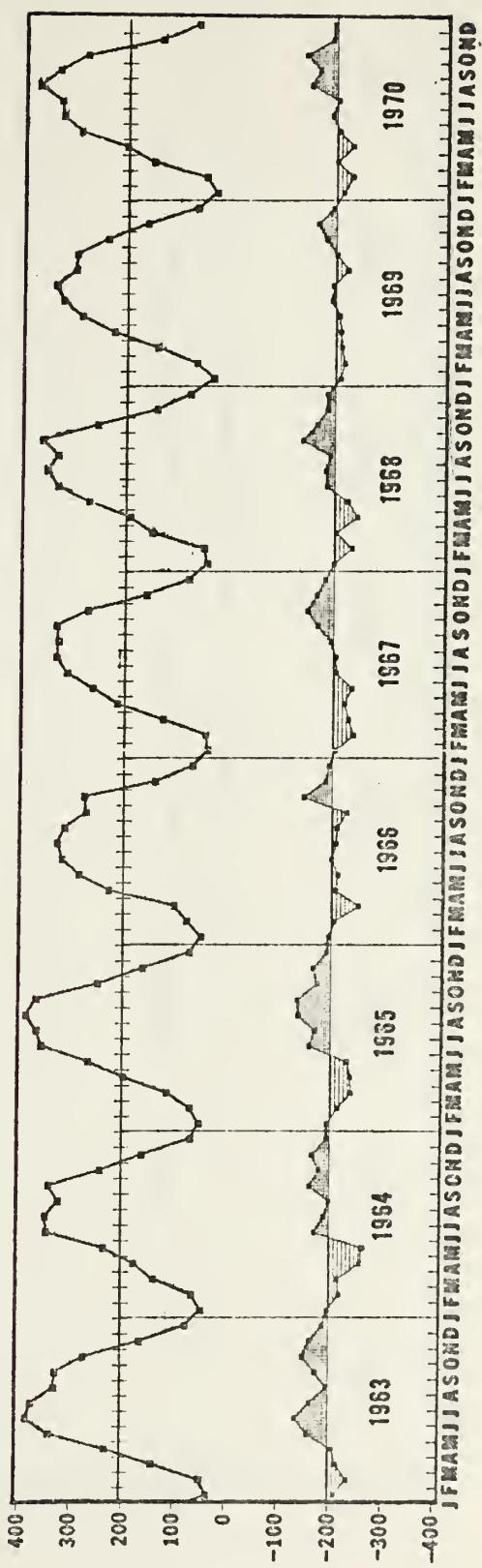
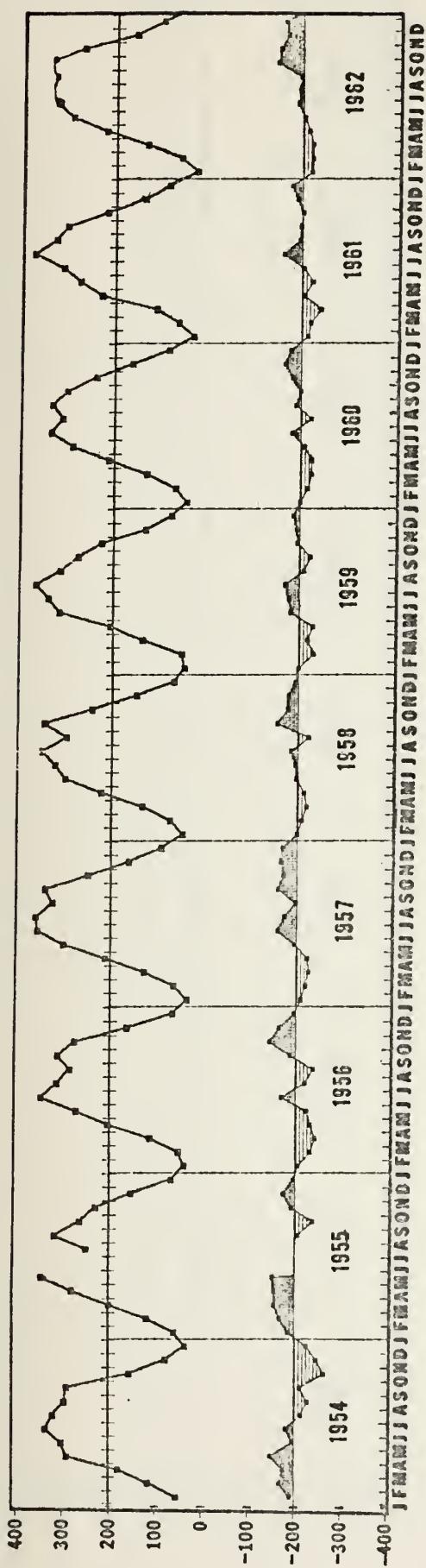
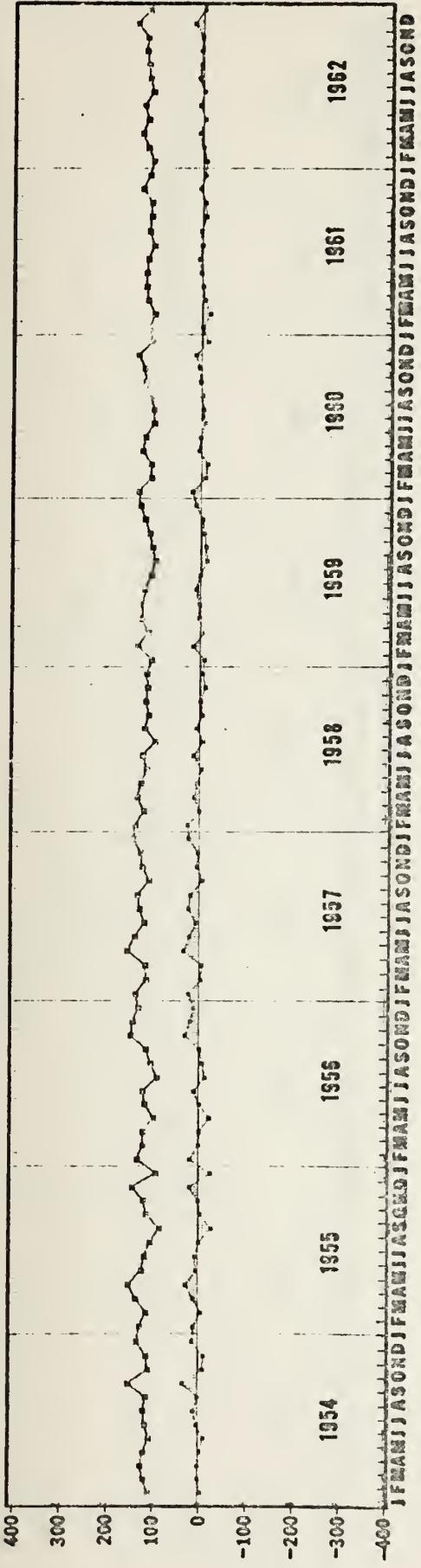


Figure B-2. Fifteen-year time series of the solar radiation and its anomaly ($\text{cal cm}^{-2} \text{ day}^{-1}$).



98

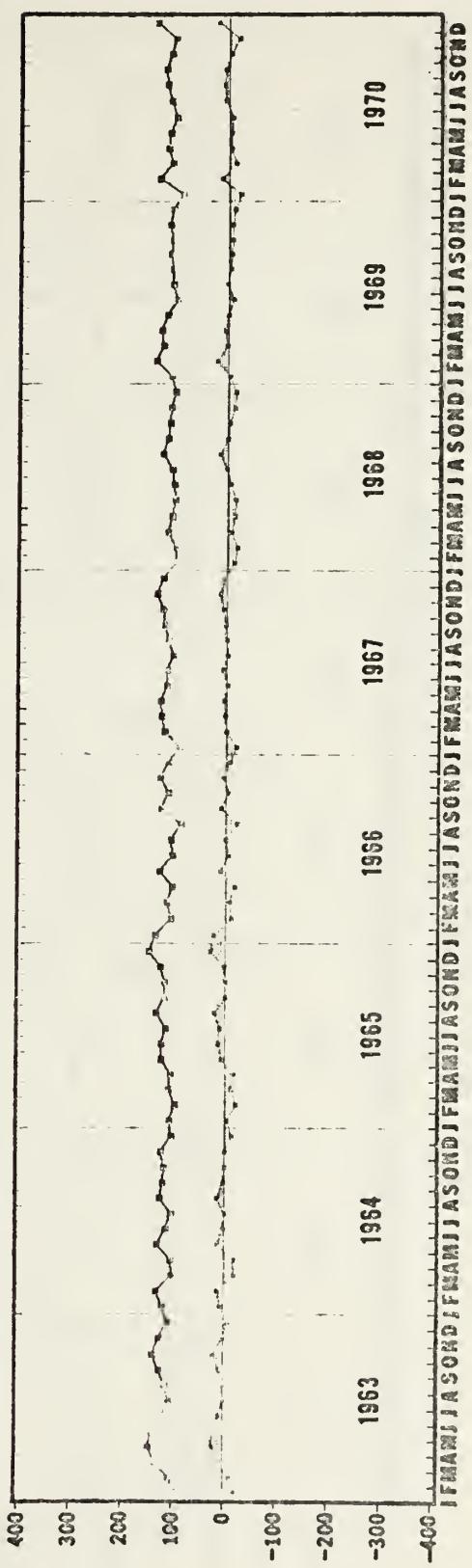


Figure B-3. Fifteen-year time series of the back radiation and its anomaly ($\text{cal cm}^{-2} \text{ day}^{-1}$).

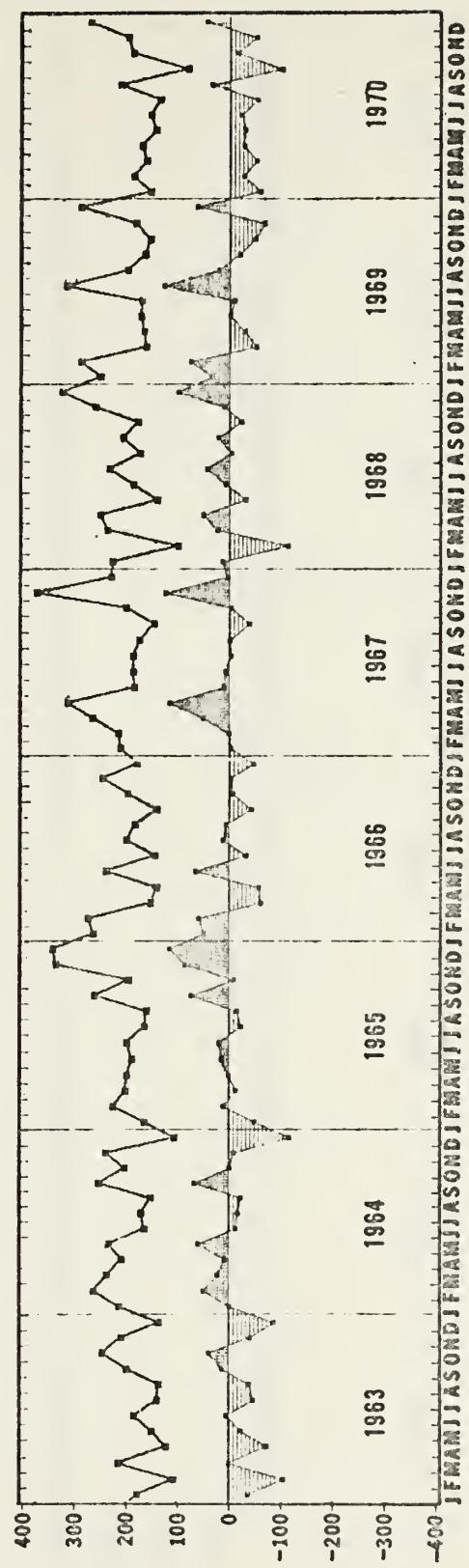
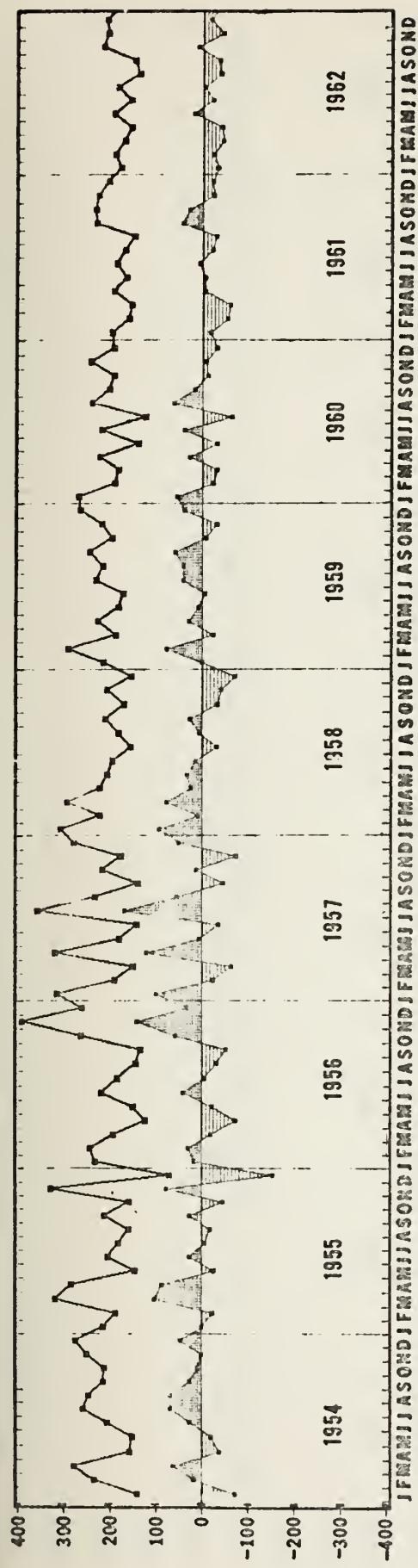


Figure B-4. Fifteen-year time series of the heat of evaporation and its anomaly ($\text{cal cm}^{-2} \text{ day}^{-1}$).

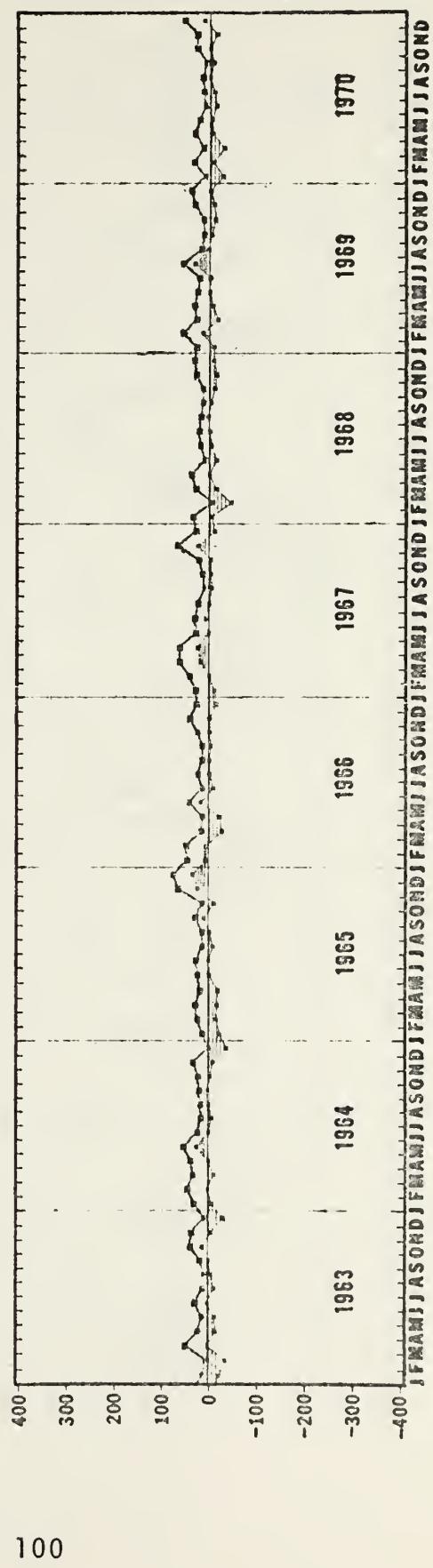
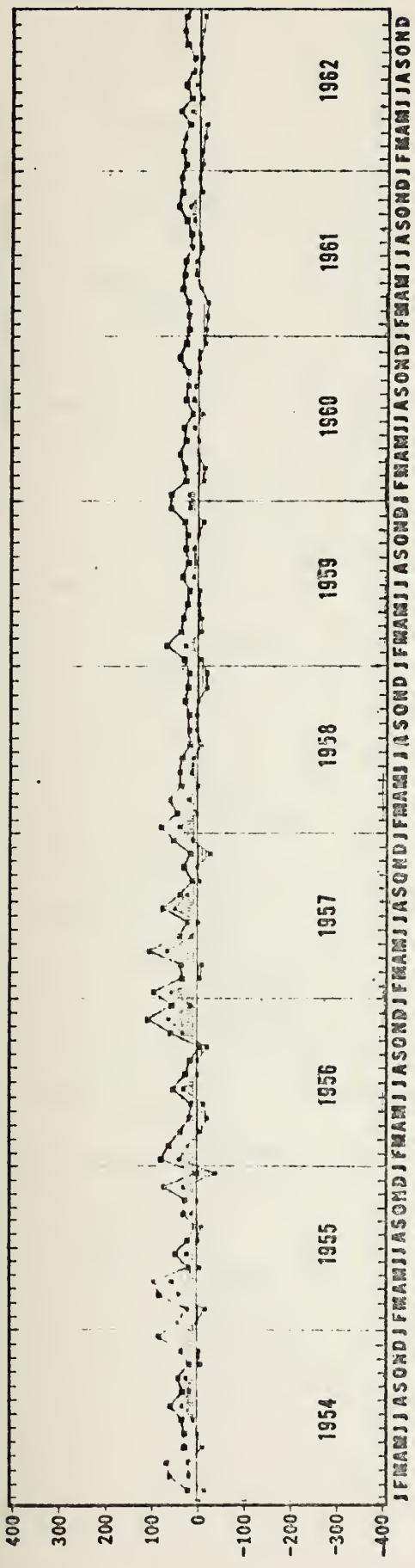


Figure B-5. Fifteen-year time series of the convective heat and its anomaly ($\text{cal cm}^{-2} \text{ day}^{-1}$).

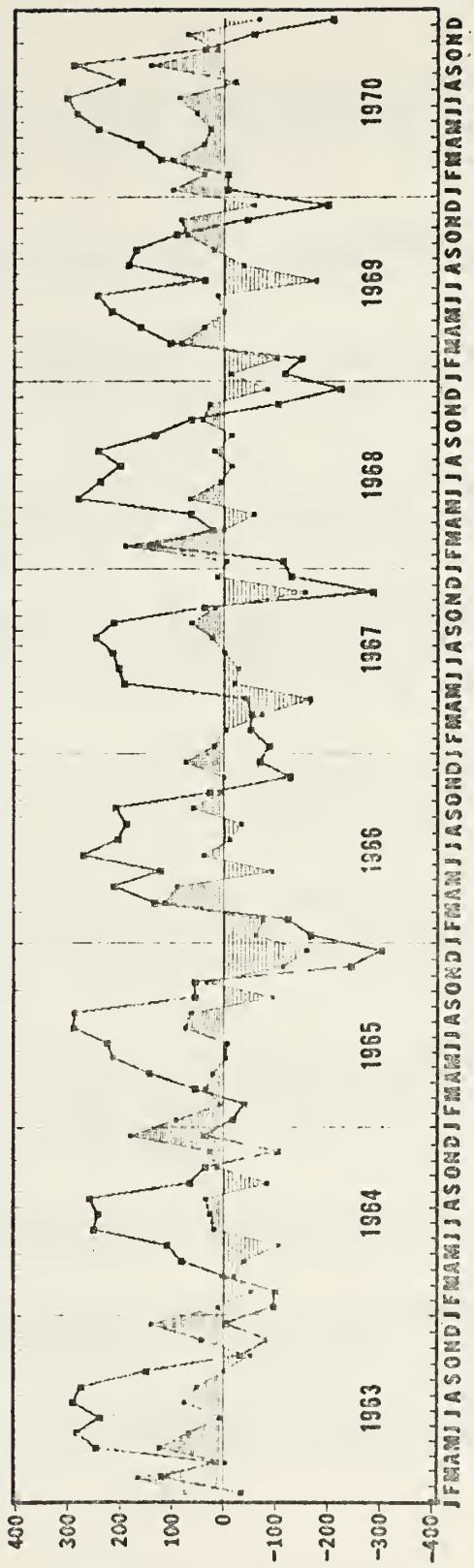
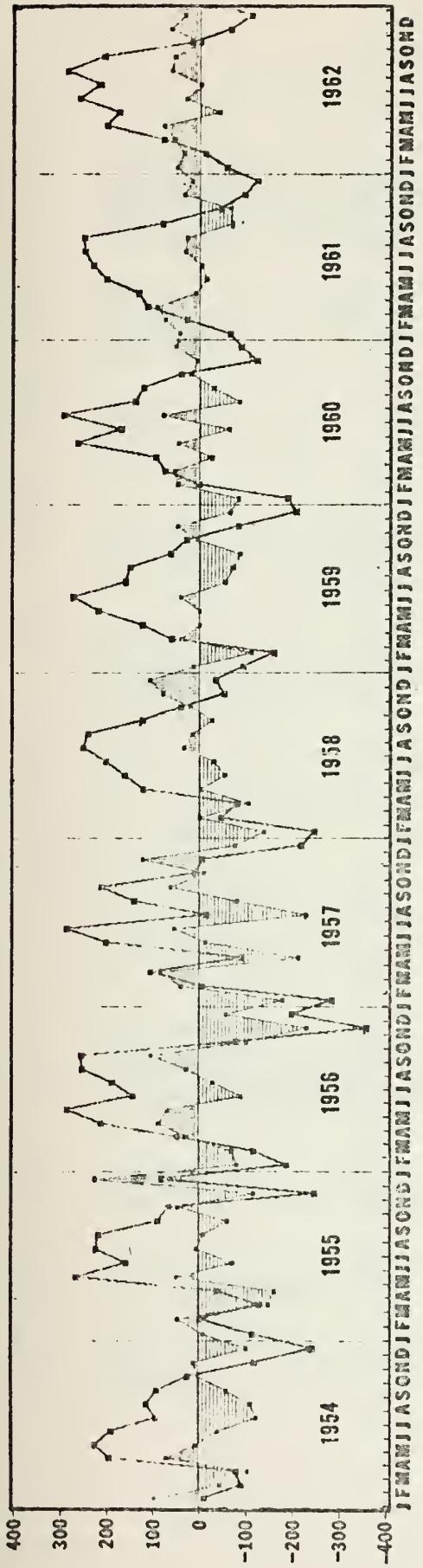


Figure B-6. Fifteen-year time series of the net radiation and its anomaly (cal cm⁻² day⁻¹).

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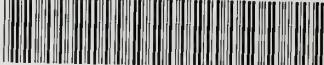
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